

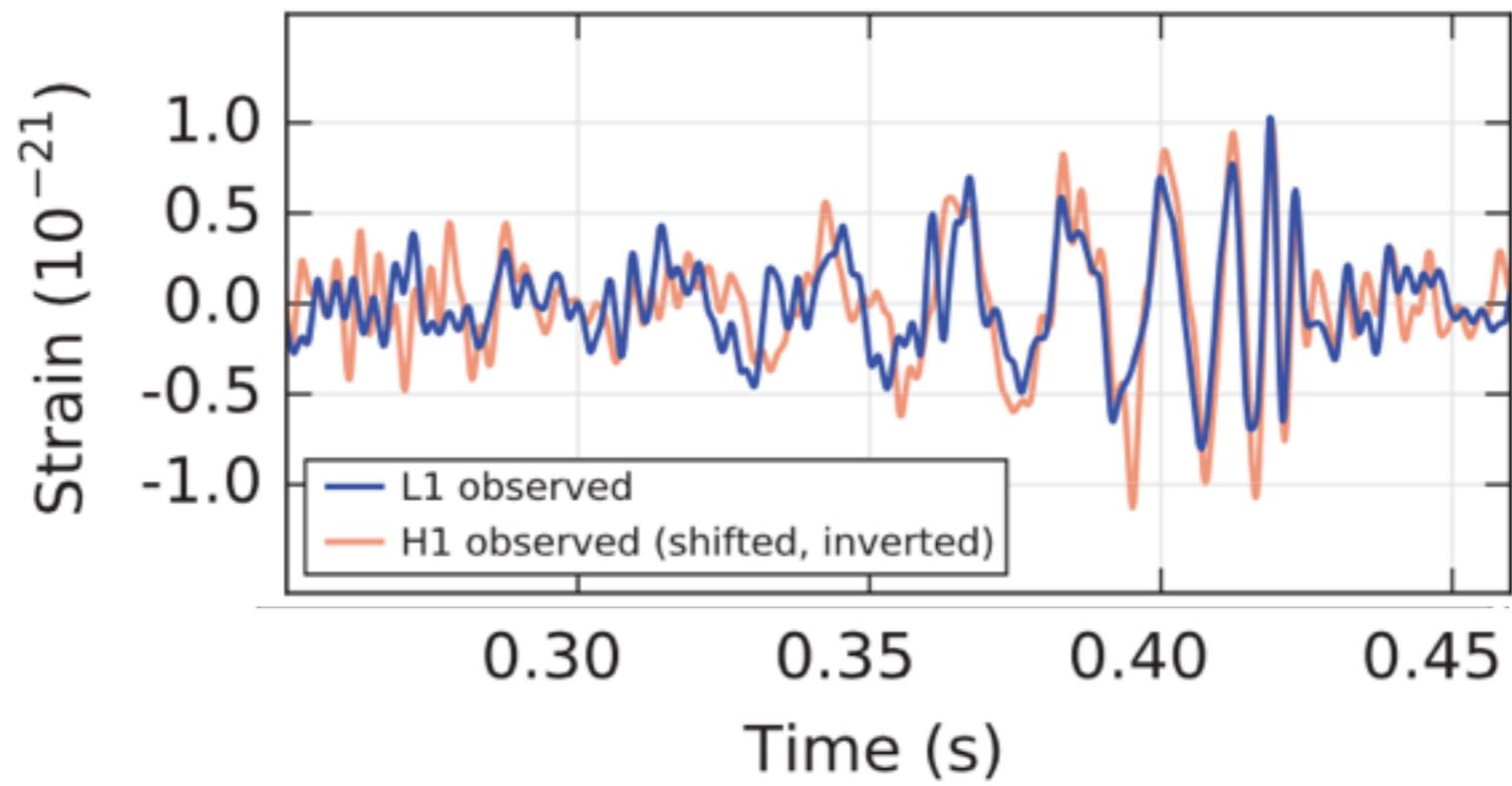
# Waveform modelling with numerical relativity: the key to decode GW150914

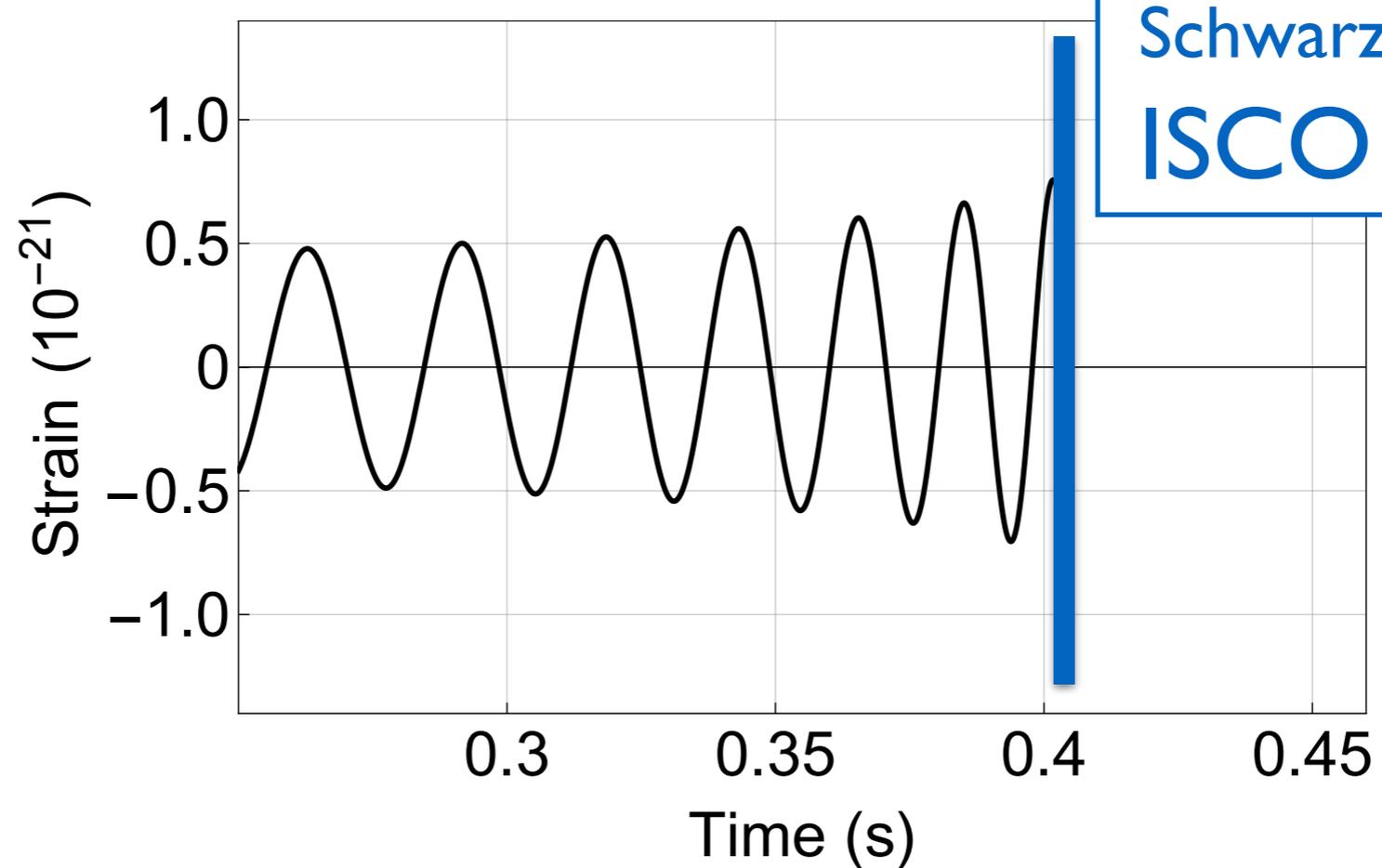
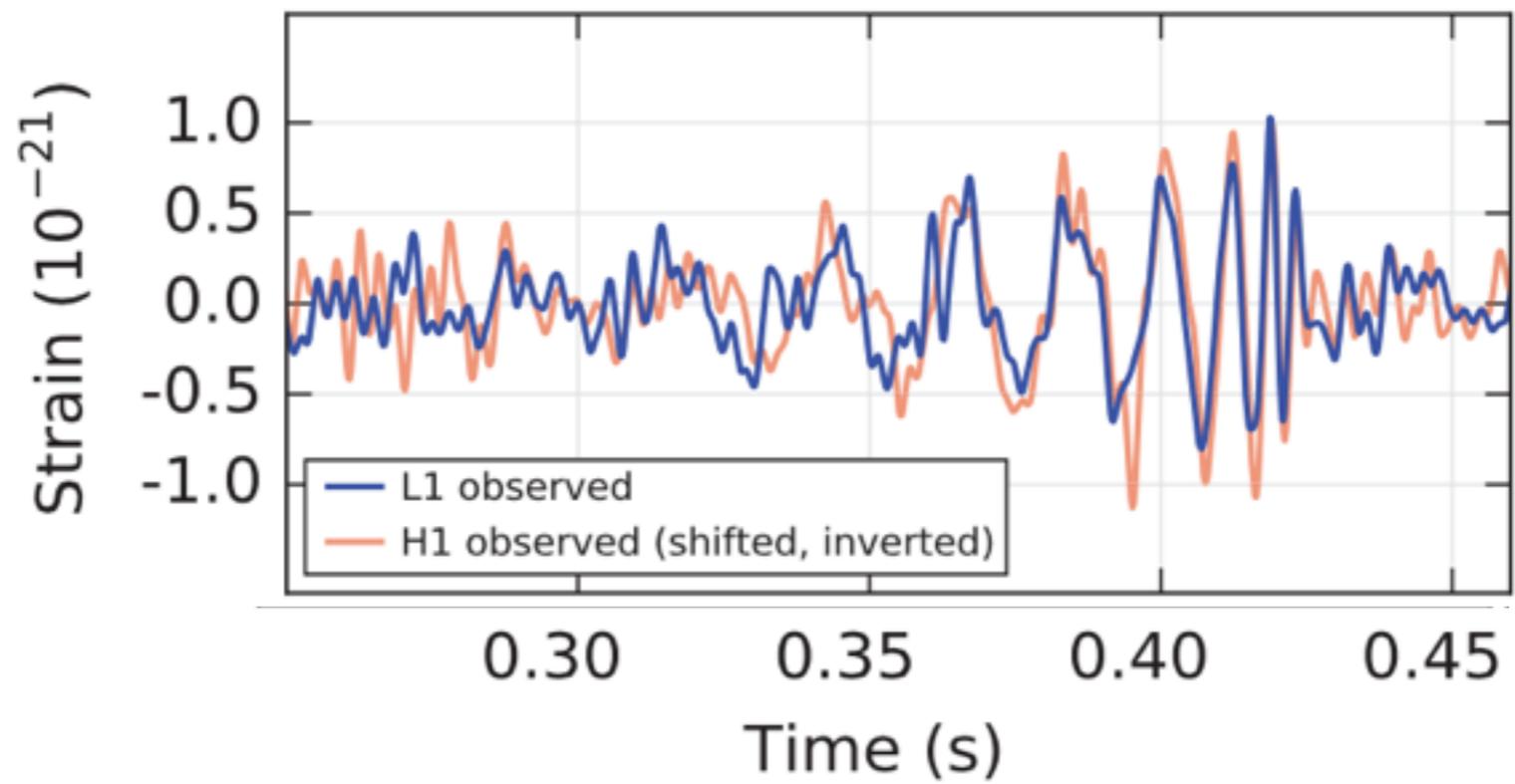
Mark Hannam  
Cardiff University



GW150914 meeting  
Hannover, May 24 2016



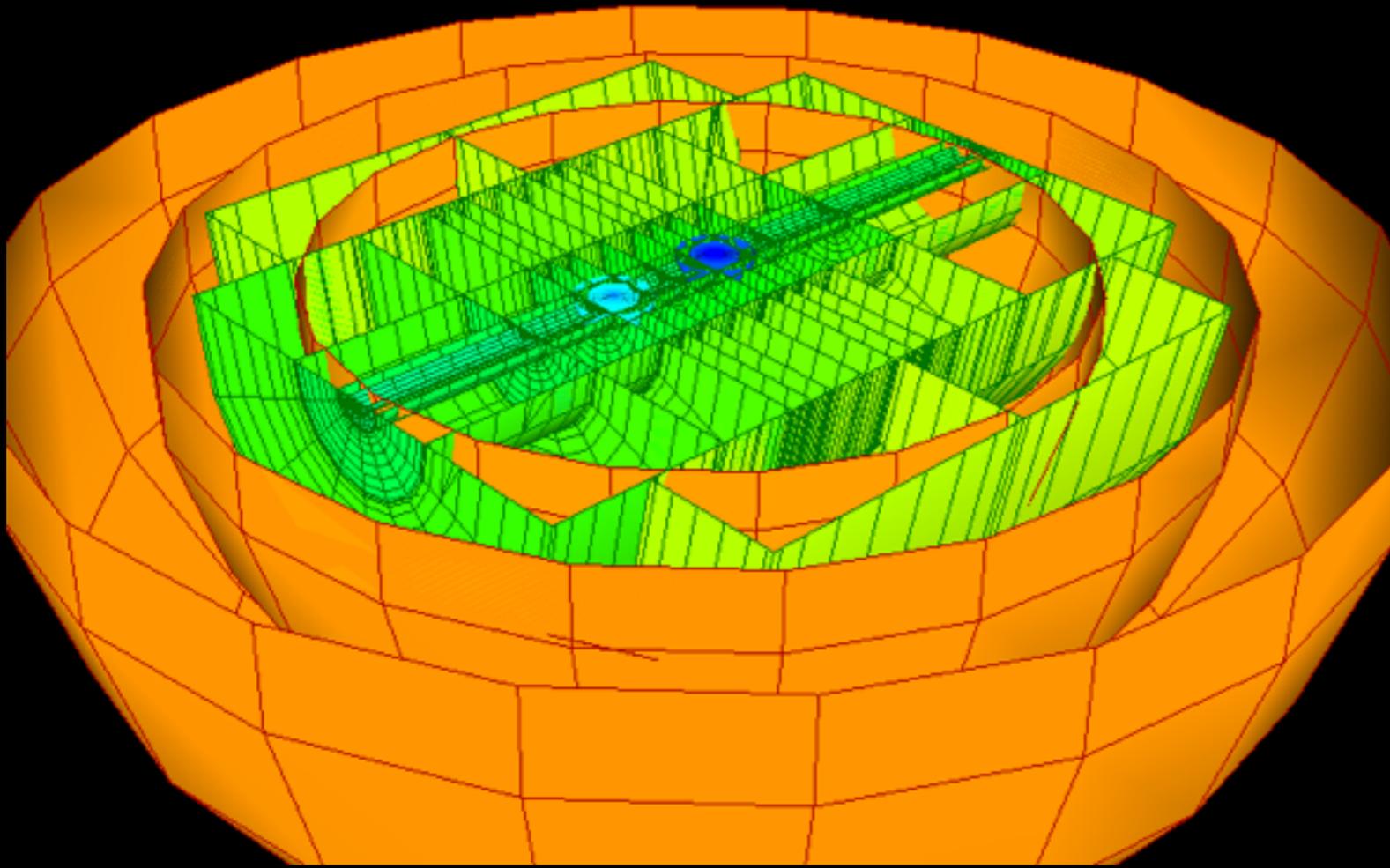




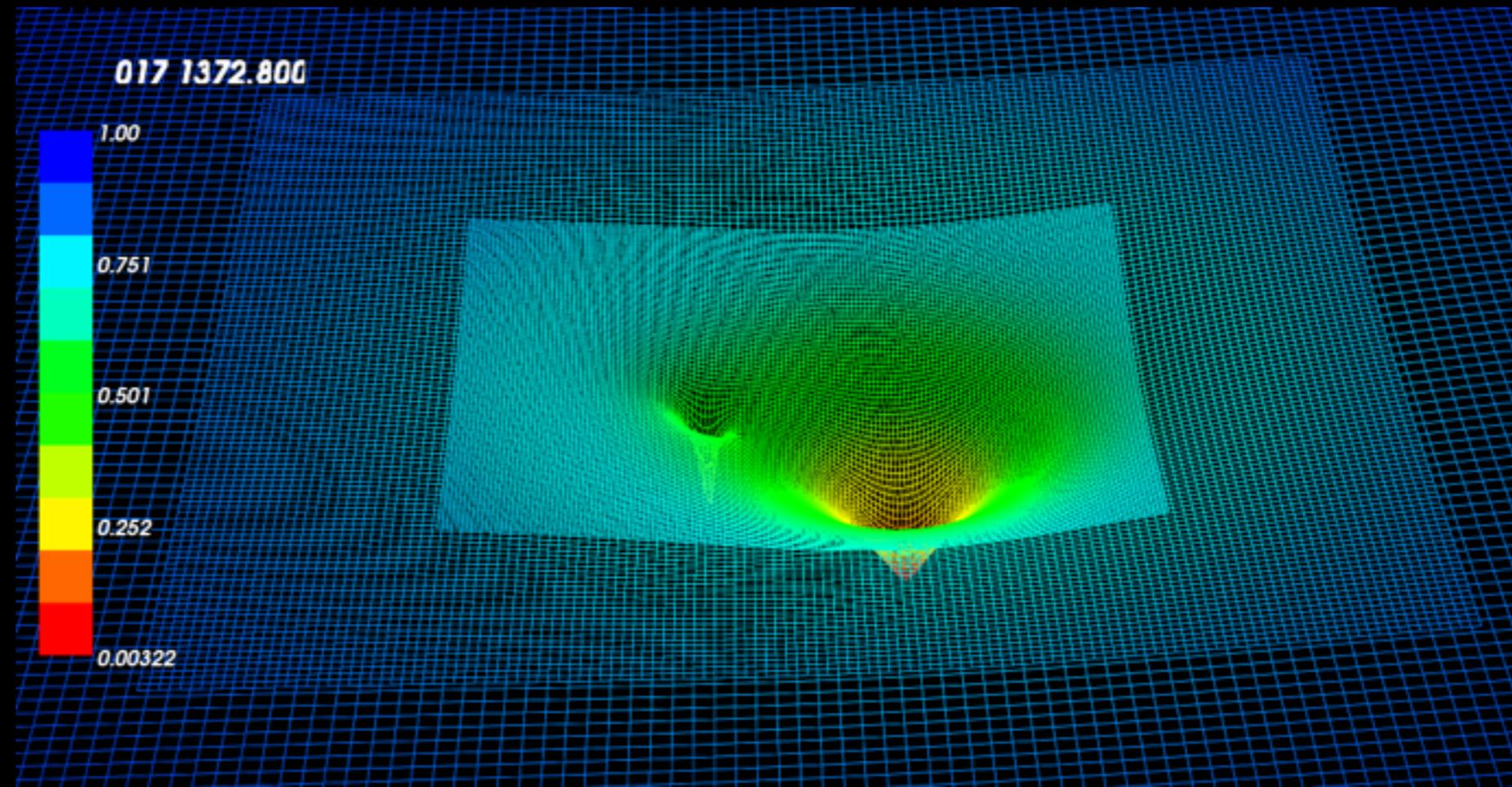
The binary black hole  
problem  
(1964—2005)

**Boosted Three-Dimensional Black-Hole Evolutions with Singularity Excision**

The accurate computational modeling of black-hole interactions is essential to the confident detection of astrophysical gravitational radiation by future space-based detectors such as LISA and by the LIGO/VIRGO/GEO complex of ground-based detectors currently under construction. The sensitivity of these detectors will be significantly enhanced if accurate computer simulations of black-hole mergers can produce predictions of radiation waveforms [1]. The Binary Black Hole Grand Challenge Alliance [2] was funded in September 1993 to develop the computational infrastructure for accurate simulations of the coalescence of black-hole binaries. The primary objective of the resulting code will be the prediction of waveforms from binary black-hole mergers. In this Letter we report on an important step towards achieving such simulations.



Computational  
Cost



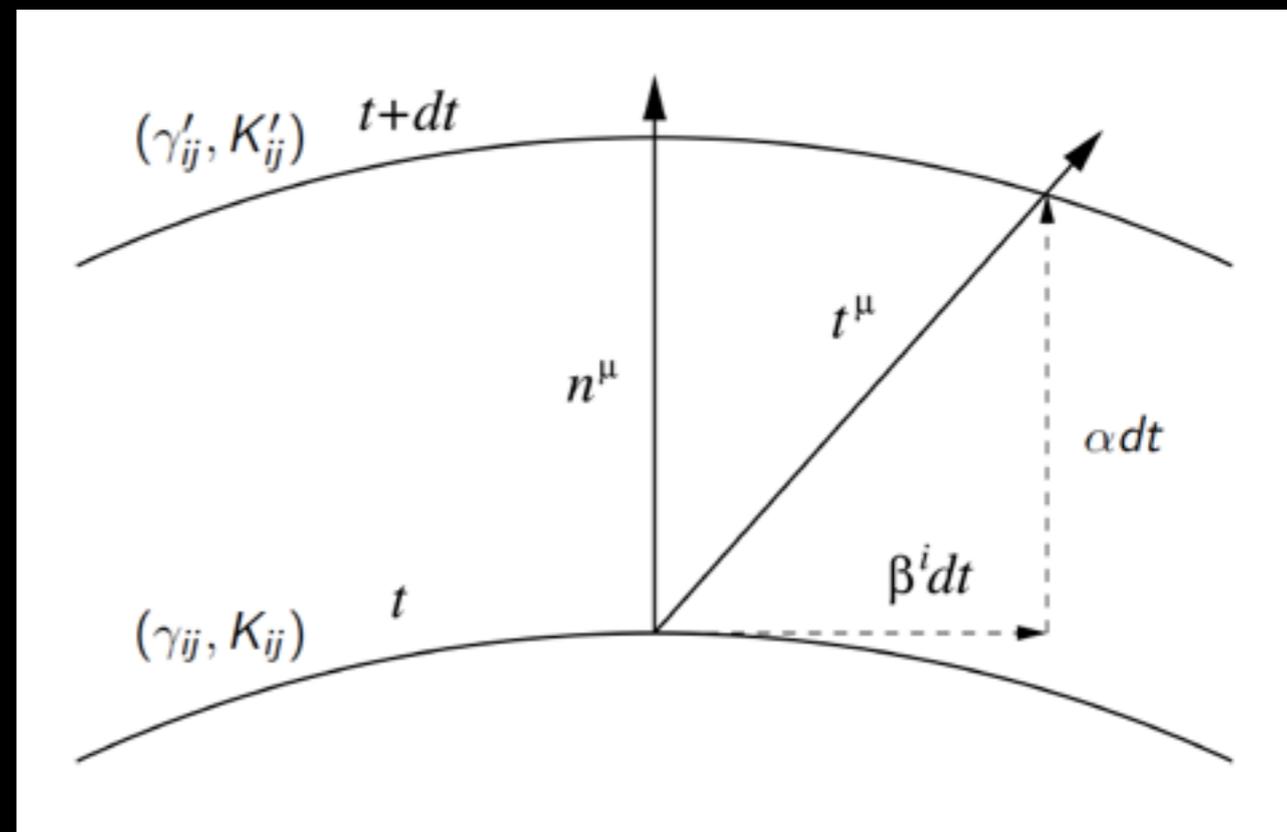
$$G_{\mu\nu} = 8\pi T_{\mu\nu}$$

Numerically  
stable  
formulations

$$G_{\mu\nu} = 8\pi T_{\mu\nu}$$



(3+1 split)



$$G_{\mu\nu} = 8\pi T_{\mu\nu}$$



(3+1 split)

$$\begin{aligned}\bar{R} + K^2 - K_{ij}K^{ij} &= 0, \\ \bar{\nabla}_j K^{ij} - \gamma^{ij}\bar{\nabla}_j K &= 0.\end{aligned}$$

$$\begin{aligned}\partial_t \gamma_{ij} &= -2NK_{ij} + \bar{\nabla}_i \beta_j + \bar{\nabla}_j \beta_i, \\ \partial_t K_{ij} &= -\bar{\nabla}_i \bar{\nabla}_j N + N(R_{ij} - 2K_{ik}K_j^k + KK_{ij}) \\ &\quad + \beta^k \bar{\nabla}_k K_{ij} + K_{ik} \bar{\nabla}_j \beta^k + K_{kj} \bar{\nabla}_i \beta^k.\end{aligned}$$

[ADM (1962), York (1979)]

$$G_{\mu\nu} = 8\pi T_{\mu\nu}$$



(3+1 split)

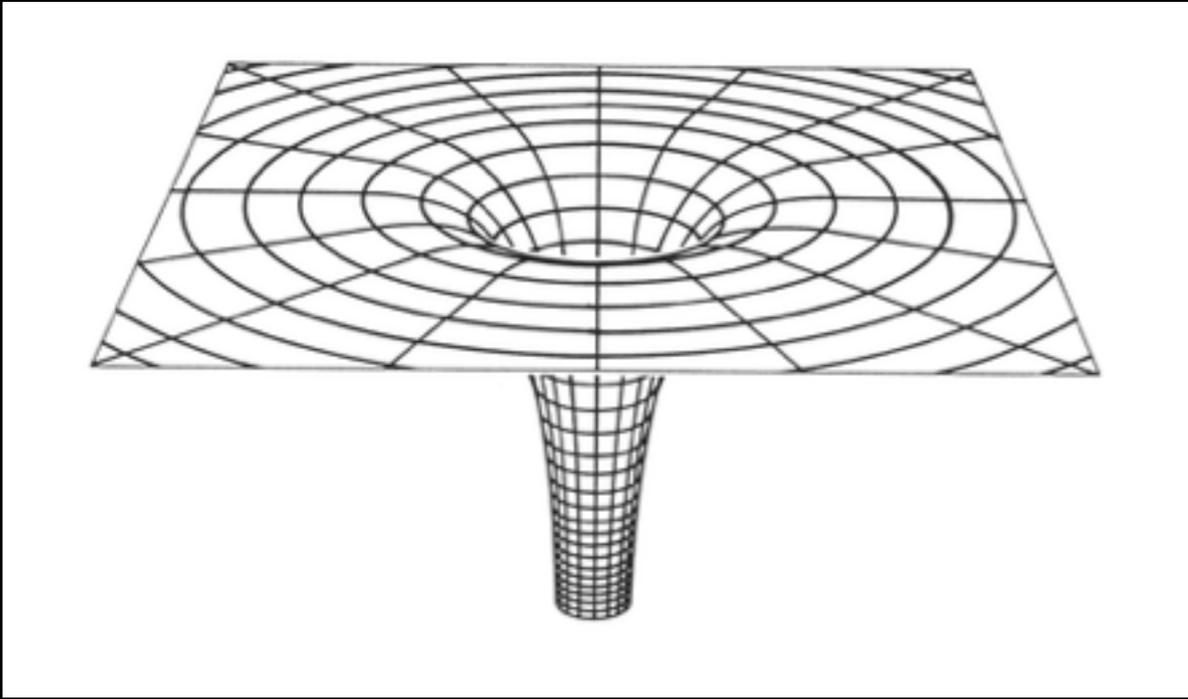
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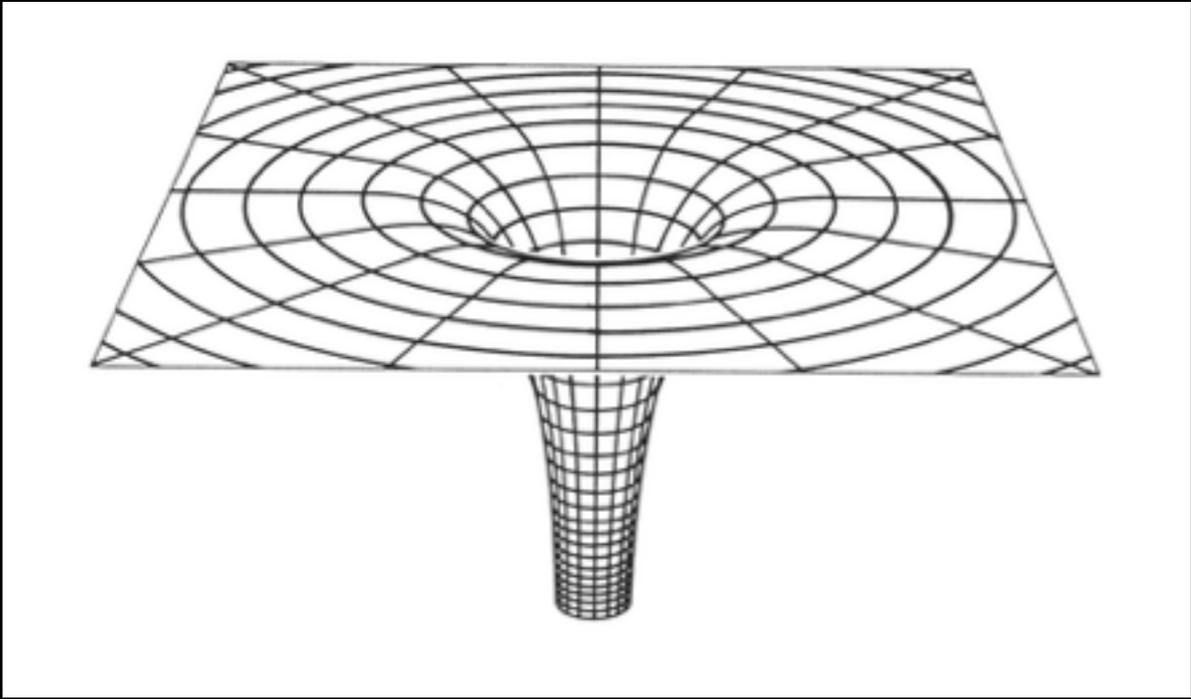


???

[ADM (1962), York (1979)]



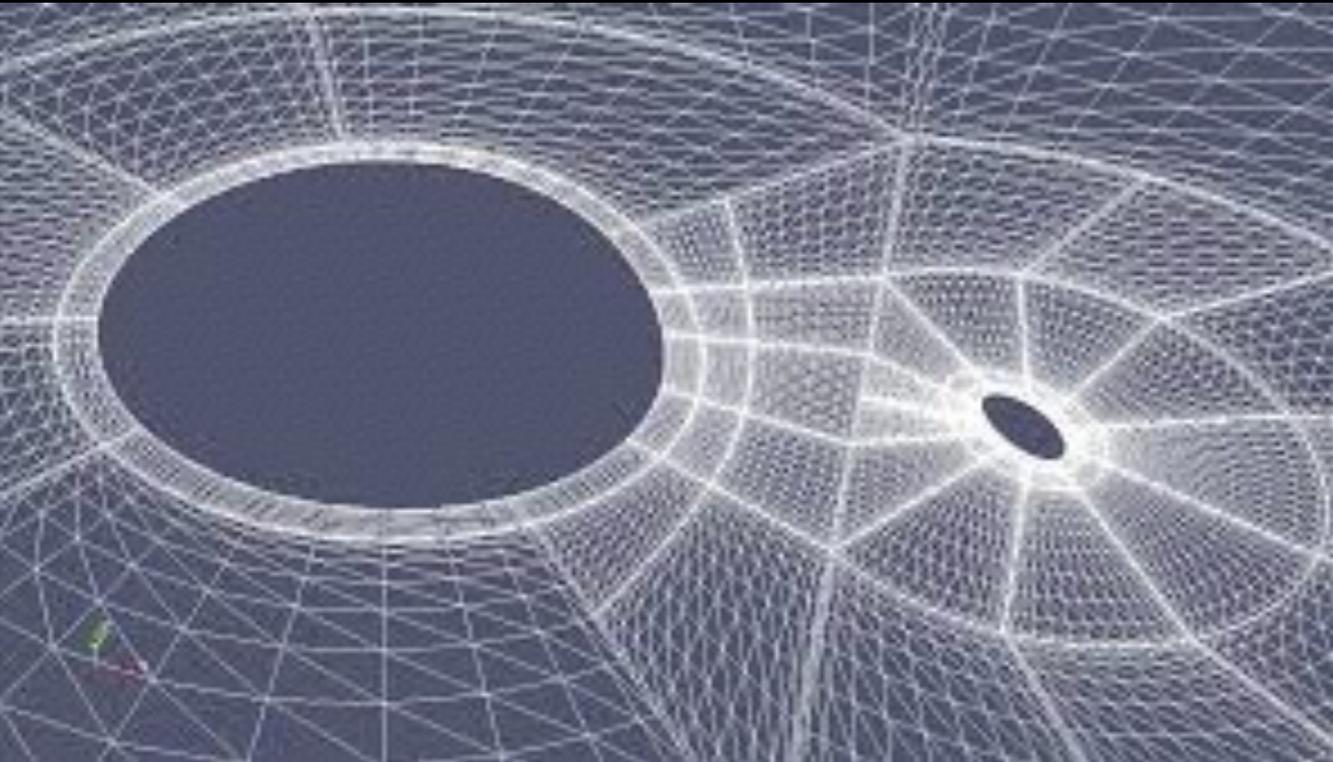
**Black-hole  
singularities**

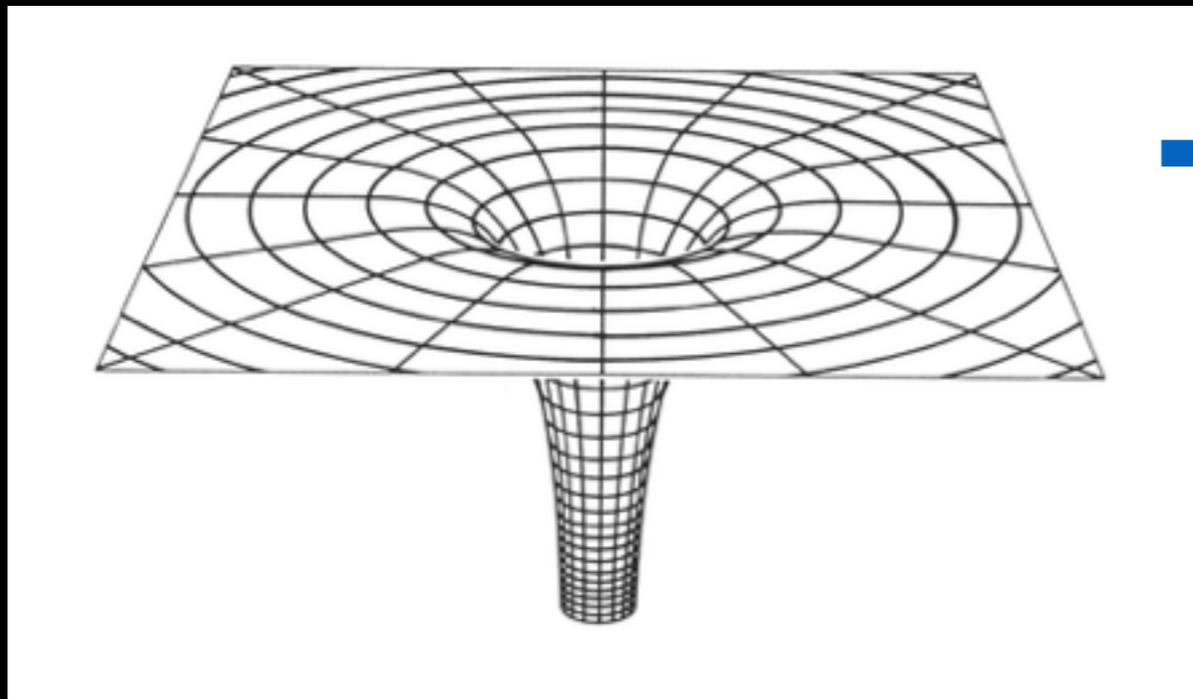


Black-hole  
singularities

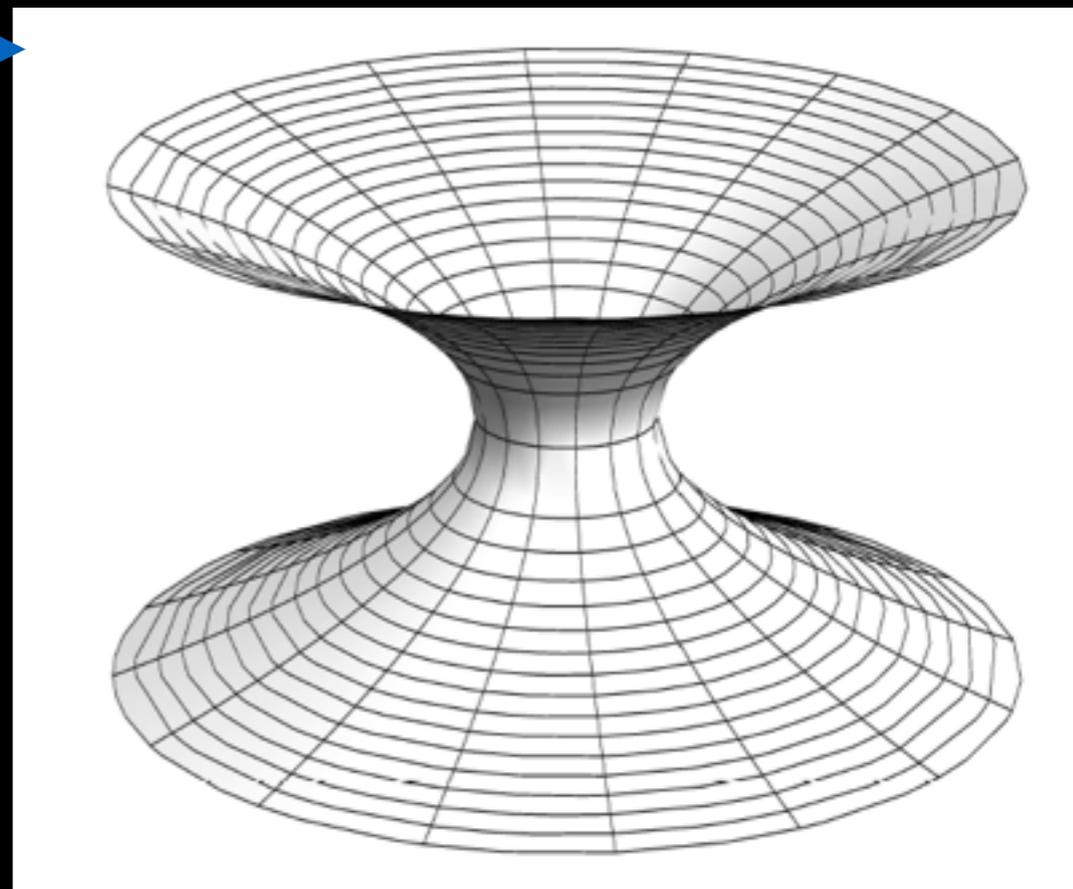


Excision

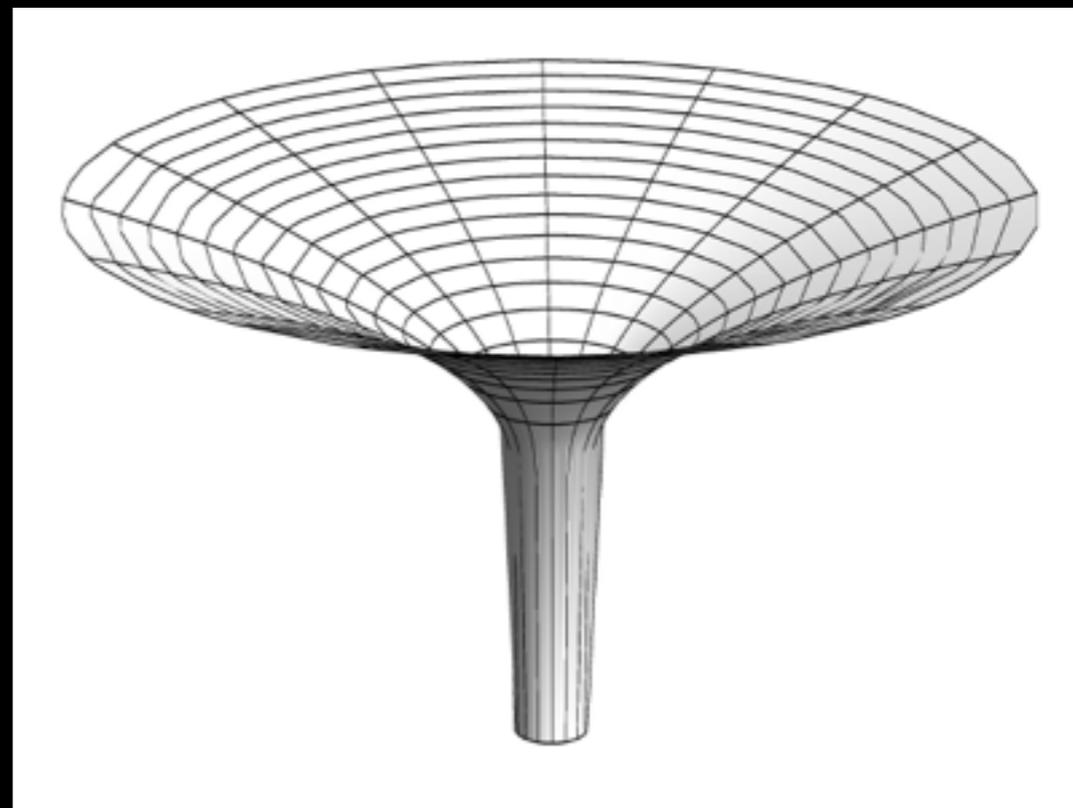
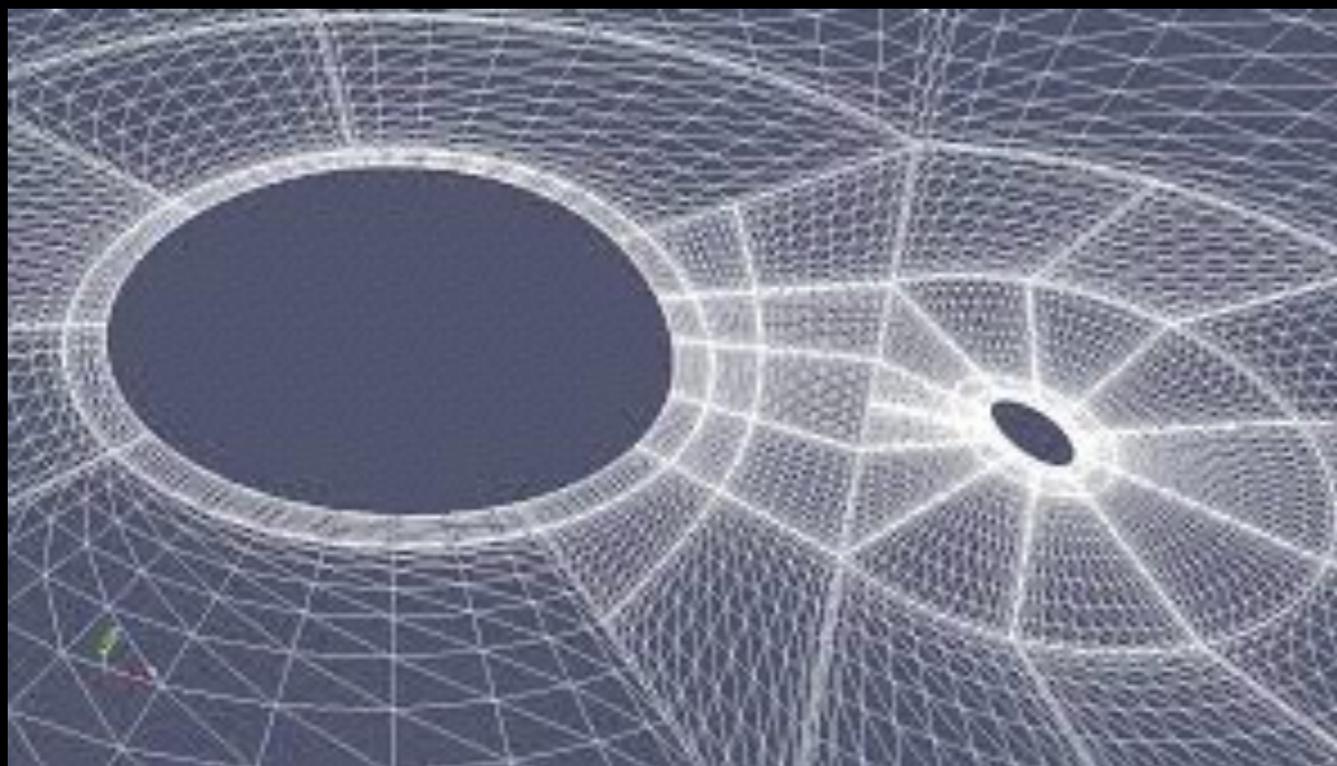




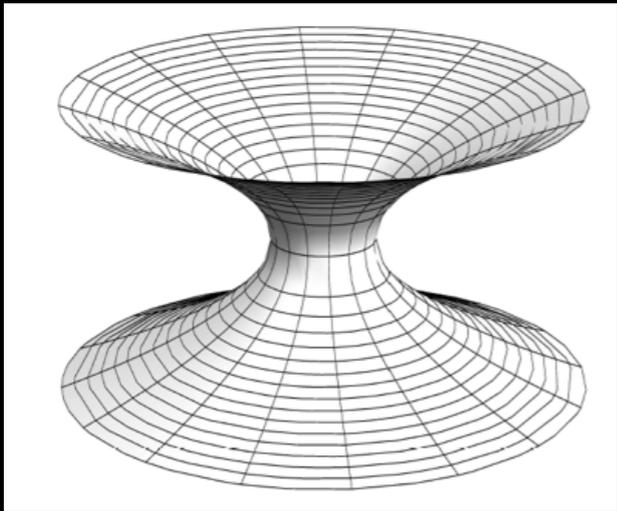
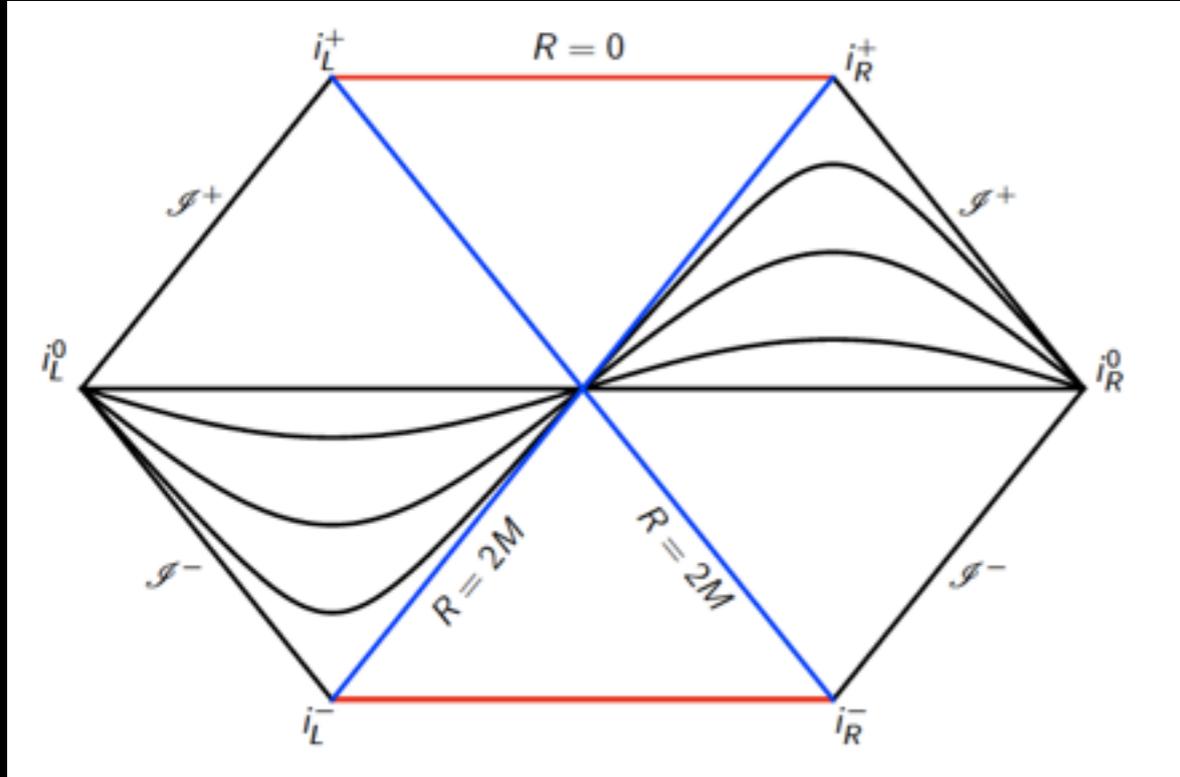
“punctures”



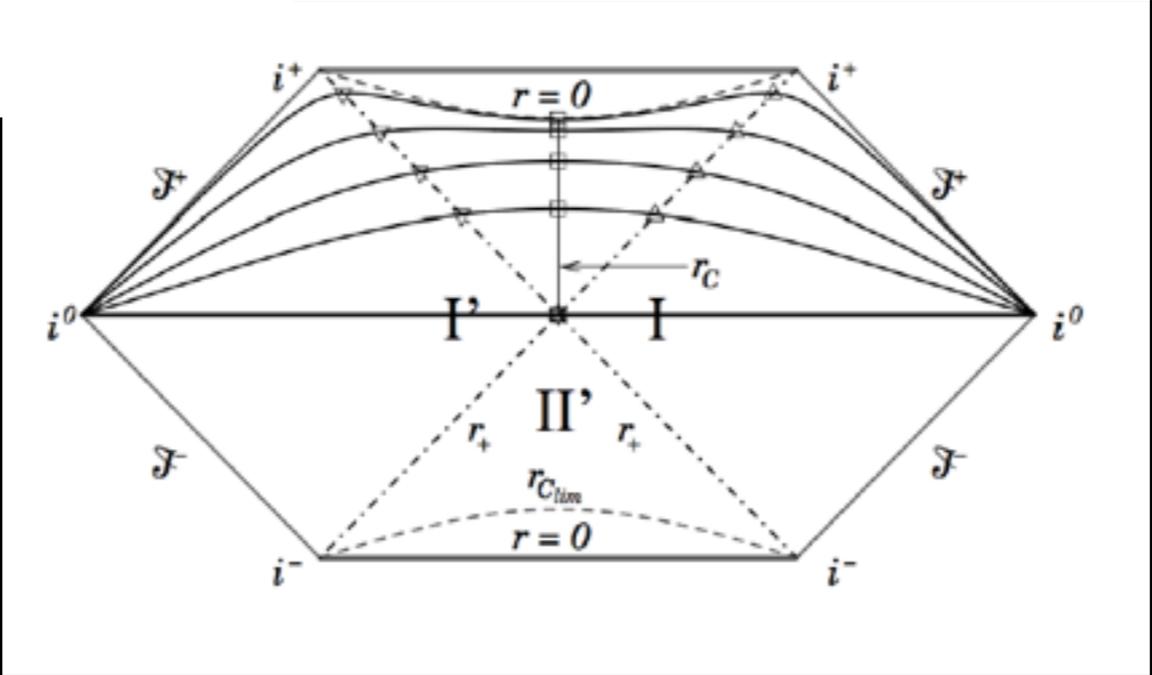
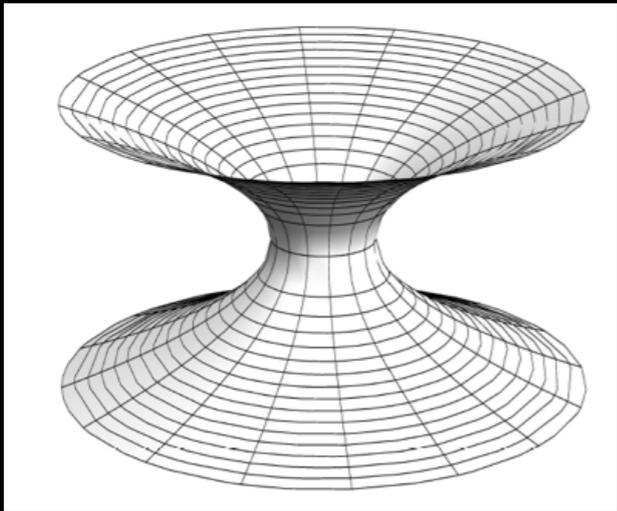
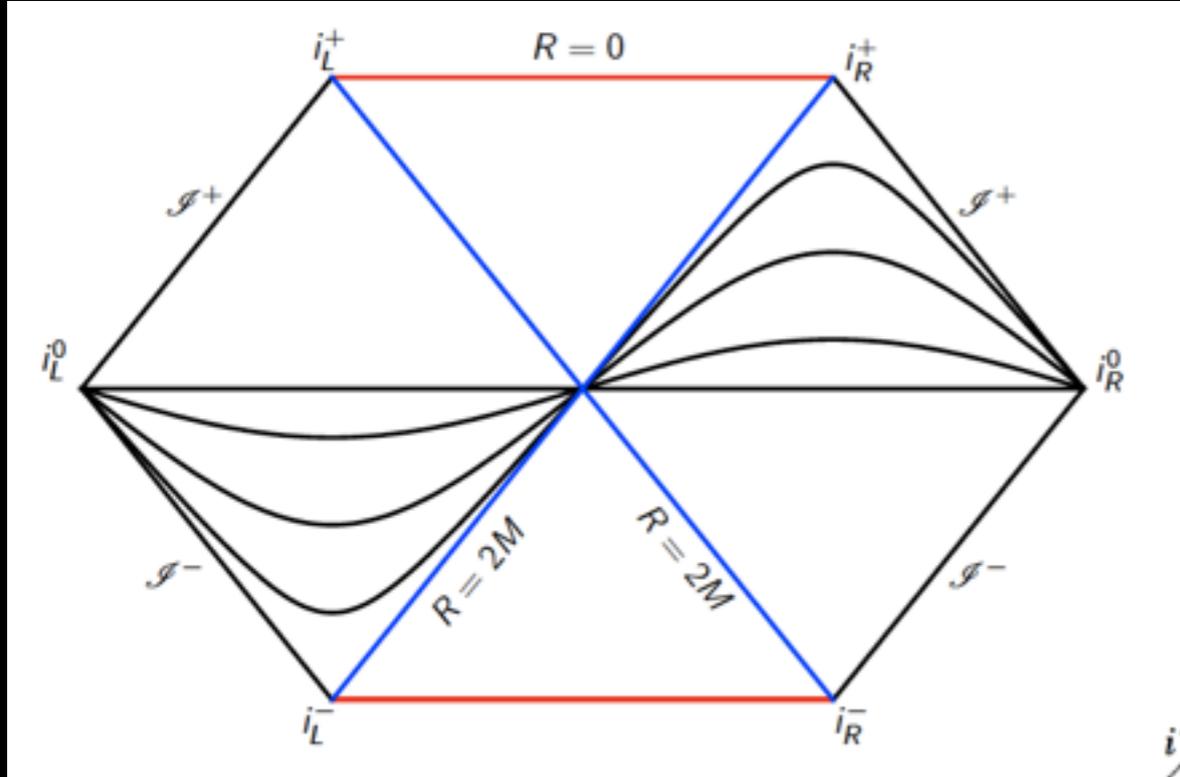
Excision



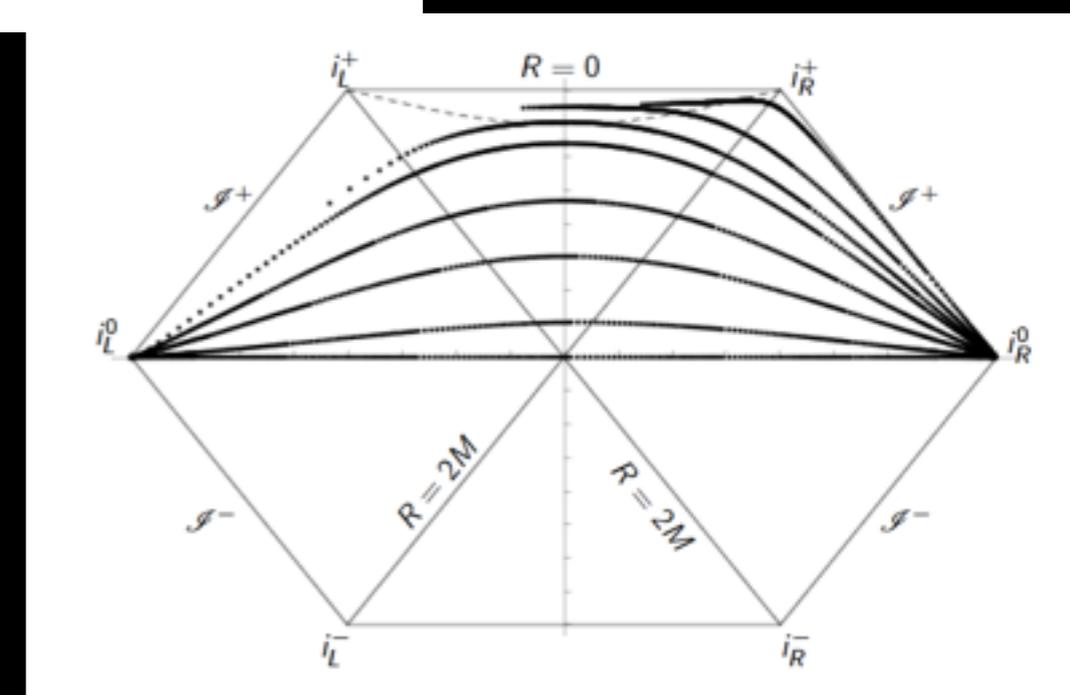
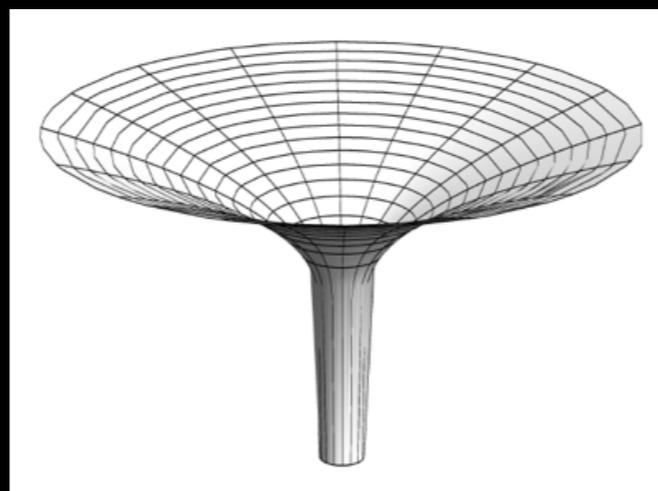
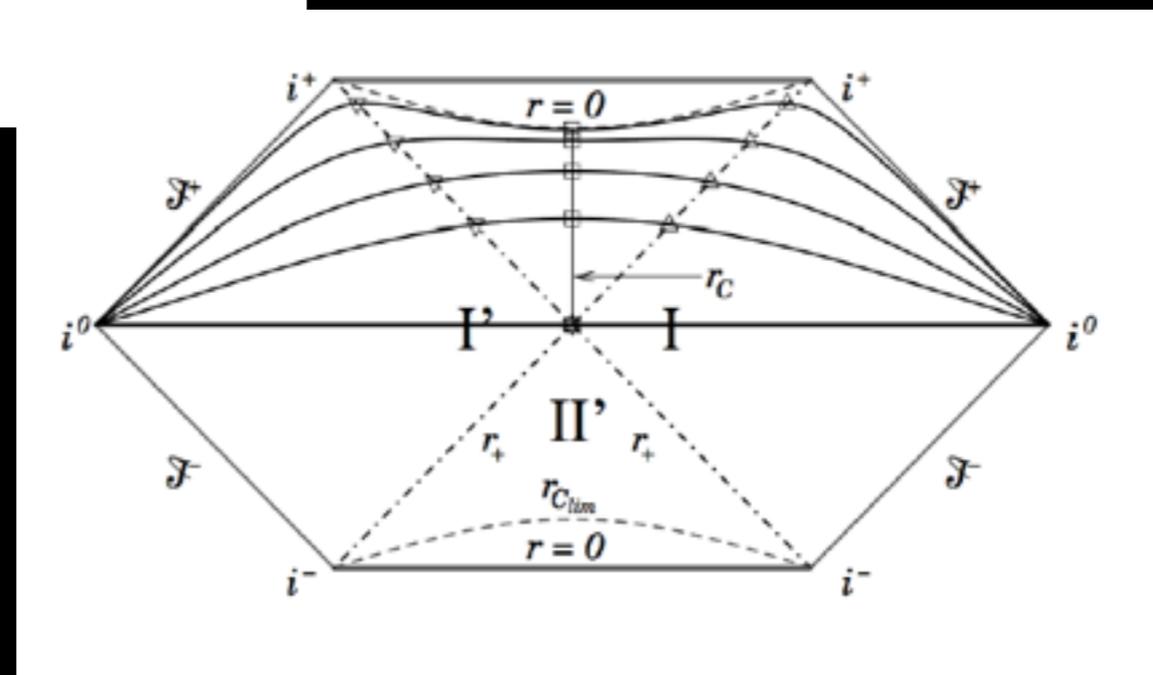
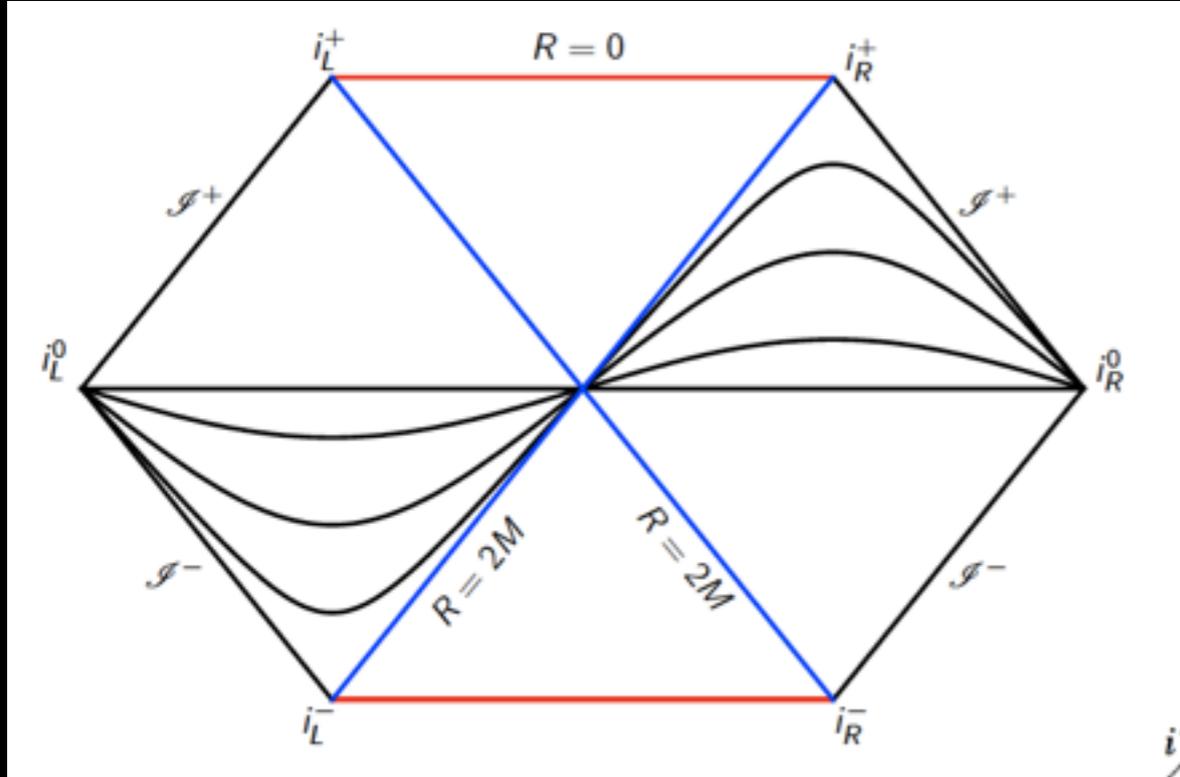
# Gauge conditions

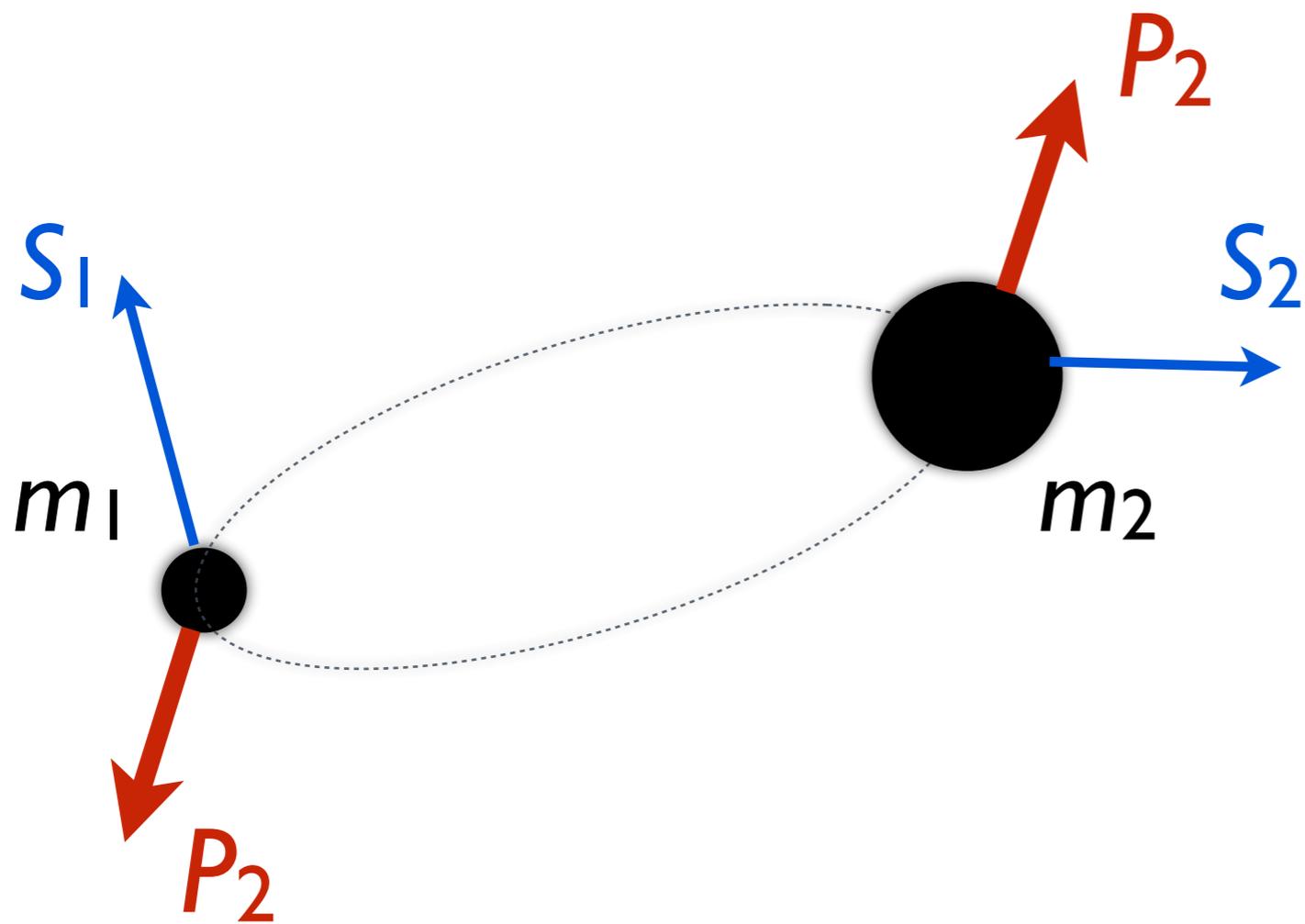


# Gauge conditions

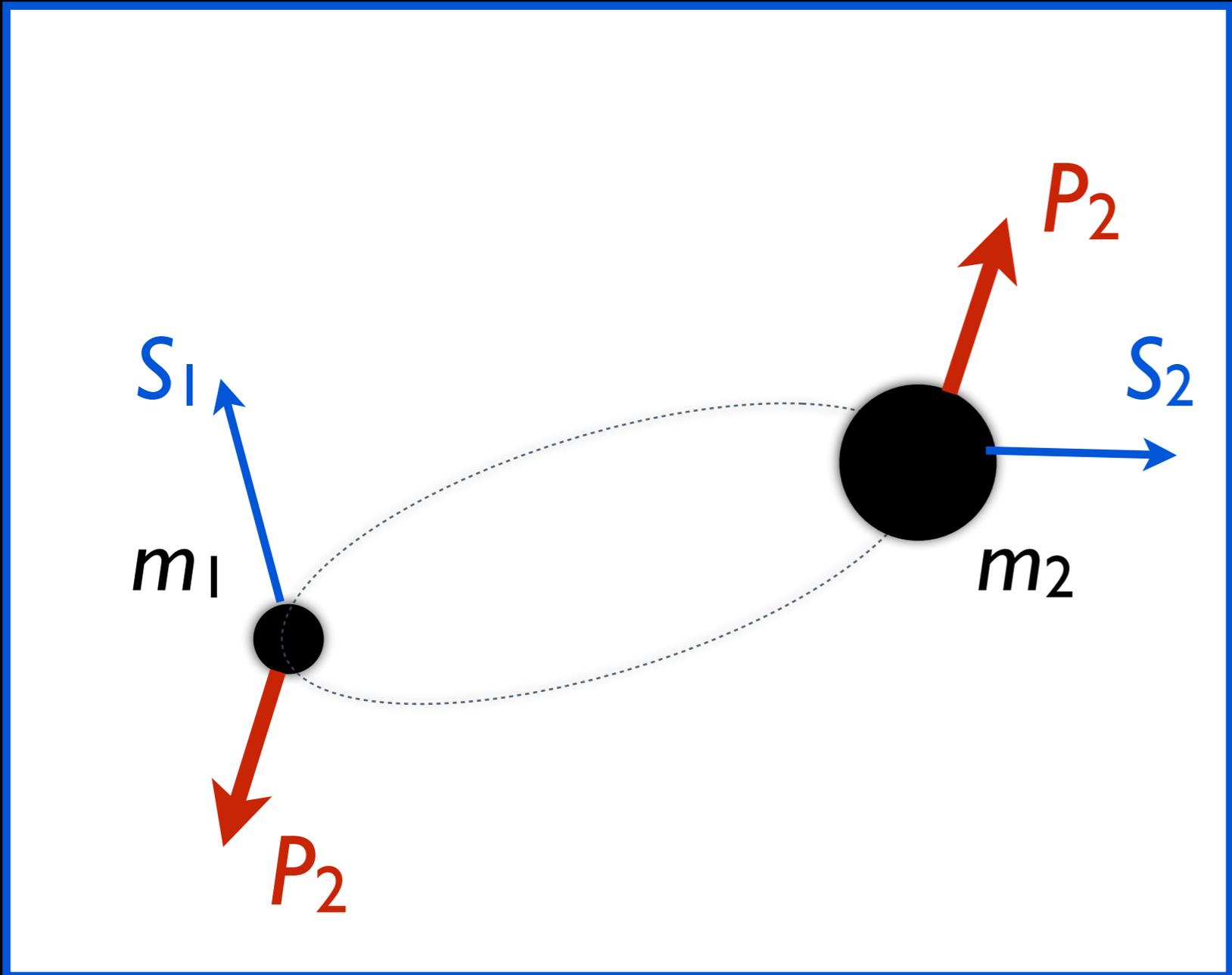


# Gauge conditions

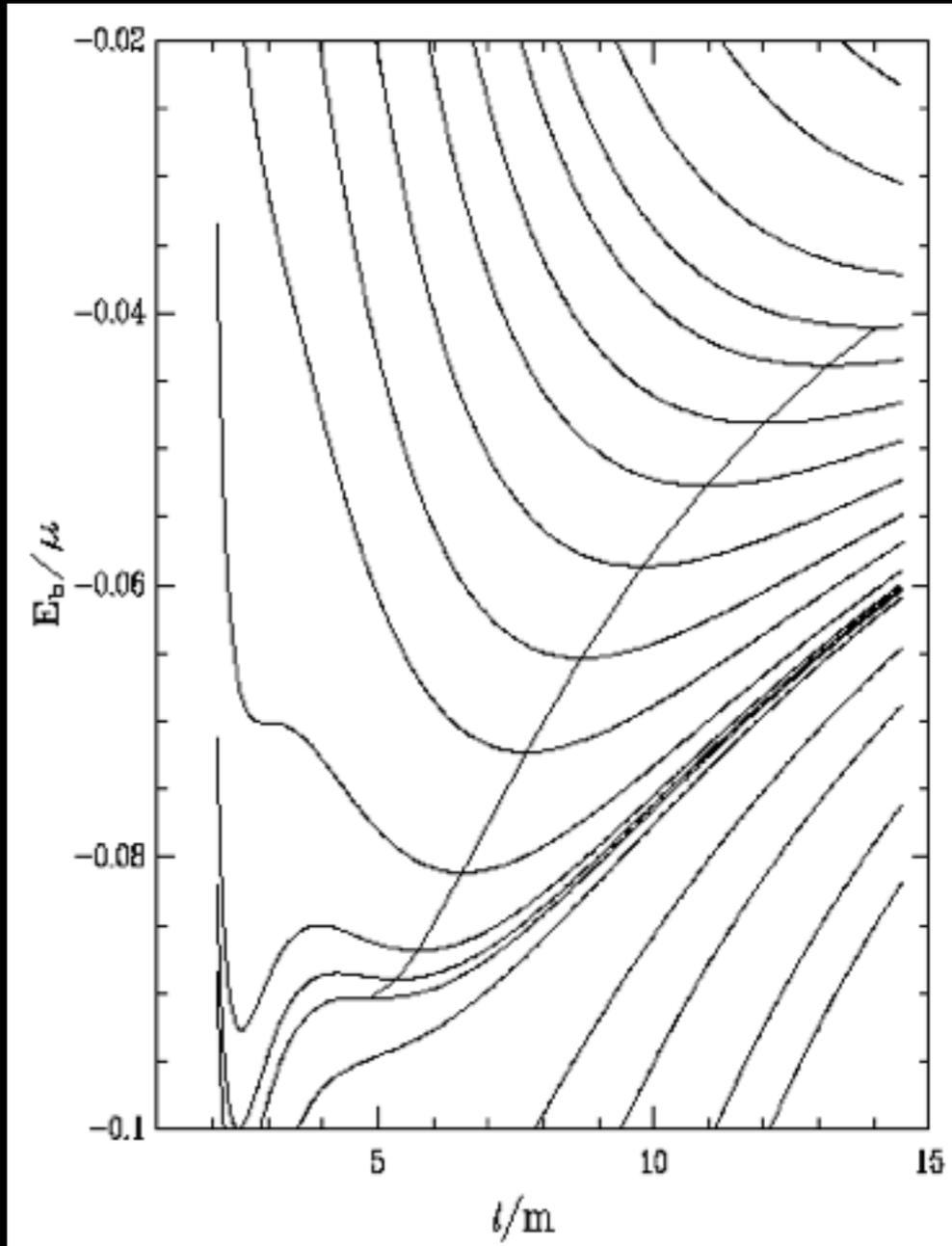




Initial  
conditions



# Initial conditions



**Boosted Three-Dimensional Black-Hole Evolutions with Singularity Excision**

The accurate computational modeling of black-hole interactions is essential to the confident detection of astrophysical gravitational radiation by future space-based detectors such as LISA and by the LIGO/VIRGO/GEO complex of ground-based detectors currently under construction. The sensitivity of these detectors will be significantly enhanced if accurate computer simulations of black-hole mergers can produce predictions of radiation waveforms [1]. The Binary Black Hole Grand Challenge Alliance [2] was funded in September 1993 to develop the computational infrastructure for accurate simulations of the coalescence of black-hole binaries. The primary objective of the resulting code will be the prediction of waveforms from binary black-hole mergers. In this Letter we report on an important step towards achieving such simulations.

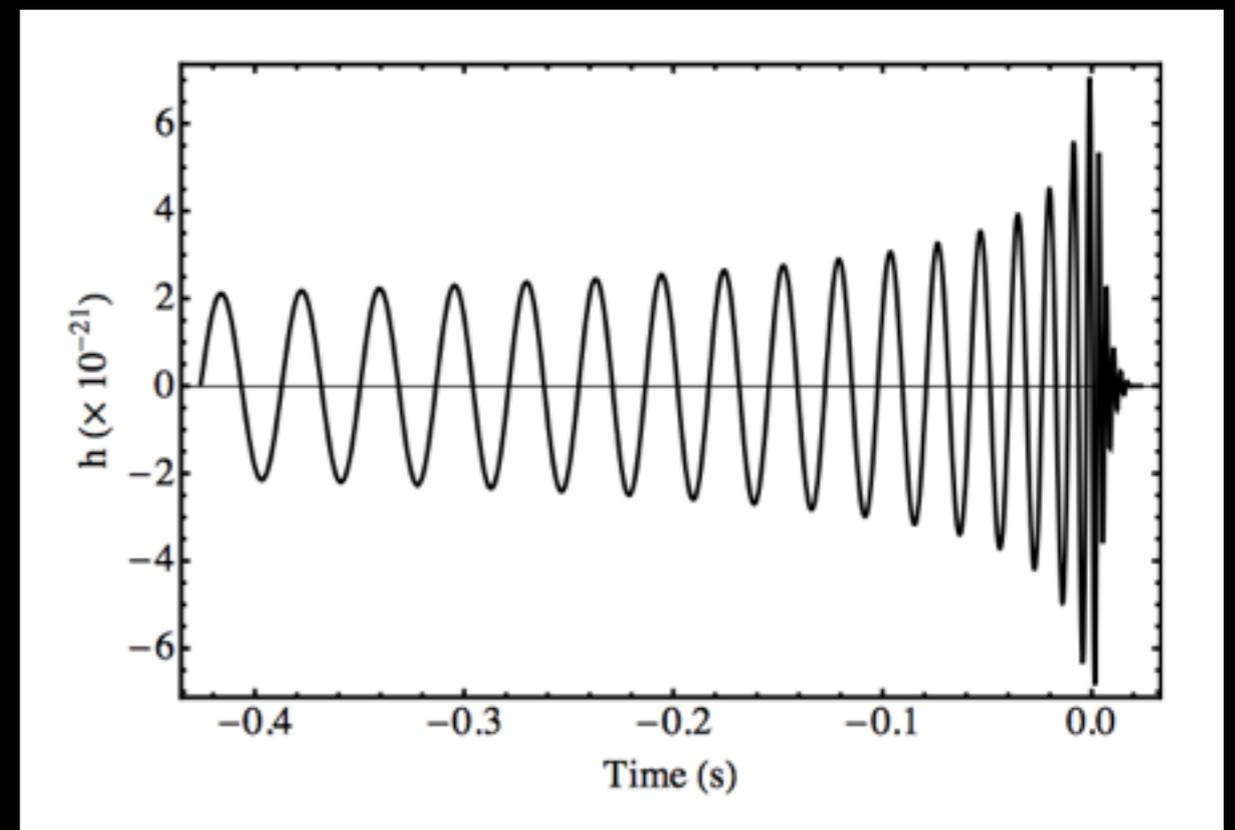
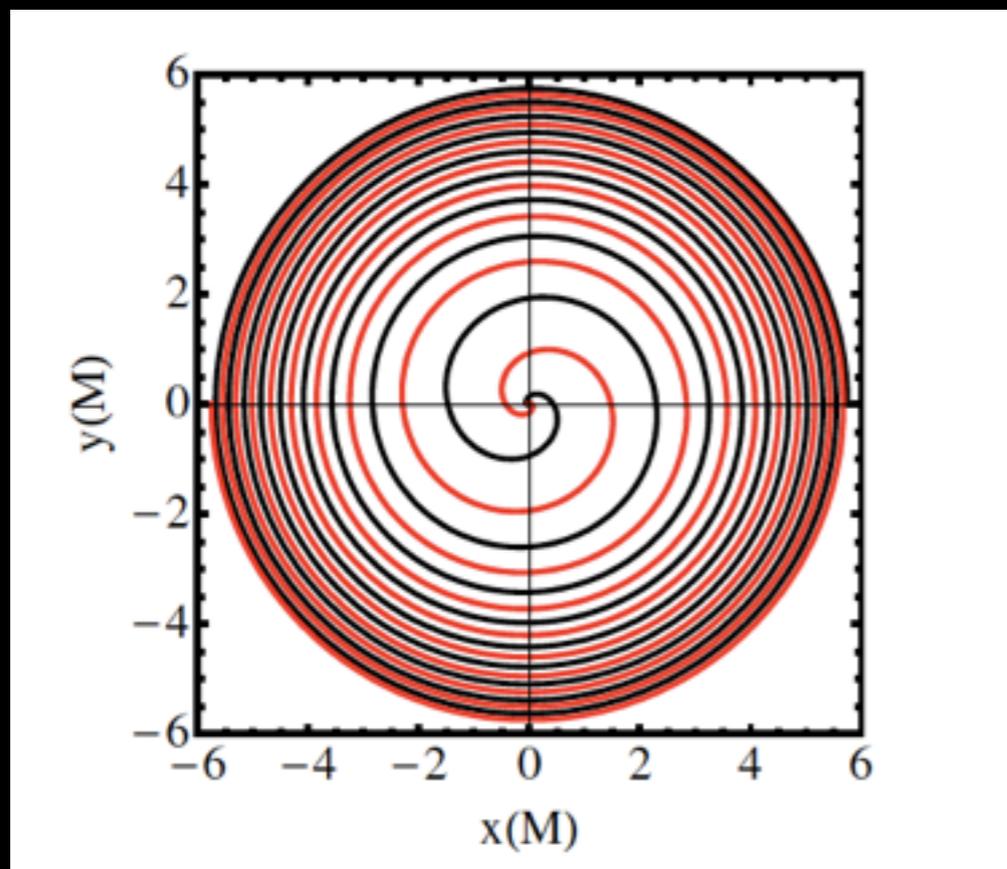
## Boosted Three-Dimensional Black-Hole Evolutions with Singularity Excision

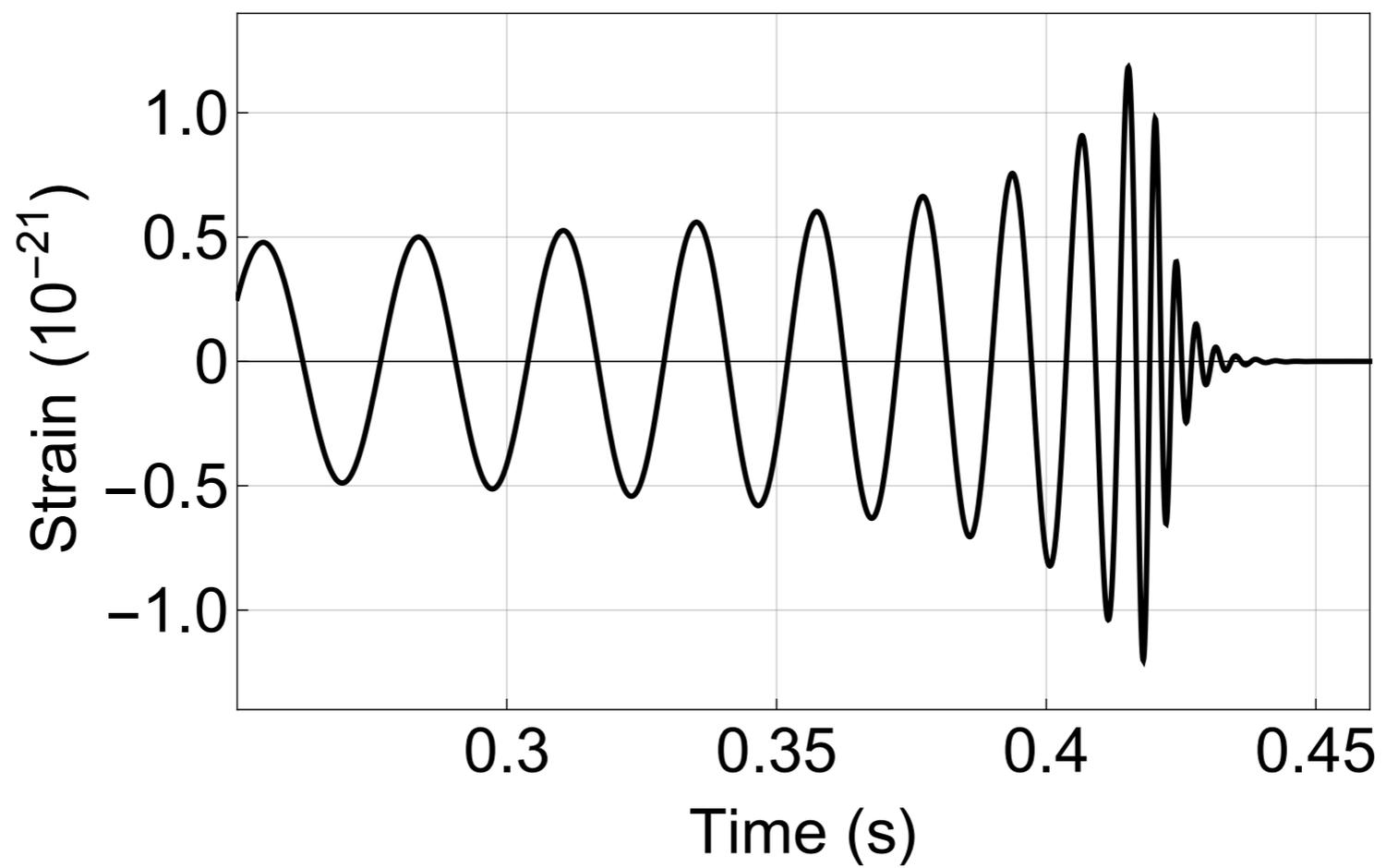
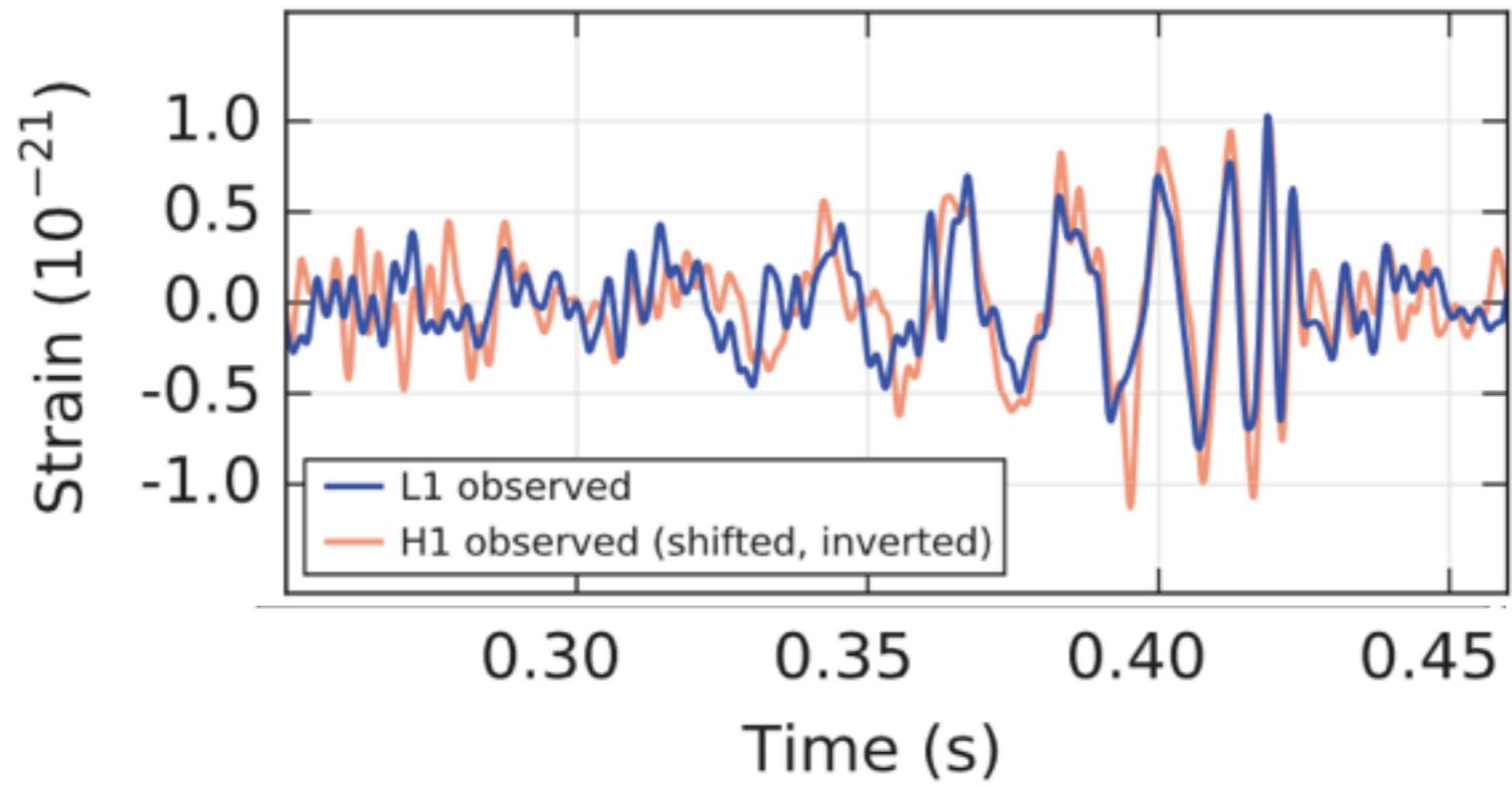
Binary black-hole interactions provide potentially the strongest source of gravitational radiation for detectors currently under development. We present some results from the Binary Black Hole Grand Challenge Alliance three-dimensional Cauchy evolution module. These constitute essential steps towards modeling such interactions and predicting gravitational radiation waveforms. We report on single black-hole evolutions and the first successful demonstration of a black hole moving freely through a three-dimensional computational grid via a Cauchy evolution: a hole moving near  $6M$  at  $0.1c$  during a total evolution of duration near  $60M$ . [S0031-9007(98)05652-X]

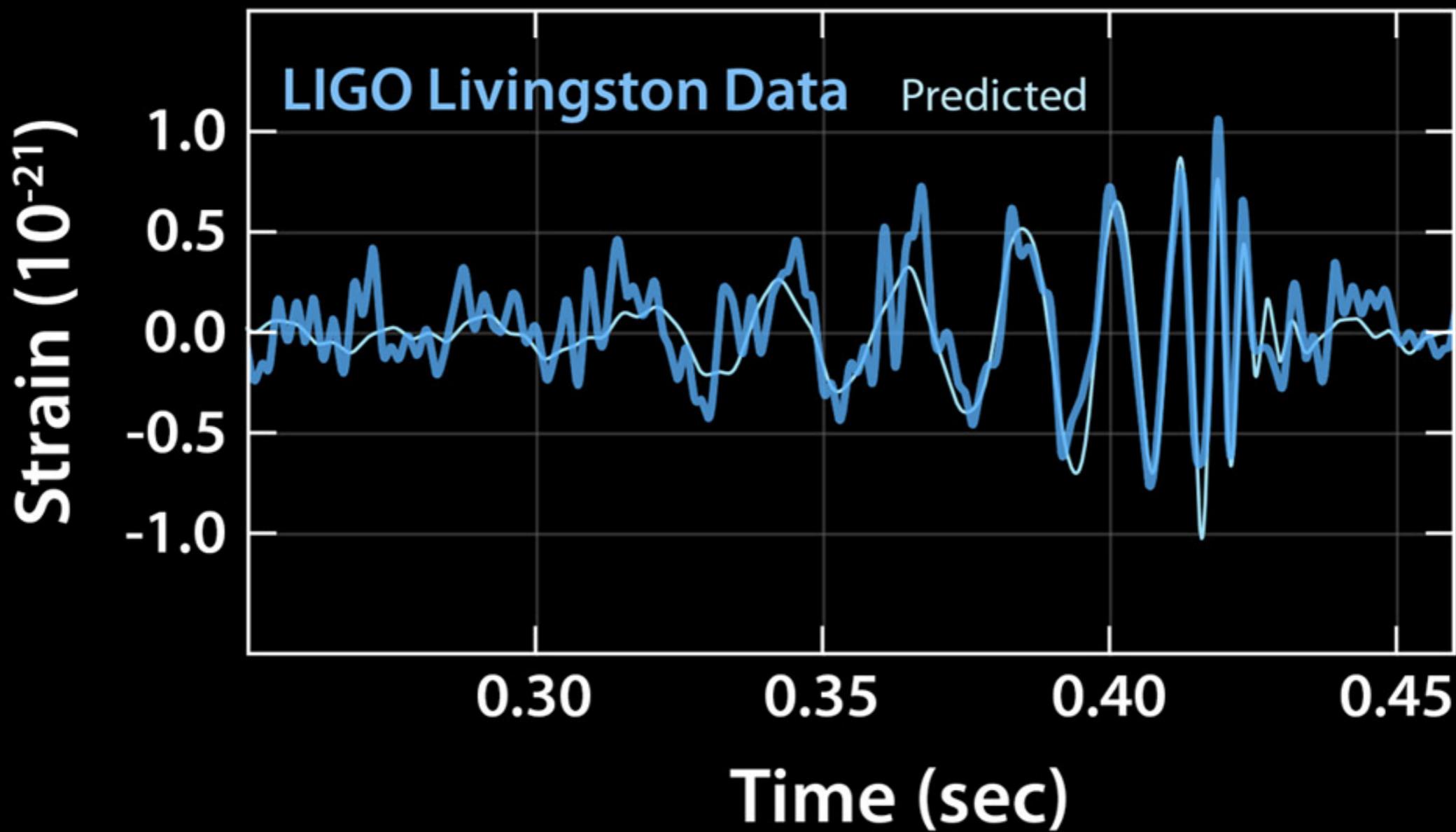
# 2005: Breakthrough!

Pretorius (July): Generalised harmonic formalism

NASA-Goddard and Brownsville-Texas (November):  
moving-puncture method.







Waveform modelling  
and  
black-hole  
measurements

Masses:  $m_1, m_2$

Spins:  $\mathbf{S}_1, \mathbf{S}_2$

(8 parameters)

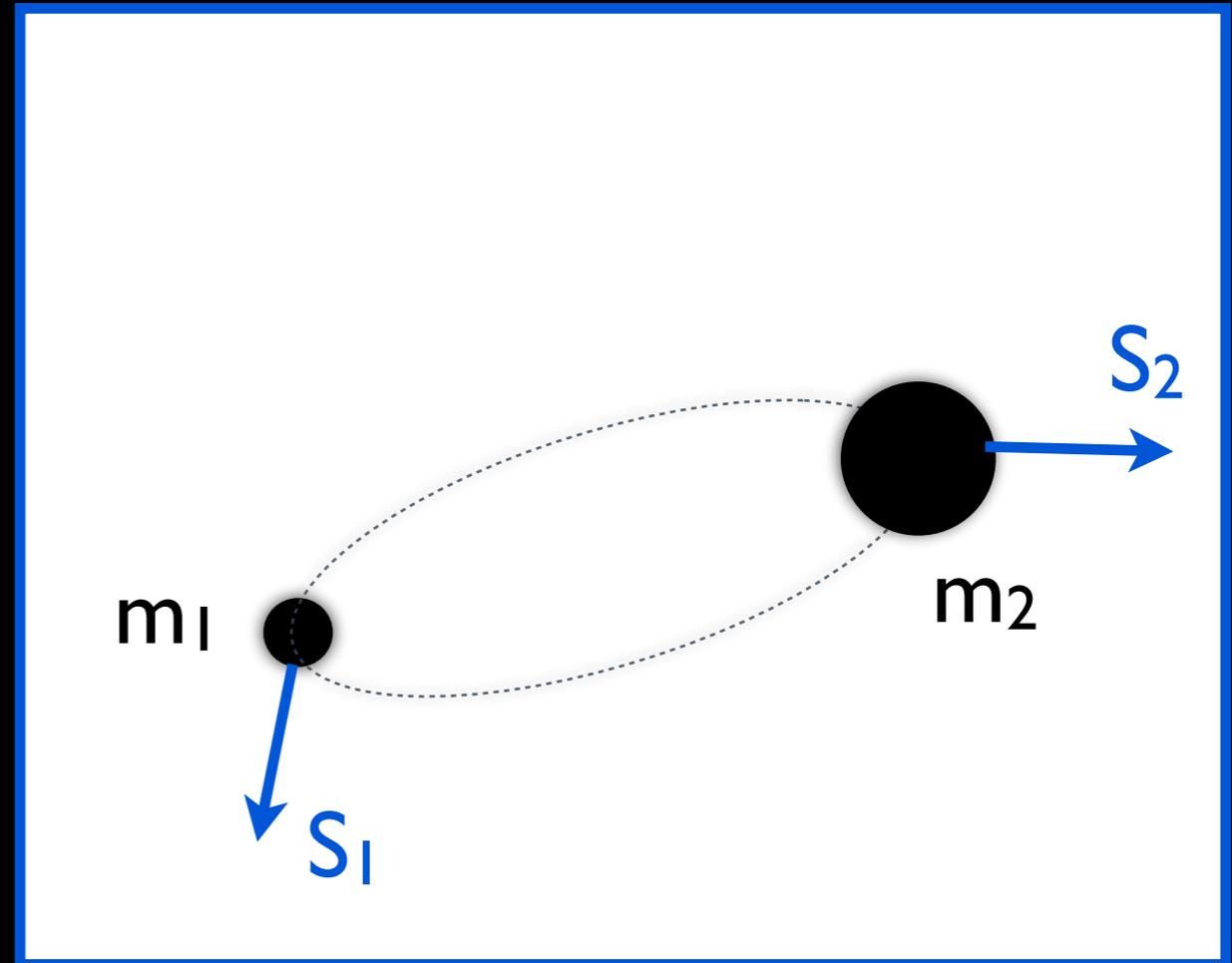
useful combinations:

$$M = m_1 + m_2$$

$$q = m_2 / m_1$$

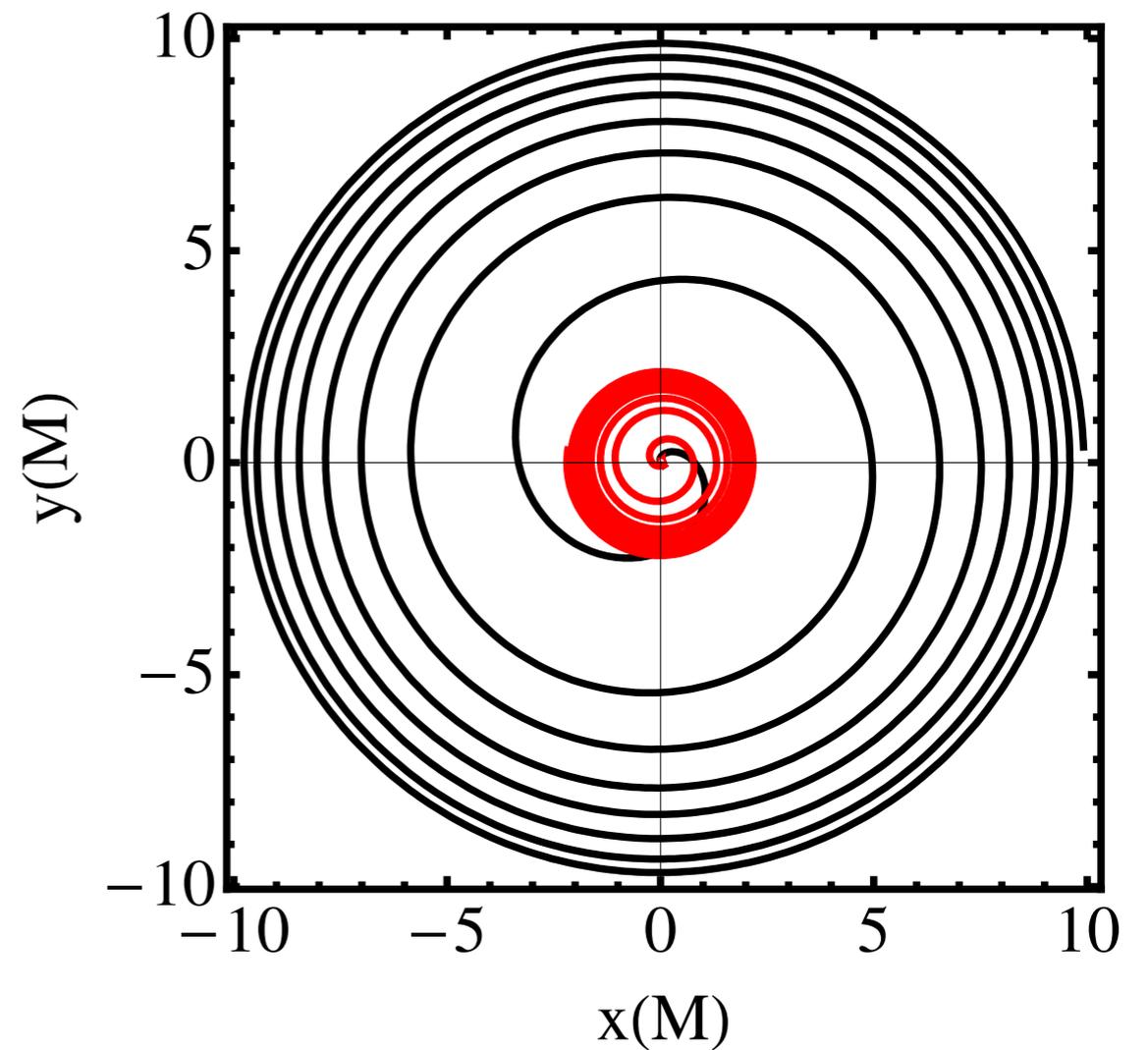
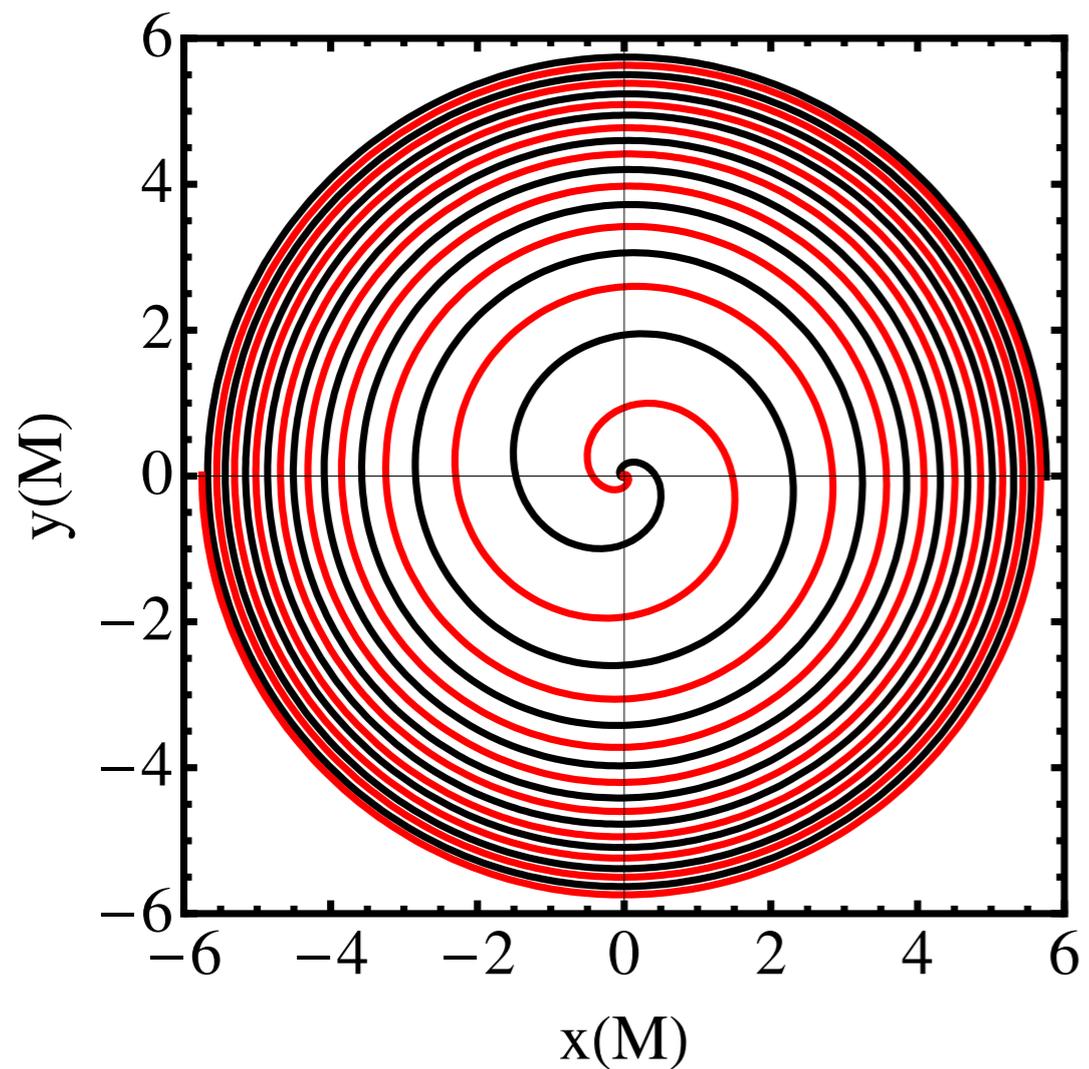
$$\eta = m_1 m_2 / M^2$$

$$\chi = S/m^2$$



**Plus:** distance, sky location,  
orientation, polarisation

# Nonspinning black holes



Amplitude:



Optimally oriented  
(face on)

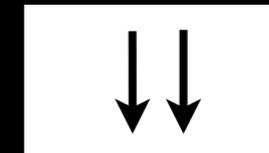
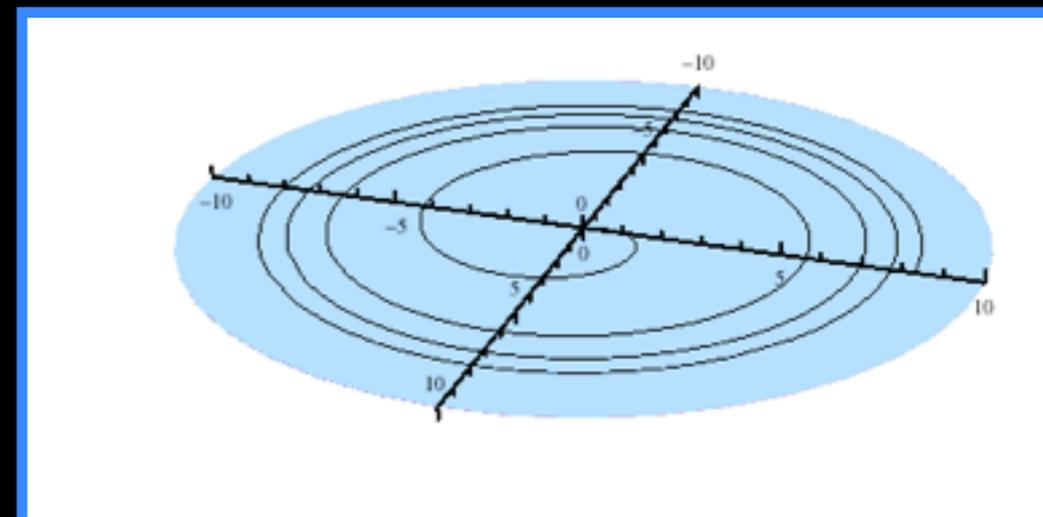
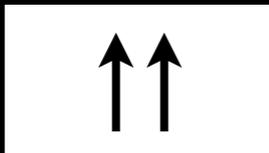
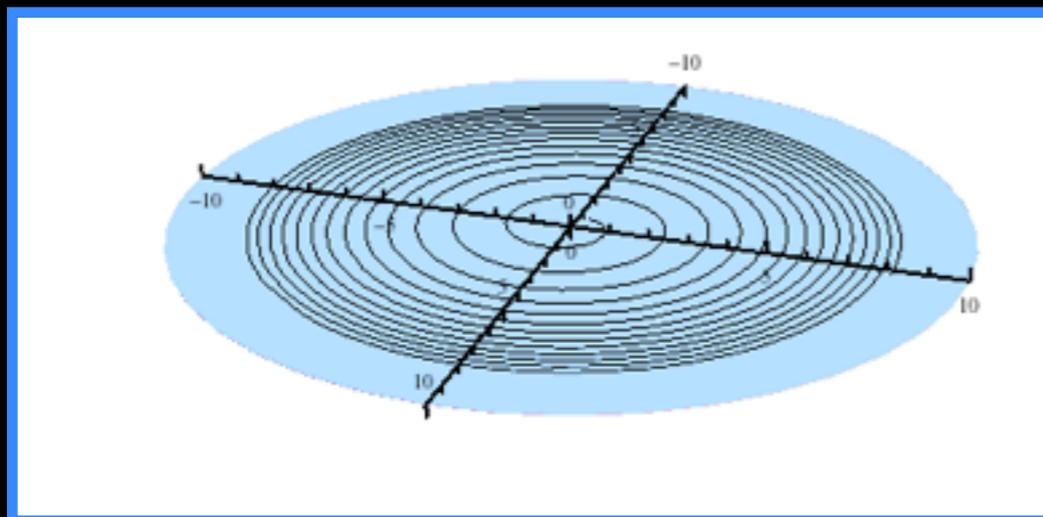
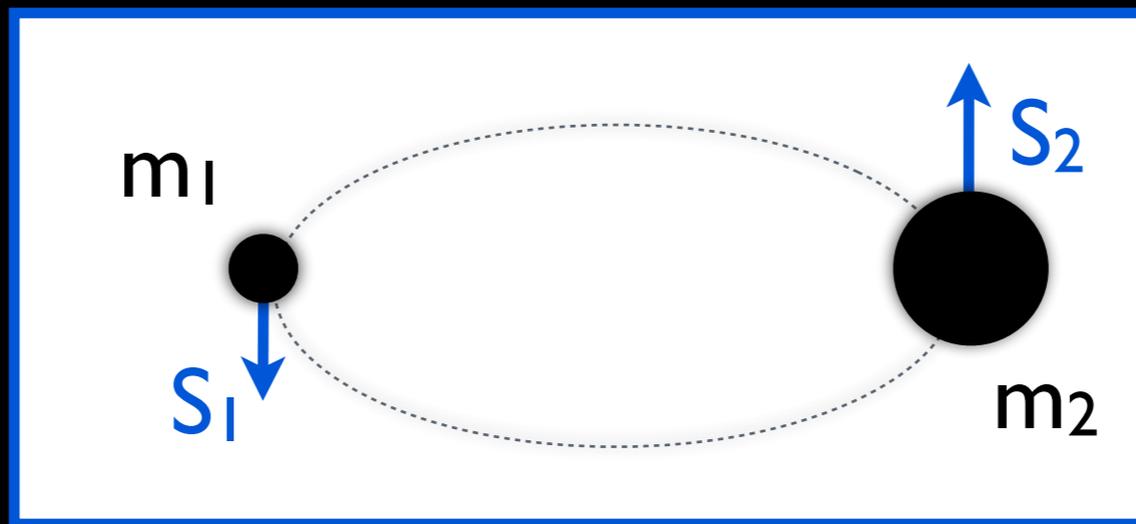


Edge on

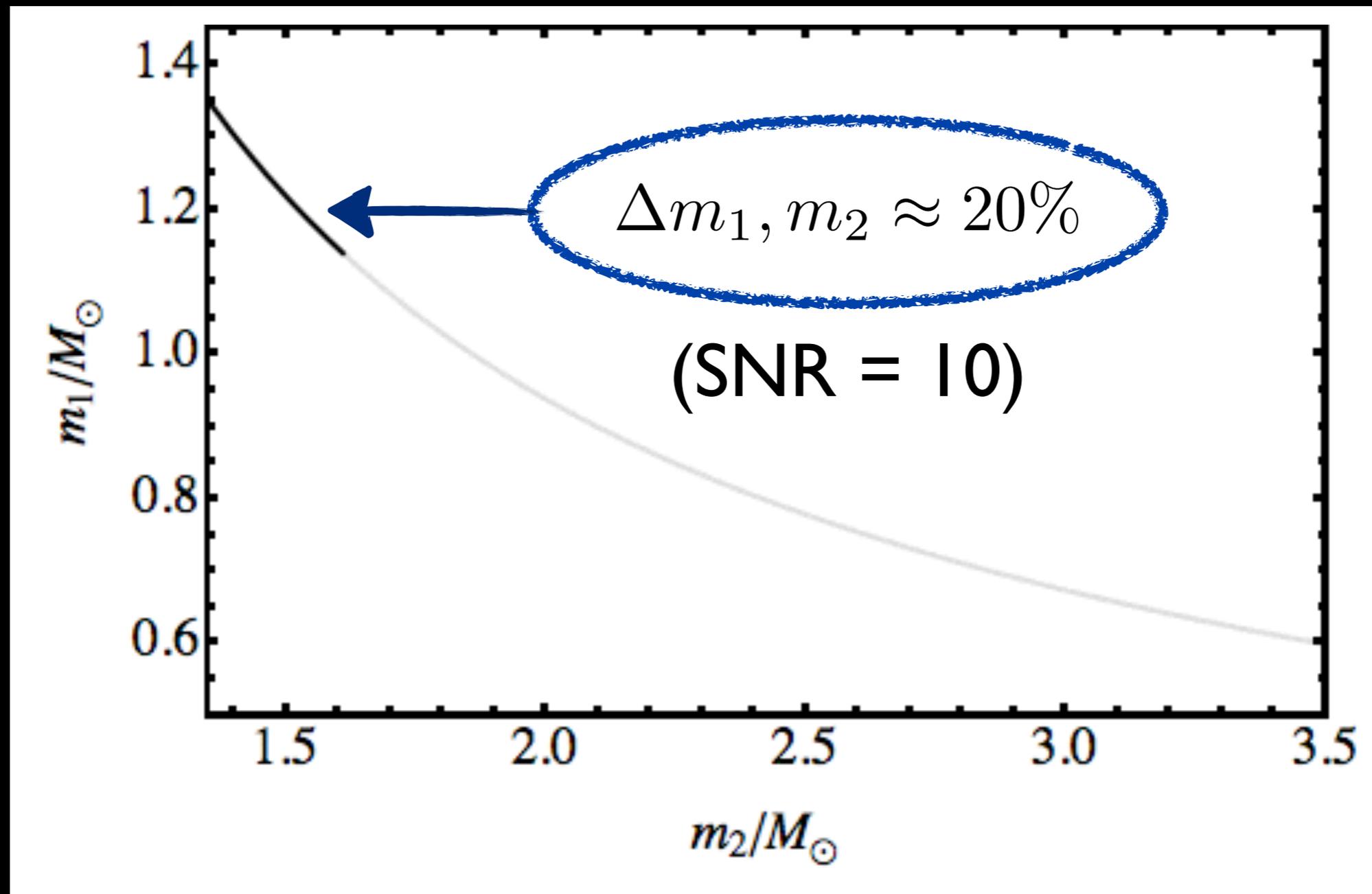
Signal shape is independent of orientation

Key information is in the phasing

# Aligned spins

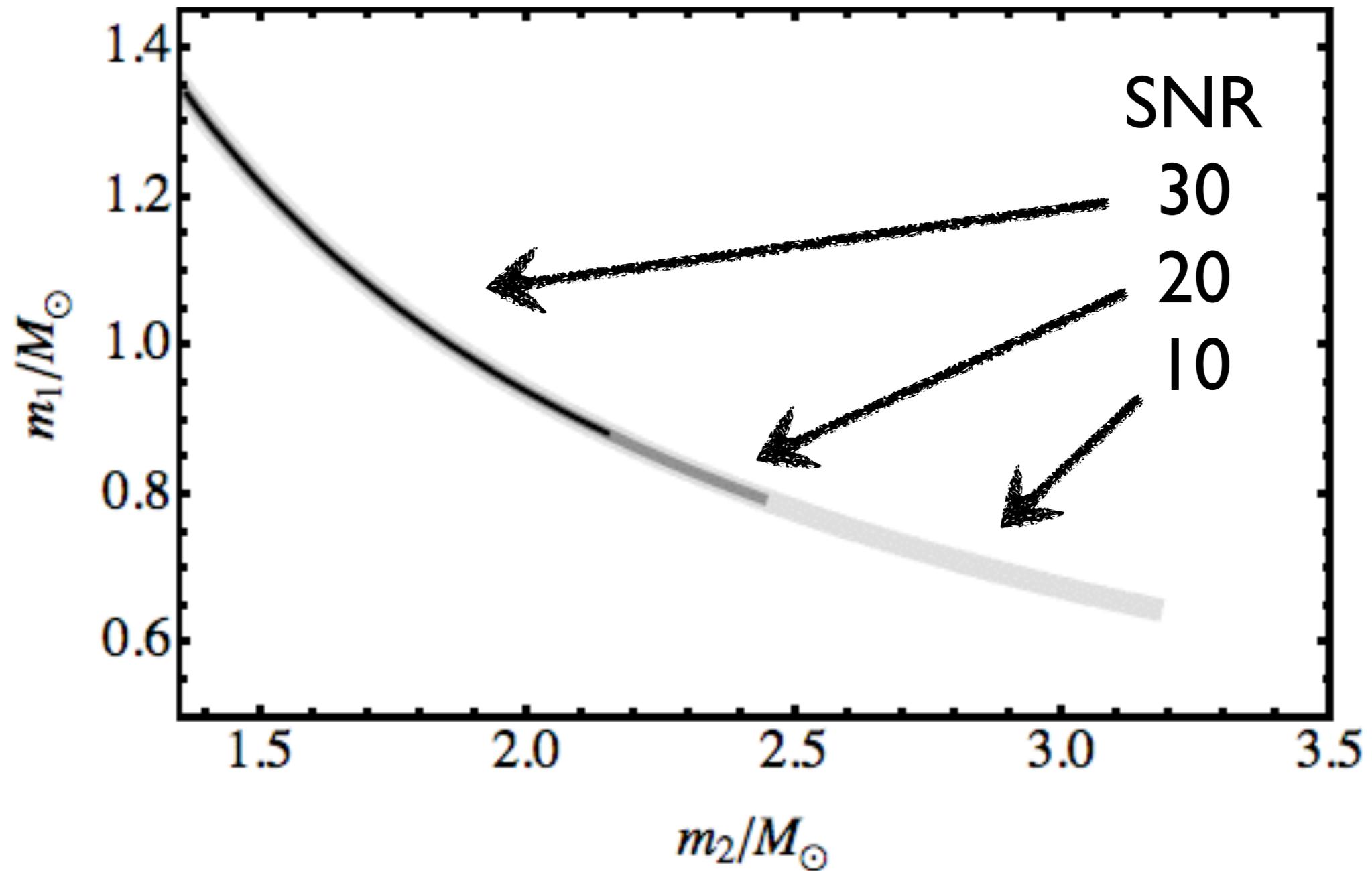


# Mass measurements (non-spinning)

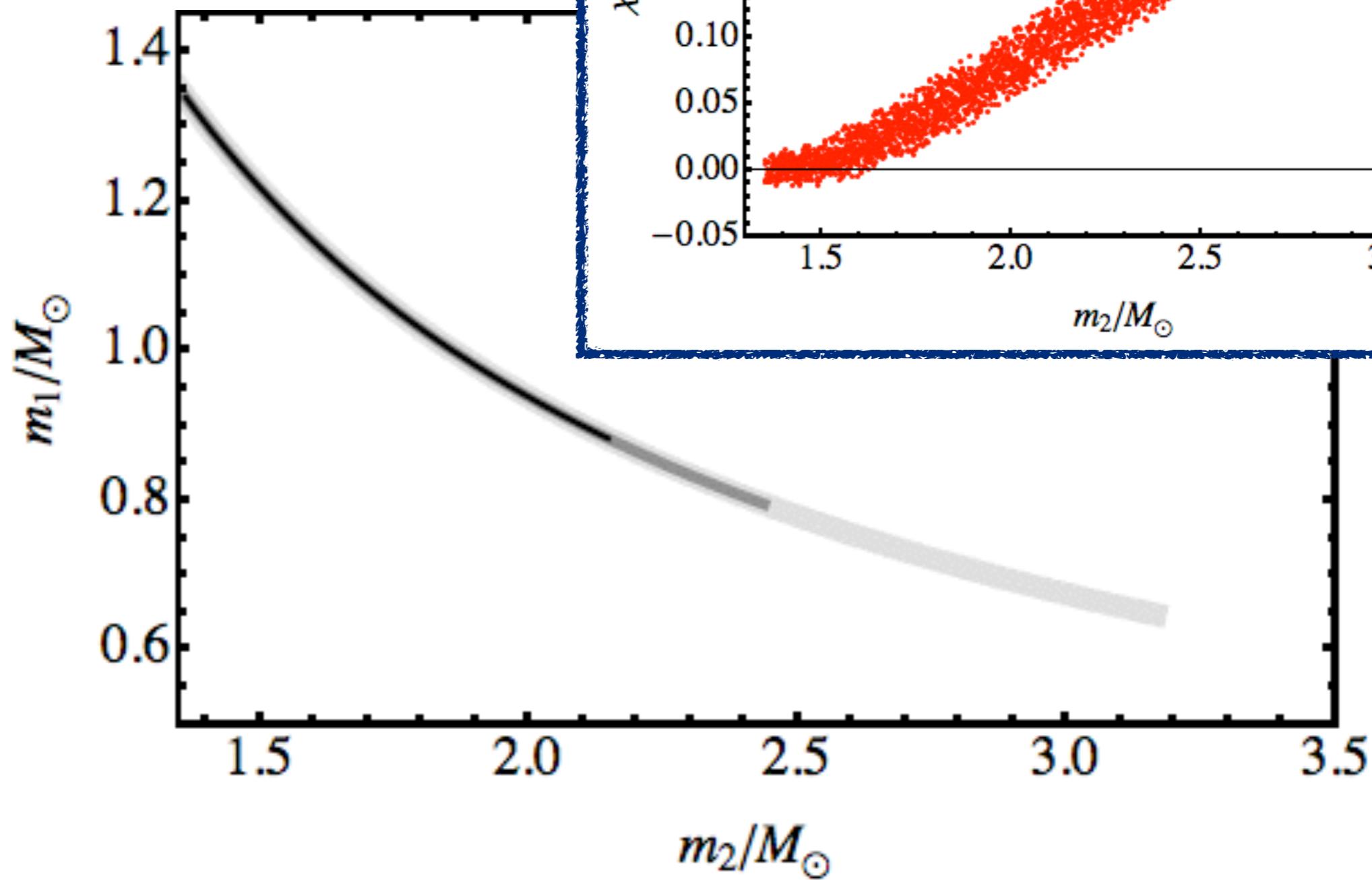


[Hannam, et. al (2013)]

# Aligned spins



[Hannam, et. al (2013)]



[Hannam, et. al (2013)]

# What is “ $\chi$ ”?

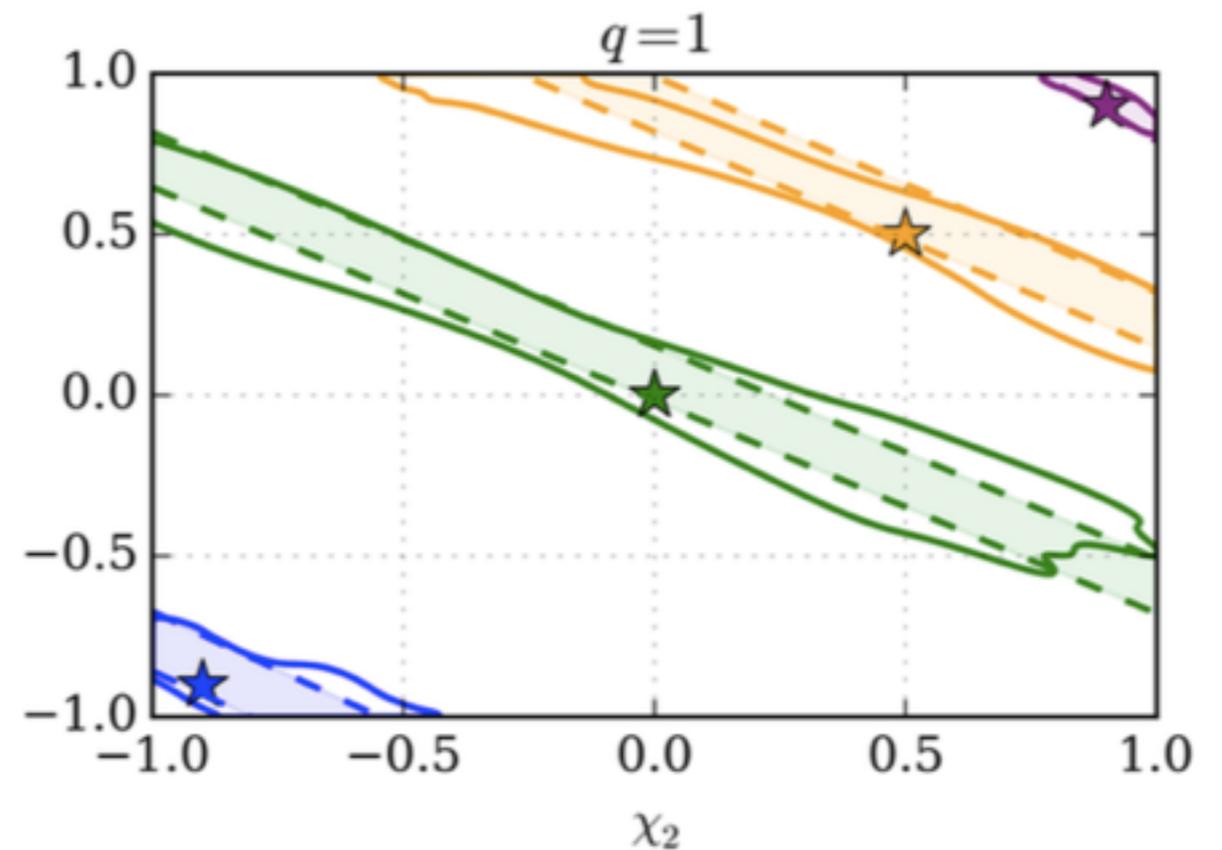
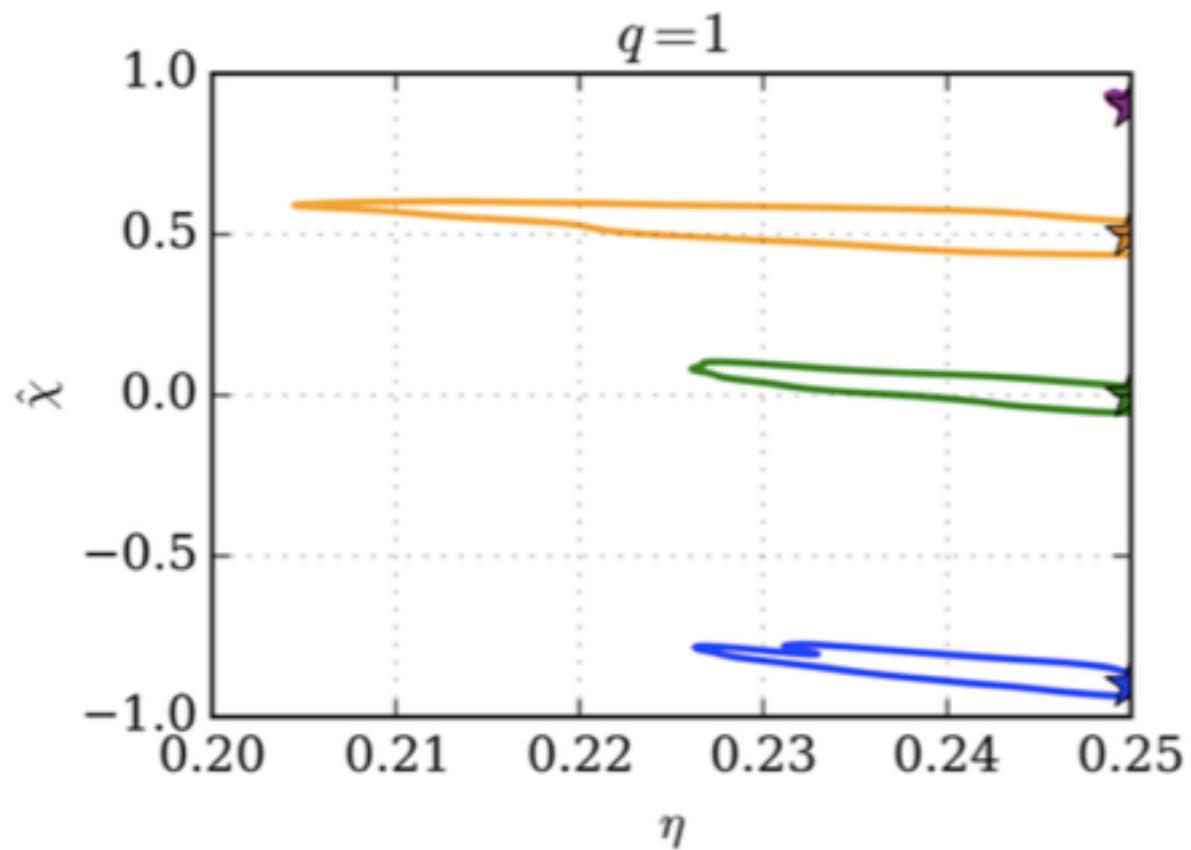
$\chi$  is a weighted sum of the two spins

$\chi$  is the dominant spin effect on the phasing

The individual spins have only a weak effect

$$\chi = \frac{m_1\chi_1 + m_2\chi_2}{M} - \frac{76\eta}{226}(\chi_1 + \chi_2)$$

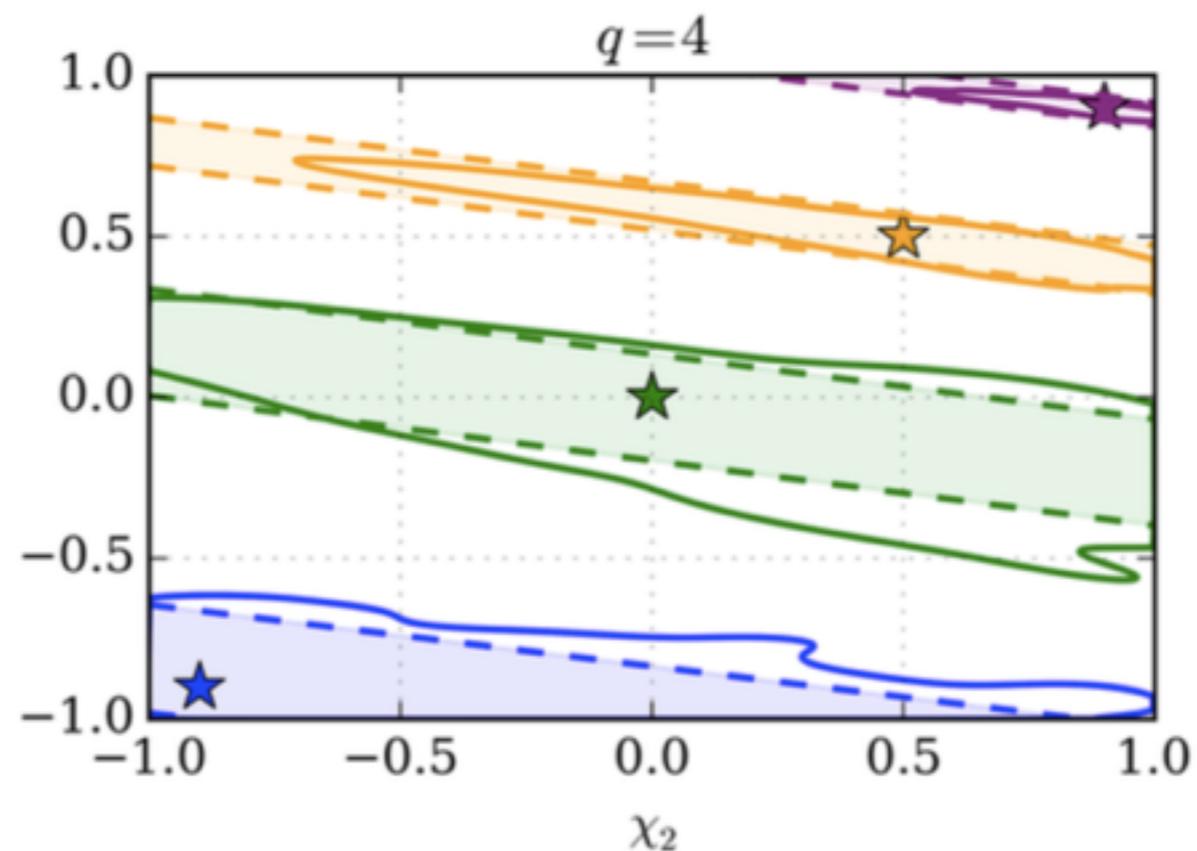
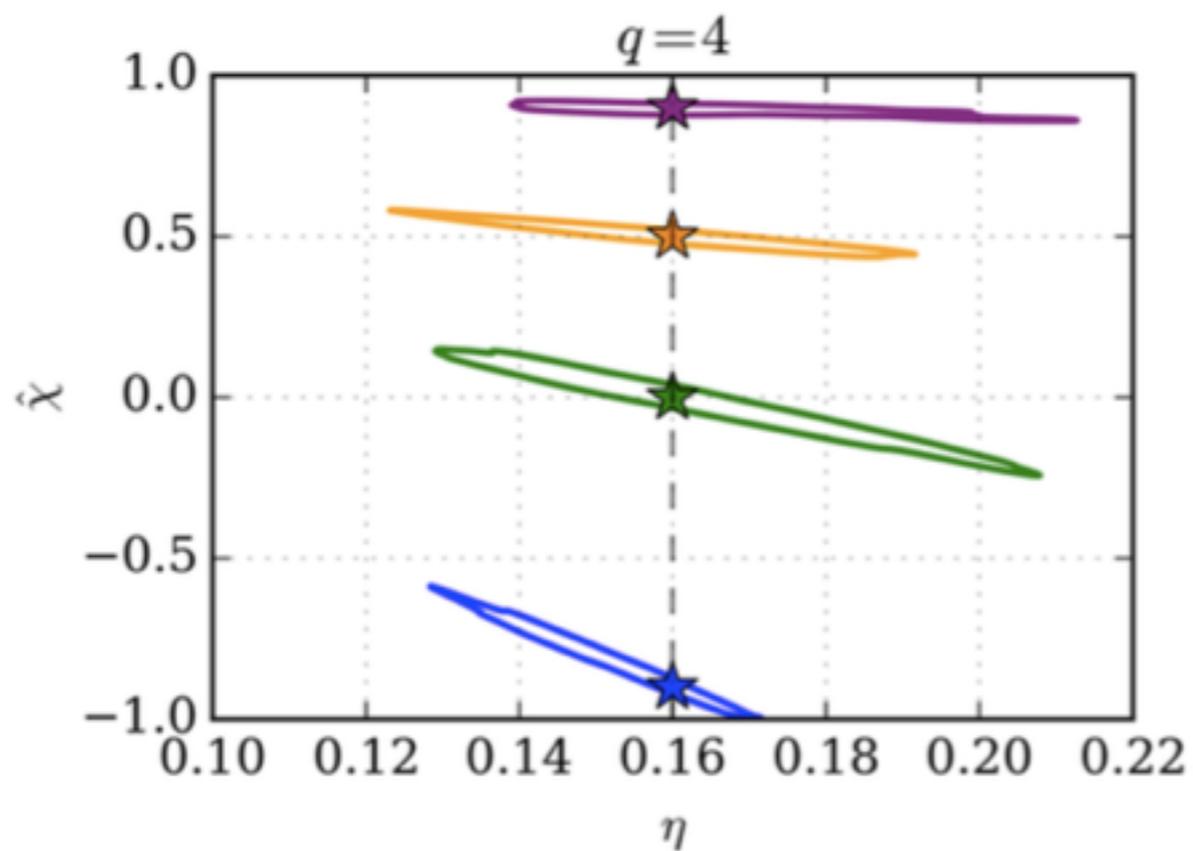
(50-solar-mass, equal spins)



[Puerrer, Hannam, Ohme (2016)]

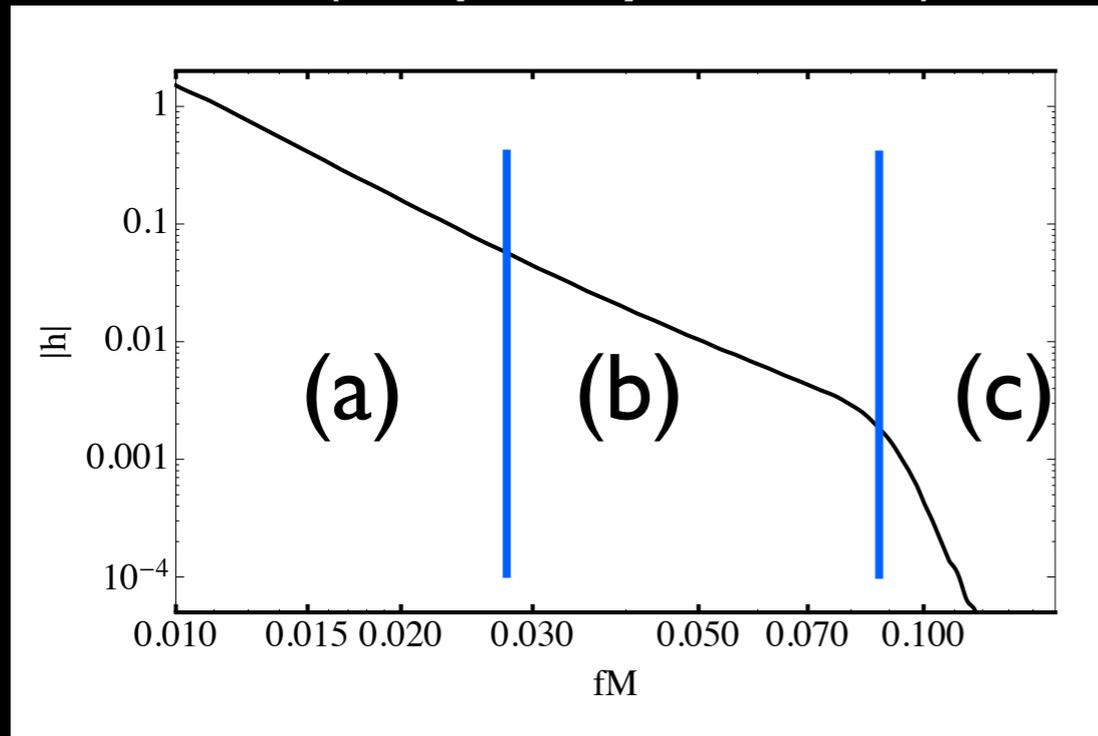
$$\chi = \frac{m_1\chi_1 + m_2\chi_2}{M} - \frac{76\eta}{226}(\chi_1 + \chi_2)$$

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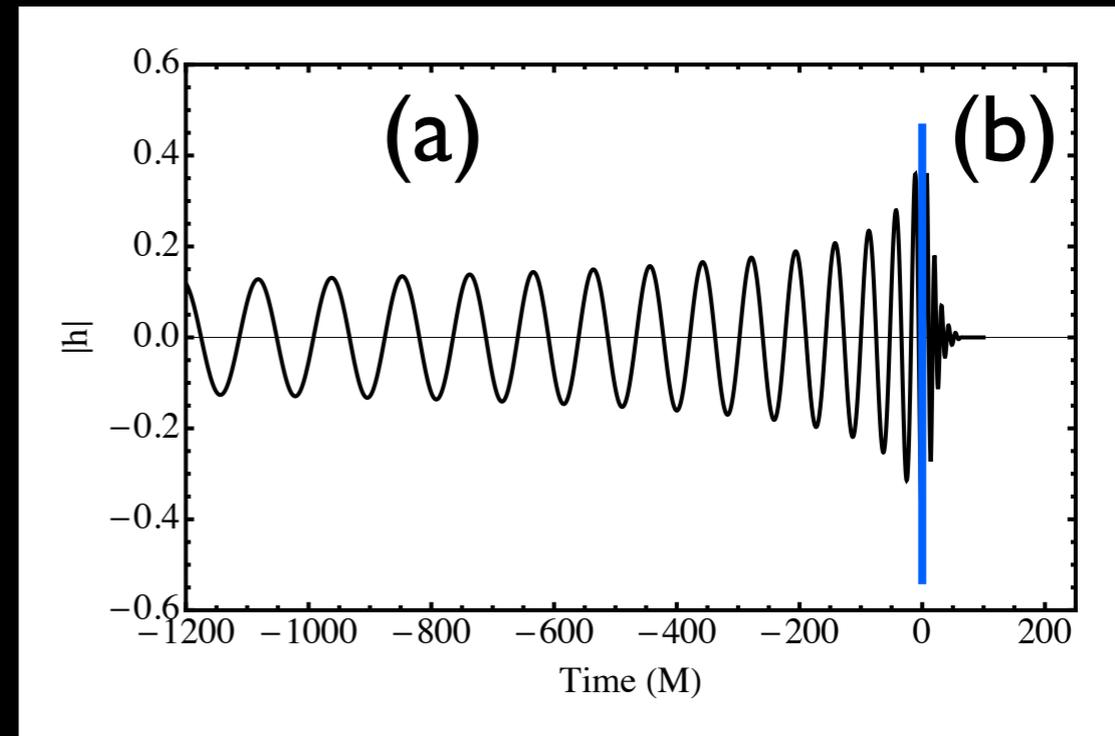
[Puerrer, Hannam, Ohme (2016)]

# Phenom (frequency domain)

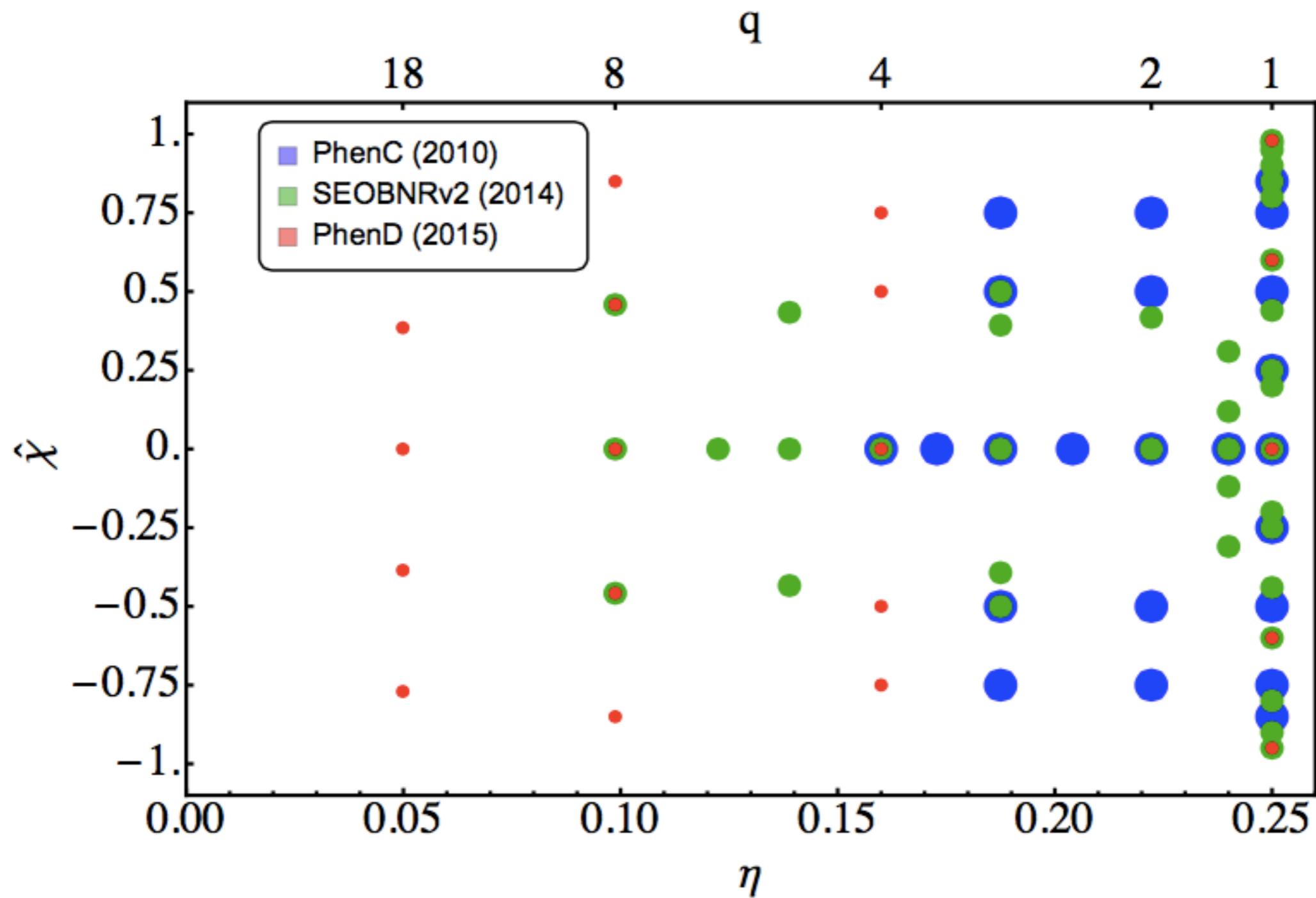


- (a) PN-based ansatz
- (b) phenomenological fit (based on NR behaviour)
- (c) FFT of ringdown waveform (Lorentzian)
- Analytic: fast

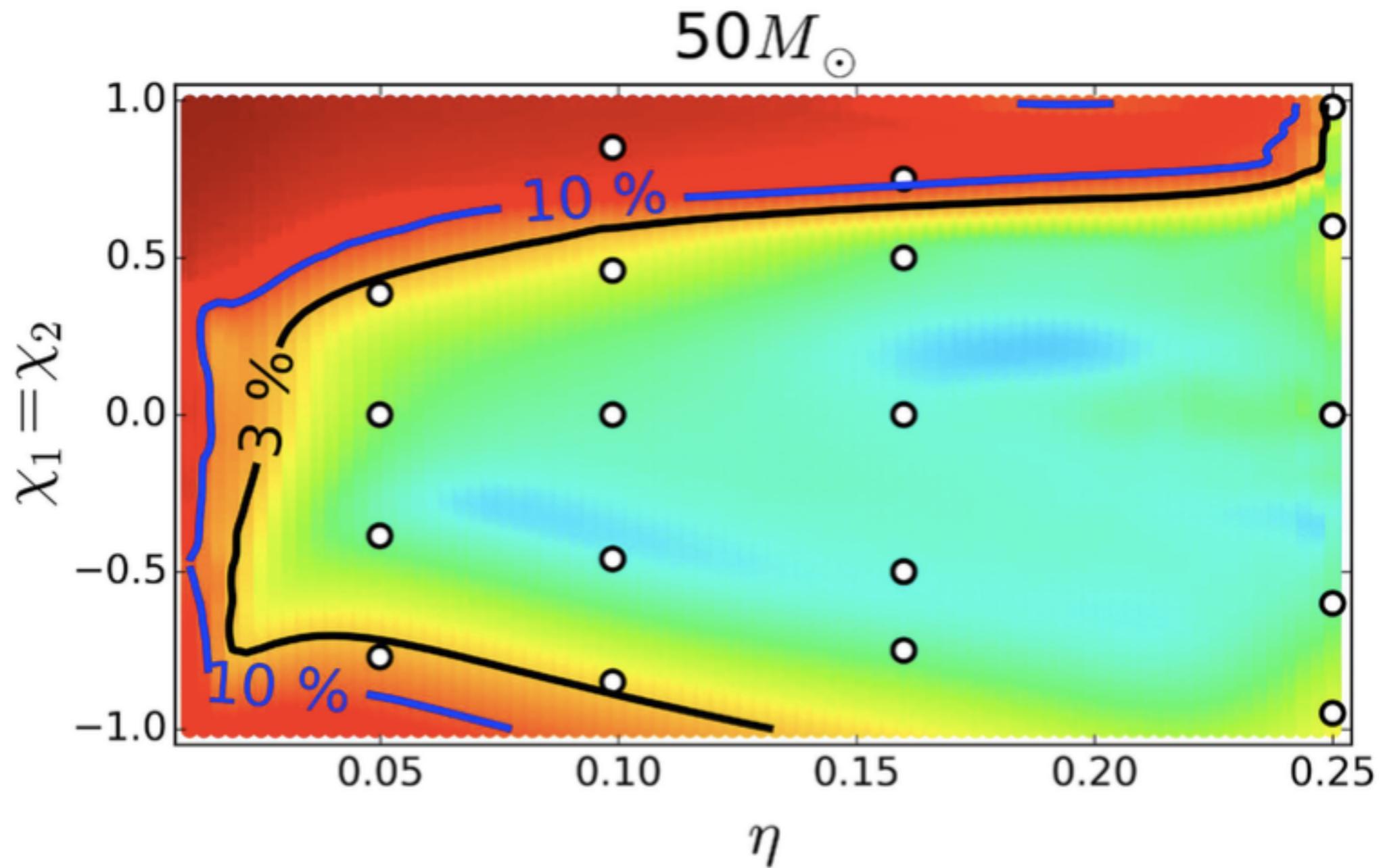
# EOB-NR (time domain)



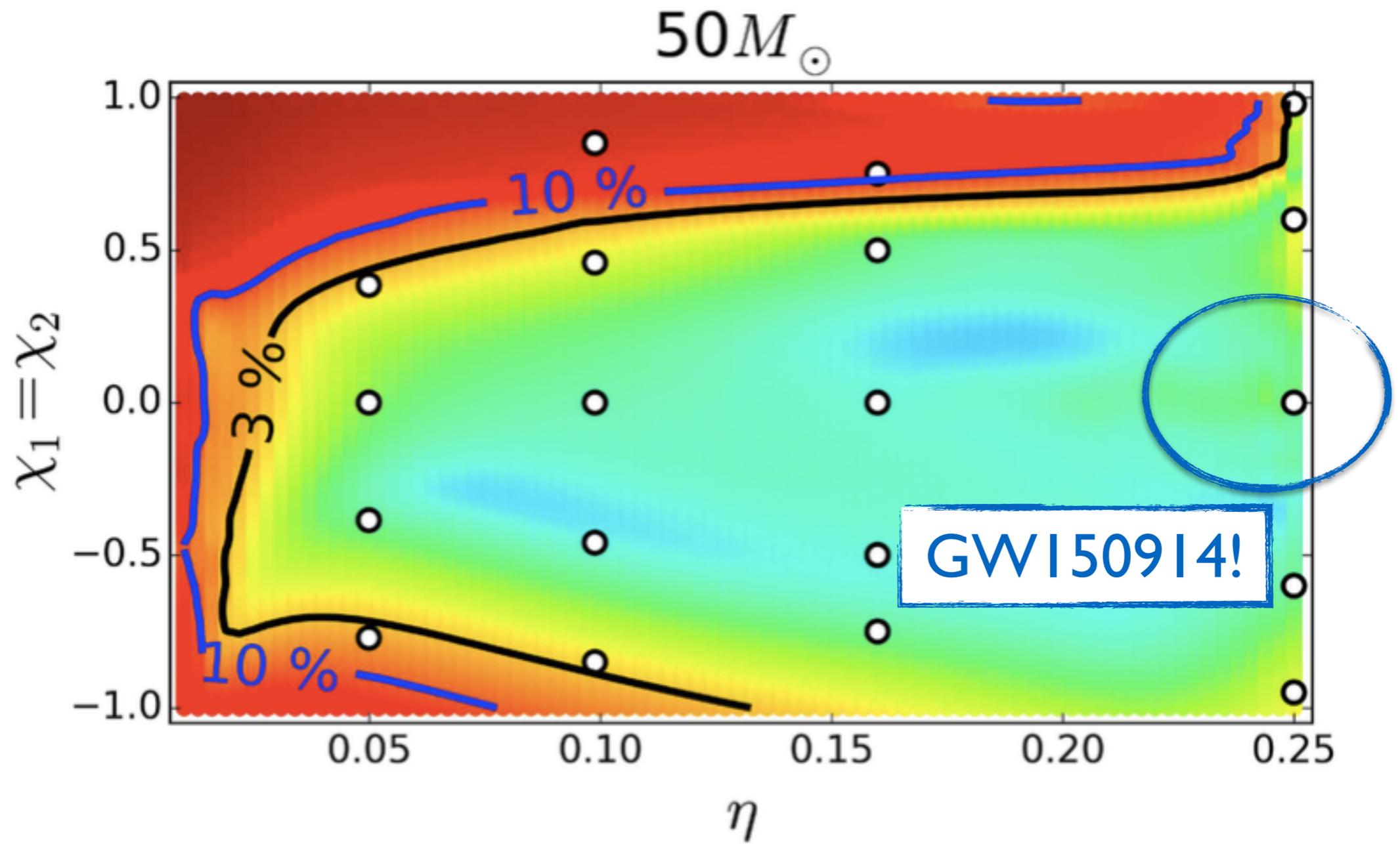
- (a) EOB + terms tuned to NR waveforms
- (b) Smooth transition to ringdown
- Includes both spins
- Numerically solve ODEs: slow
- Speed-up: Reduced-order models



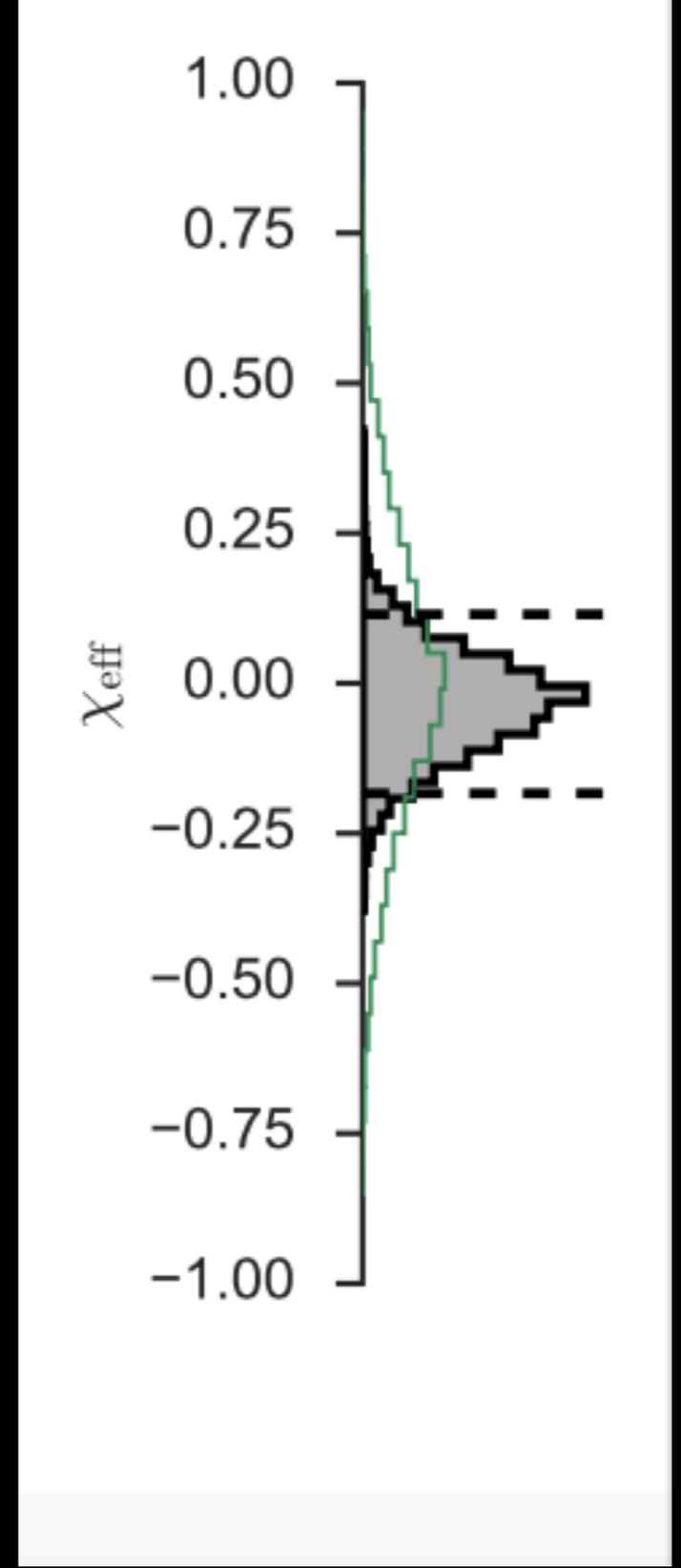
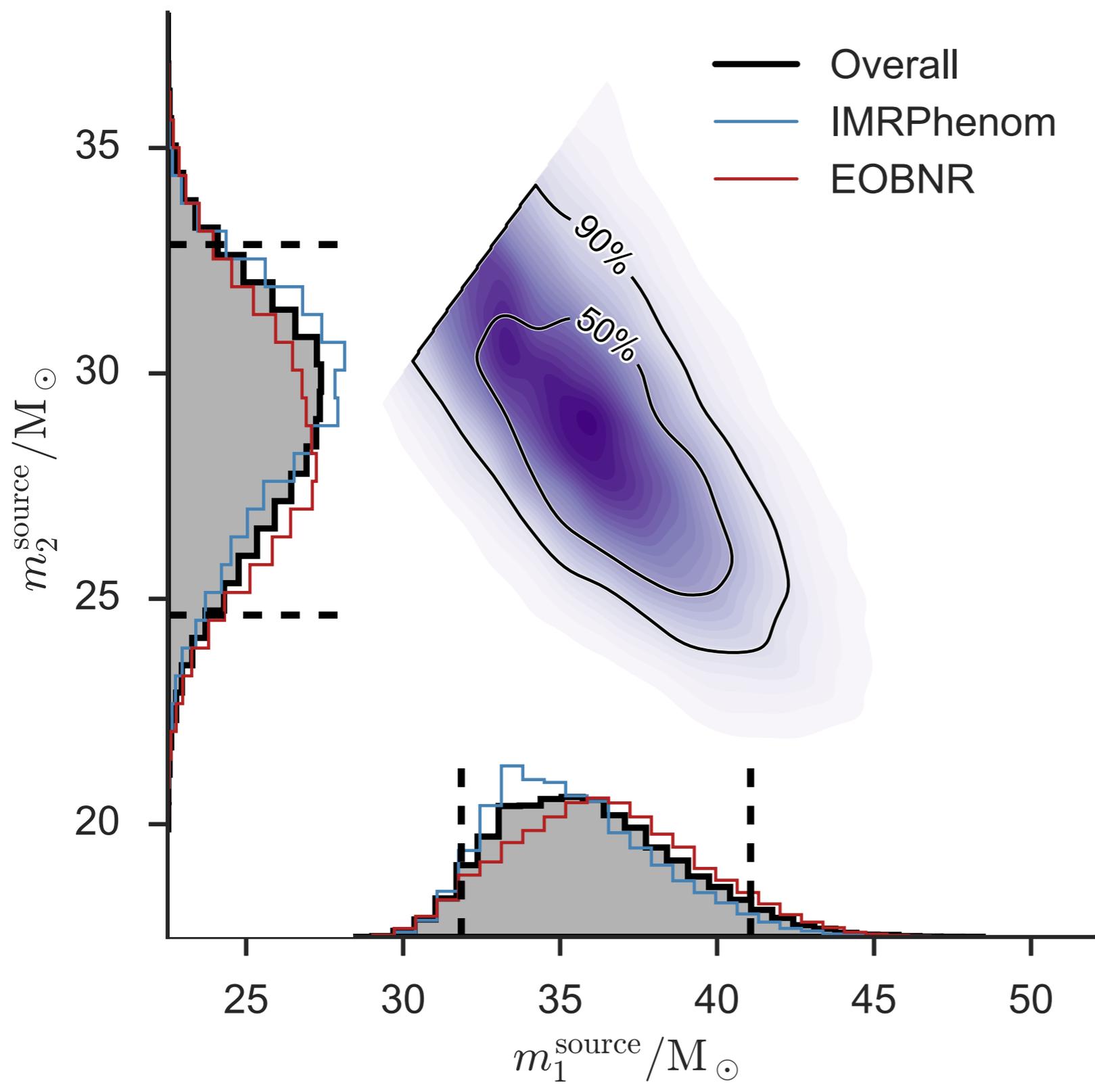
[Khan, *et. al* (2016)]



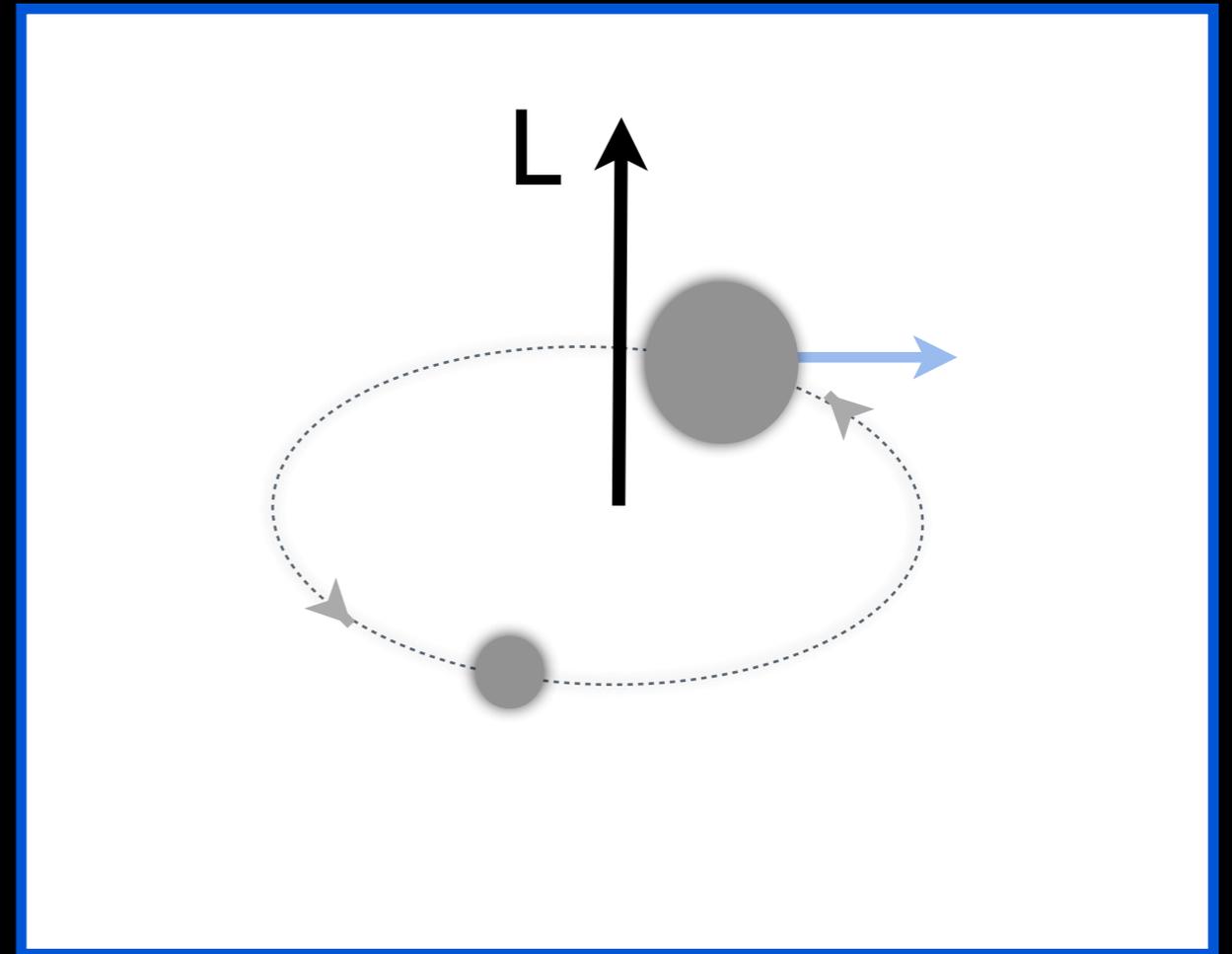
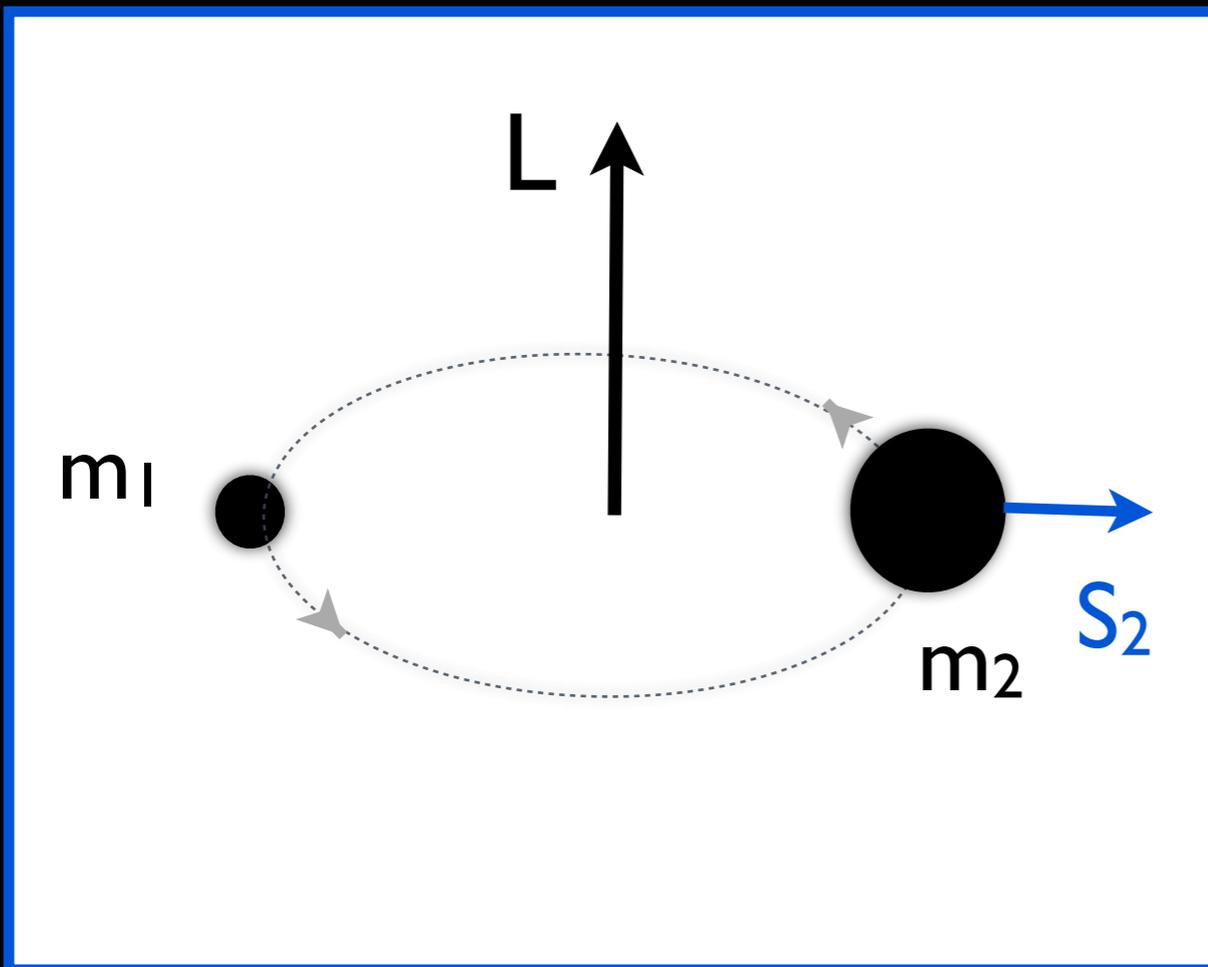
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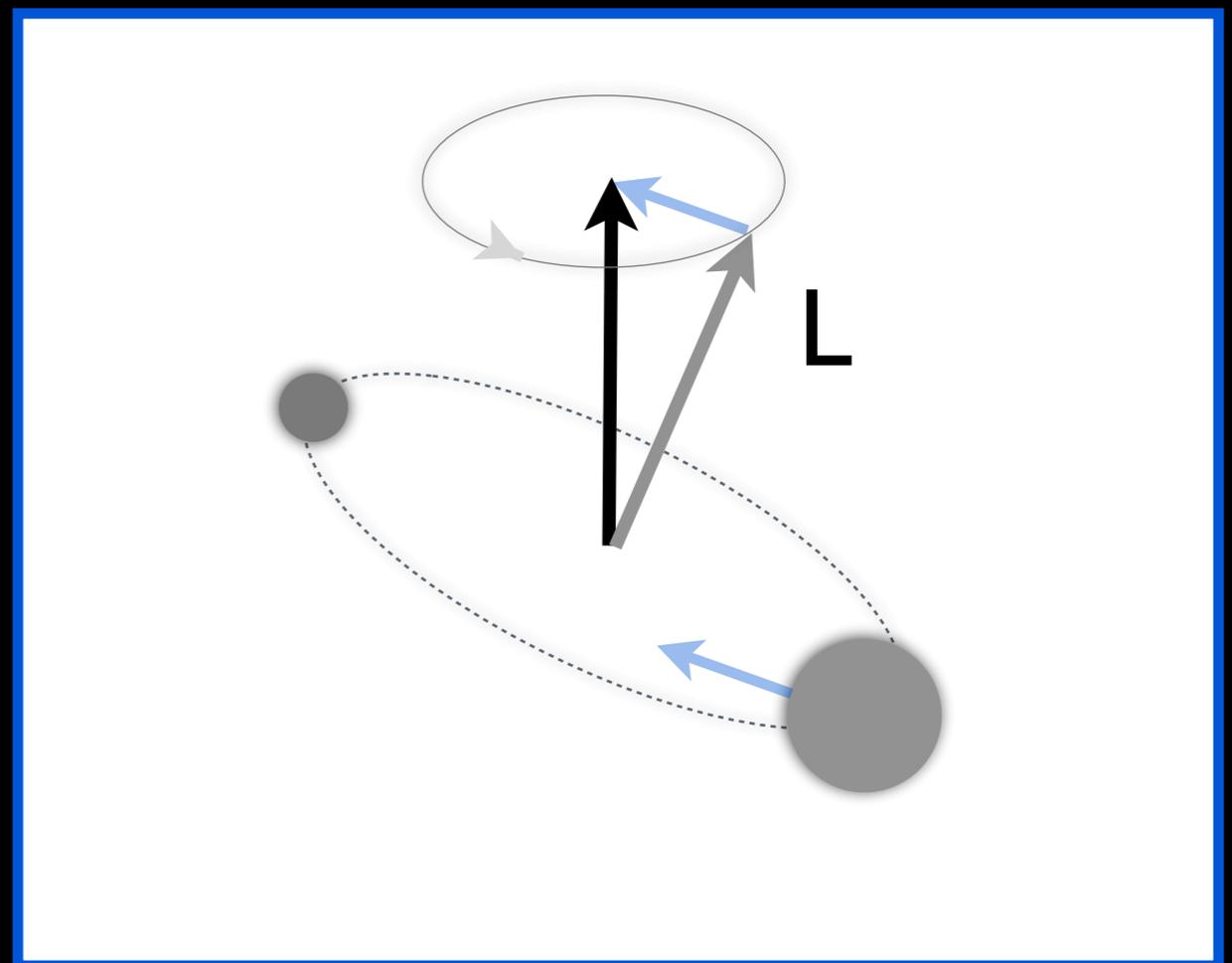
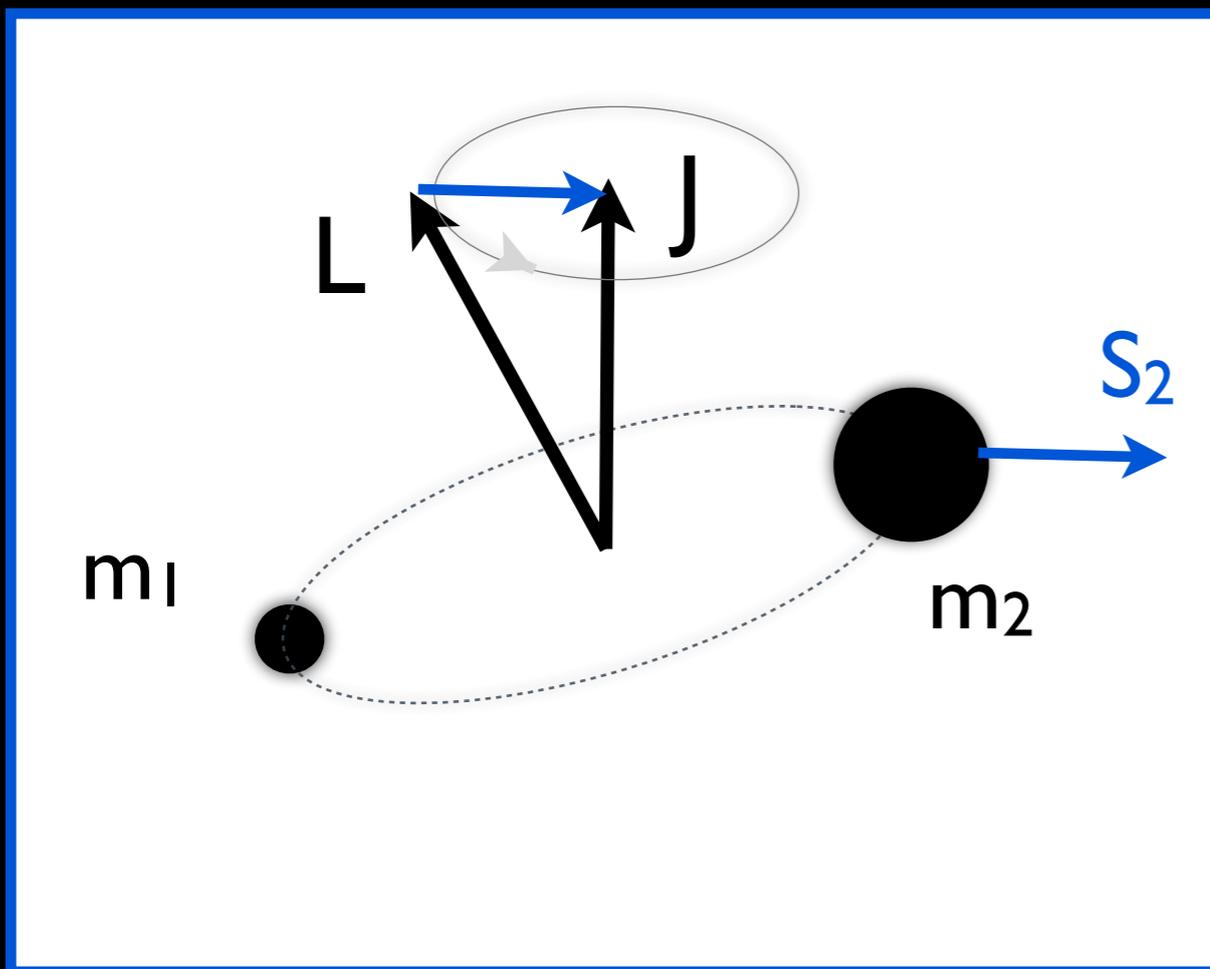


# Orbital precession



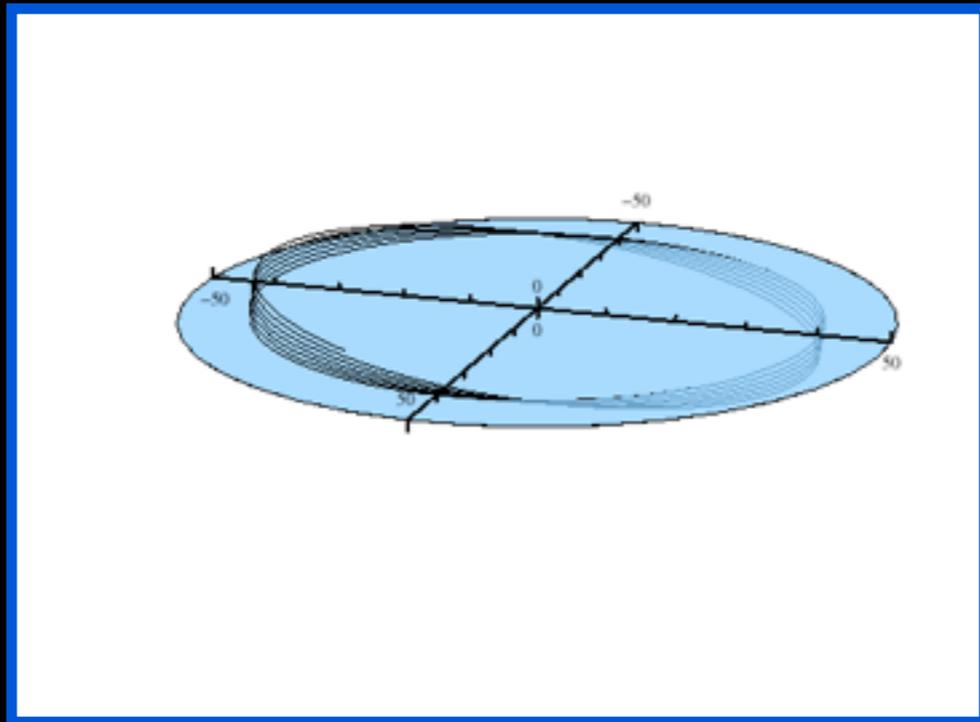
Newtonian gravity:  
 $L, S_1, S_2$  remain fixed

# Orbital precession

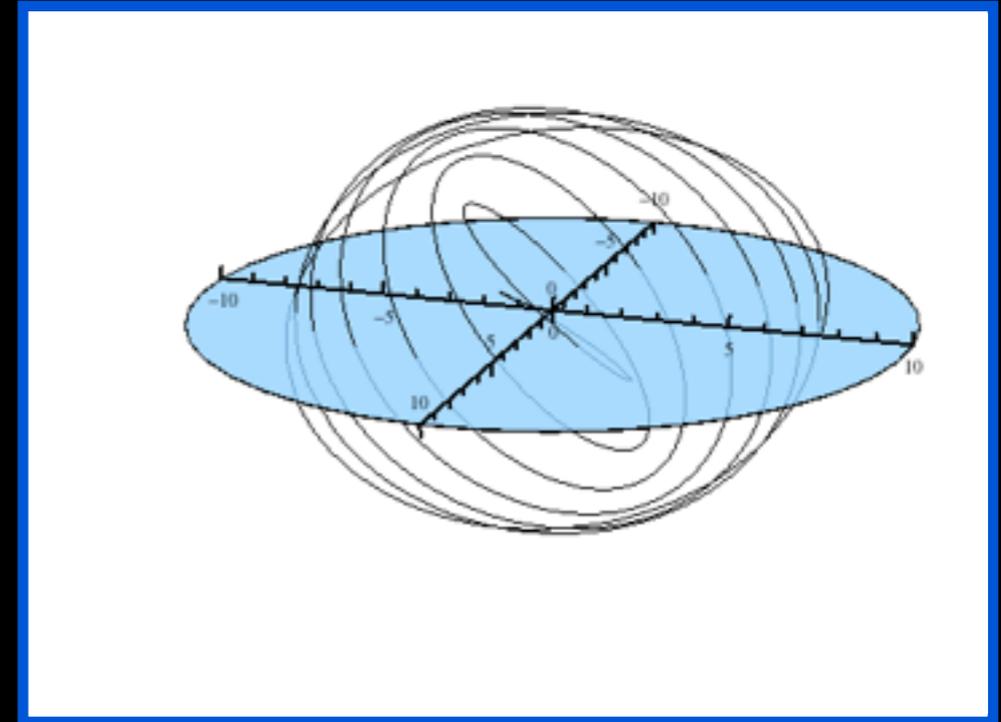


General relativity  
( $L, S_1, S_2$ ) precess around  $J$

# Precessional dynamics



Large separation

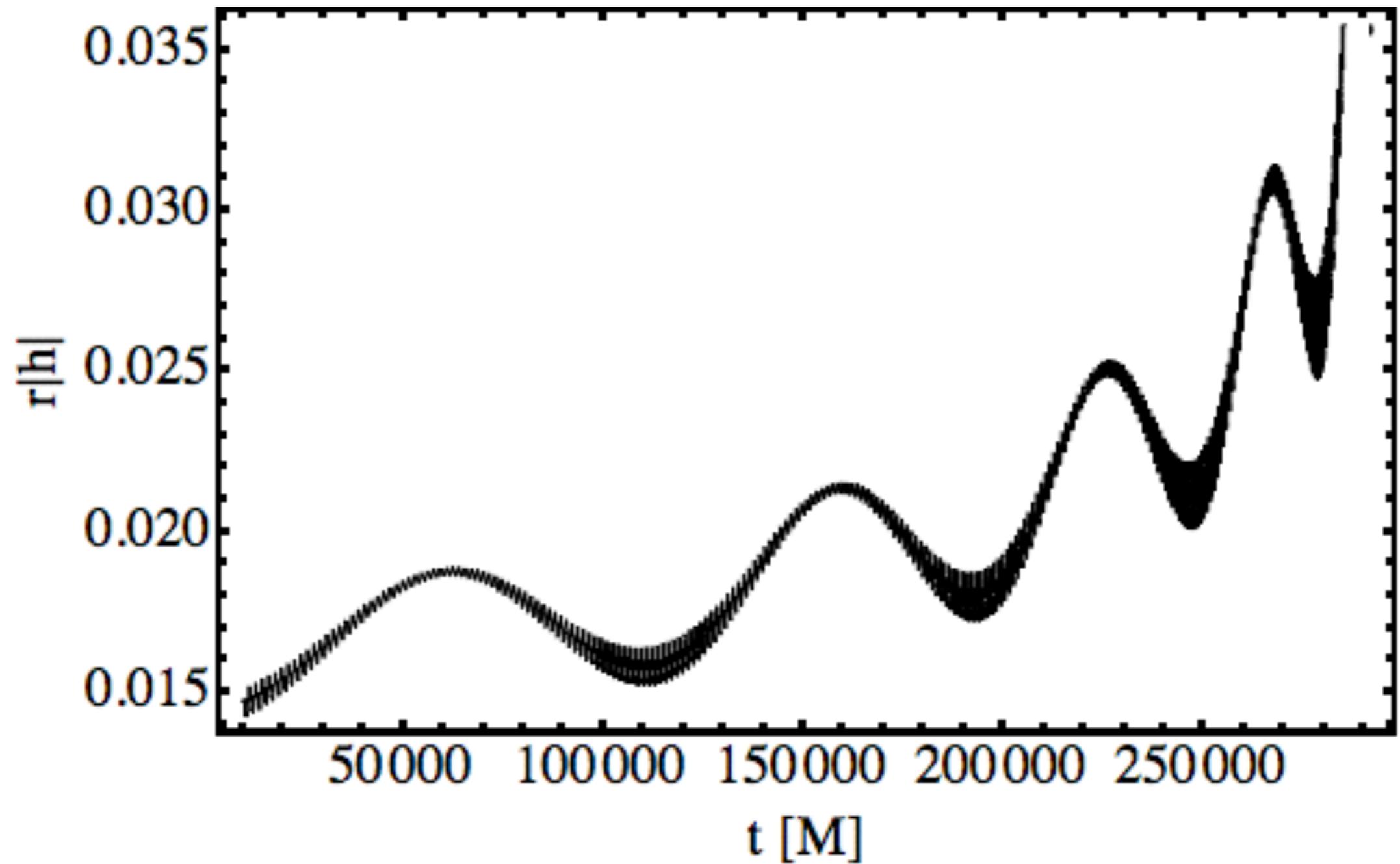


Merger

# Aside: modelling precession

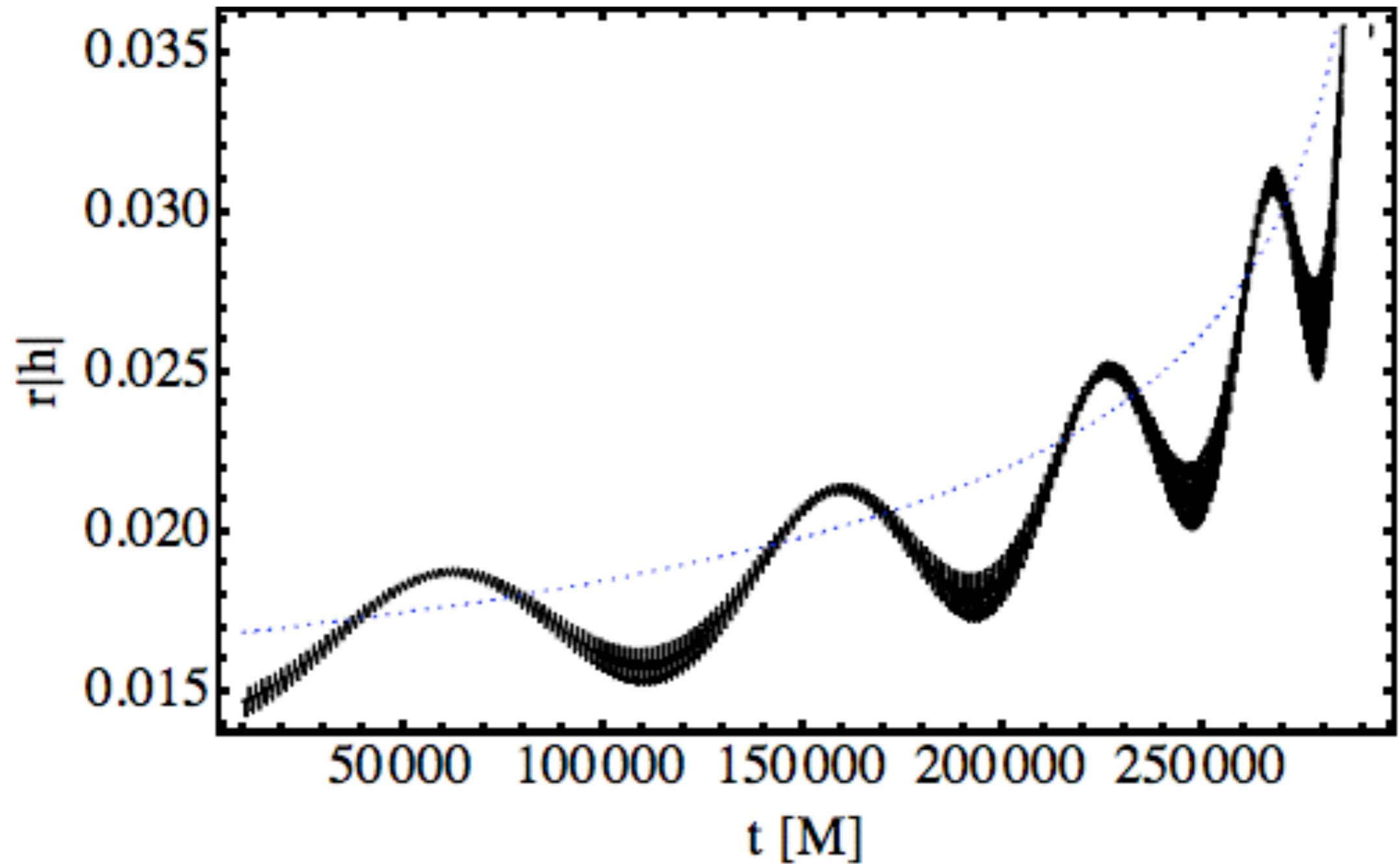
Precessing waveform =  
(non-precessing waveform)  
x (time dependent rotation)

( $q=3$  precessing binary, inclination 2.8 rad)



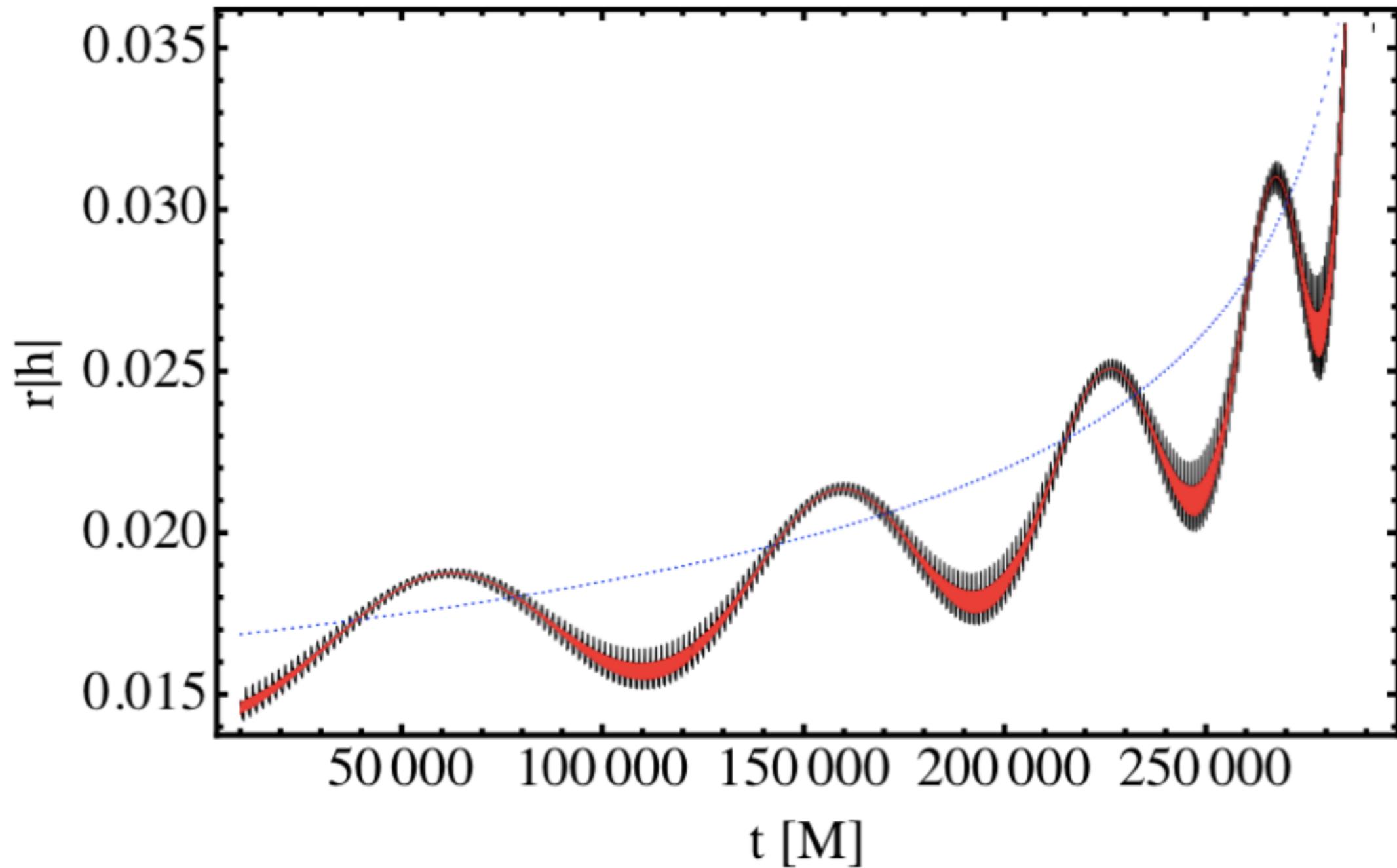
[Schmidt, Hannam, Husa (2012)]

( $q=3$  precessing binary, inclination 2.8 rad)



[Schmidt, Hannam, Husa (2012)]

( $q=3$  precessing binary, inclination 2.8 rad)



[Schmidt, Hannam, Husa (2012)]

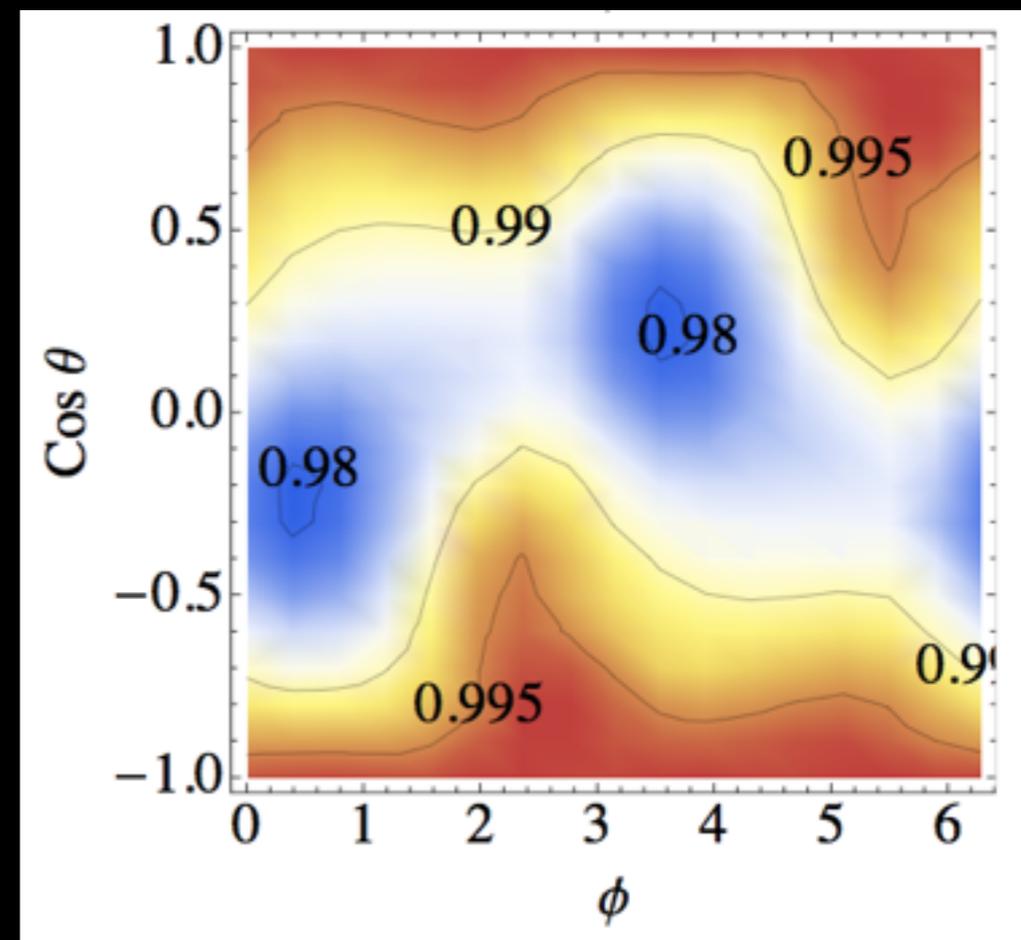
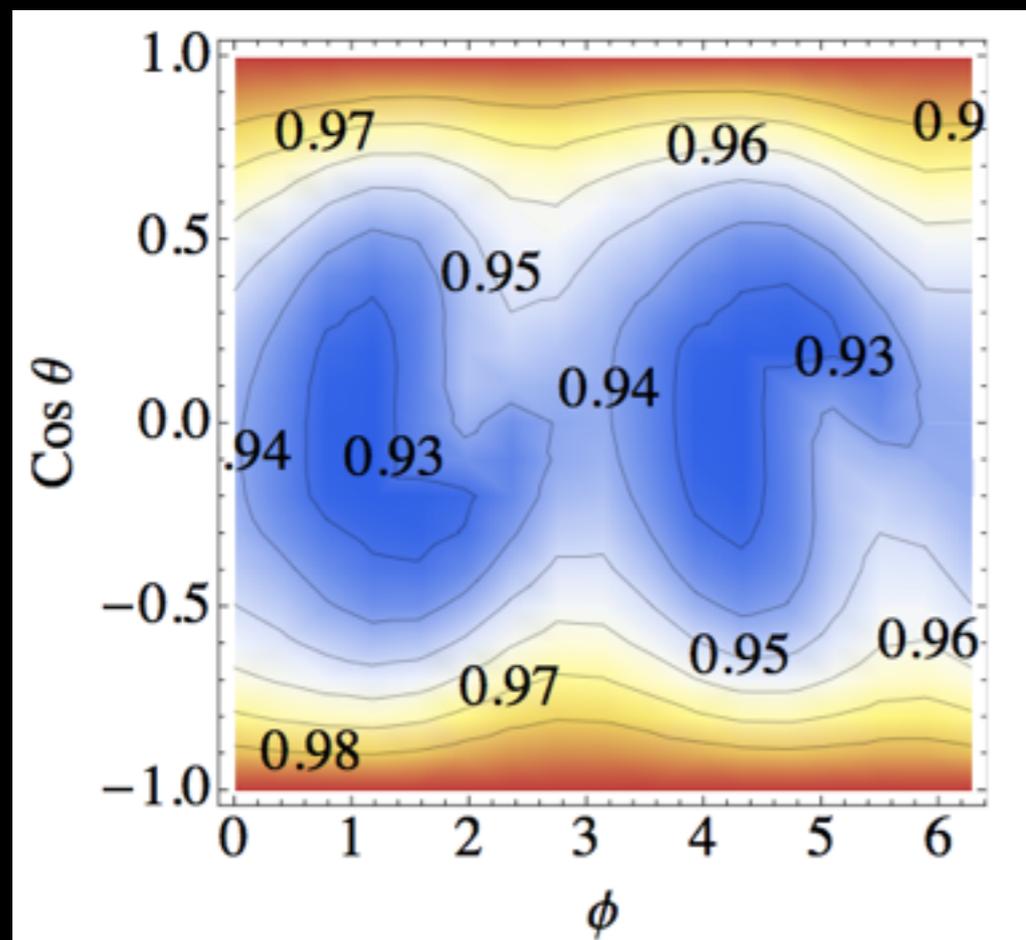
For non-precessing binaries spin effects dominated by only *one* key spin parameter

Does something similar happen for precession?  
i.e., can we replace  
the four in-plane spin components  
with *one* “precession spin”?

Yes!

# PhenomP

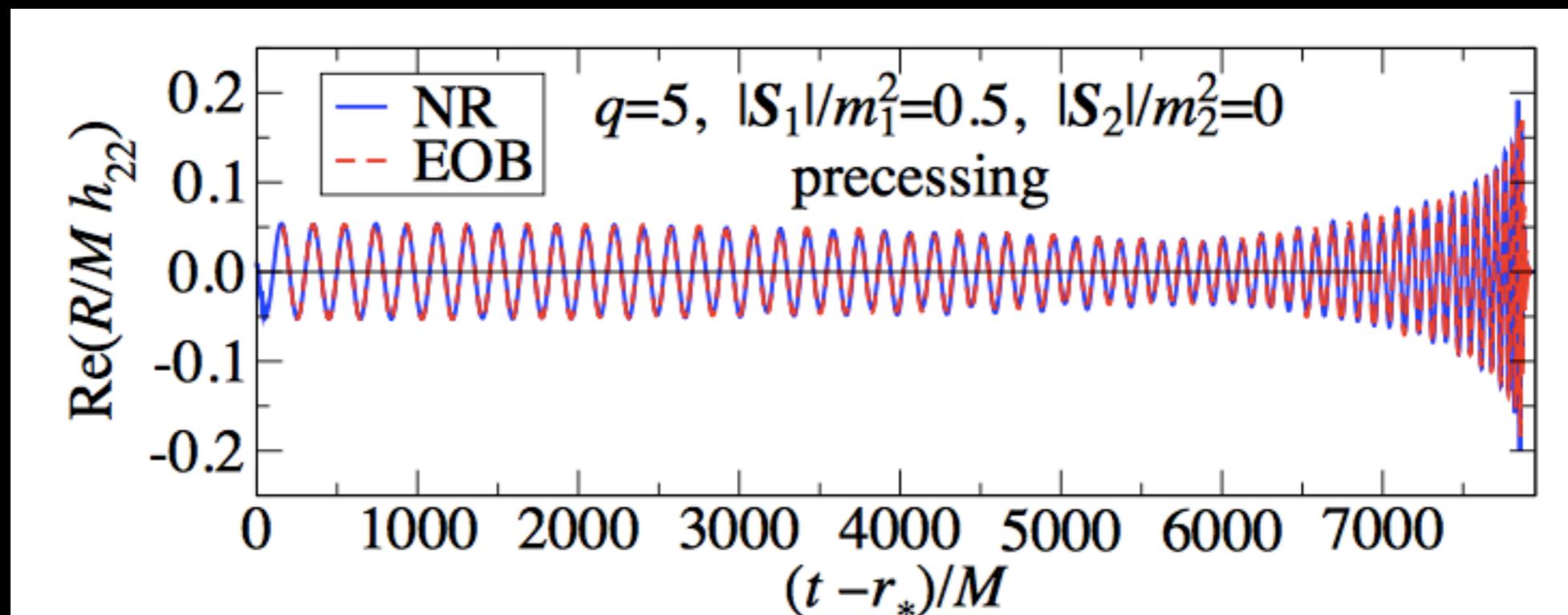
- Non-precessing model: PhenomD
- Twist with (analytic) PN precession angles
- Approximation: use PN angles through ringdown.



[Hannam, et al. (2014)]

# SEOBNRv3

- Non-precessing model: inspiral part of SEOBNRv2
- Twist with solution of precessing-EOB dynamics
- Attach ringdown
- Includes all 6 spin components



[Pan, et al. (2014), Taracchini, et. al. (2014)]

Neither precessing model  
is tuned to  
precessing NR simulations!

# Orientation dependence

$q=3$ ,  $|S_2| = 0.75$  (in plane)



Observer aligned  
with  $J$

# Orientation dependence

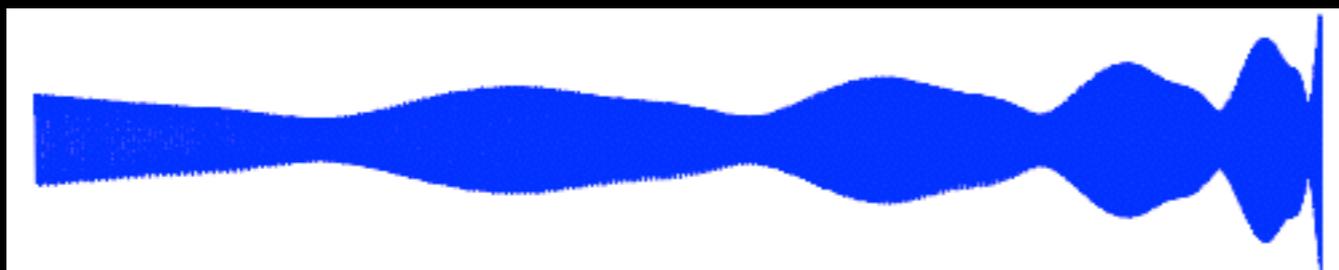
$q=3, |S_2| = 0.75$  (in plane)



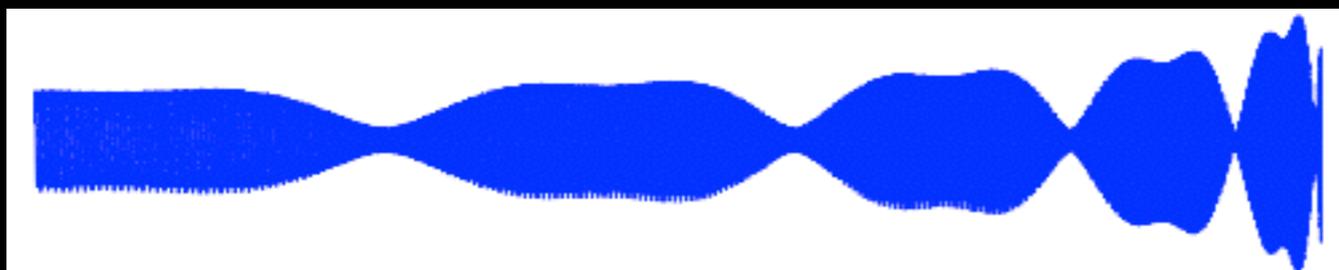
Observer aligned  
with  $J$



Observer inclined  
 $\pi/6$  to  $J$

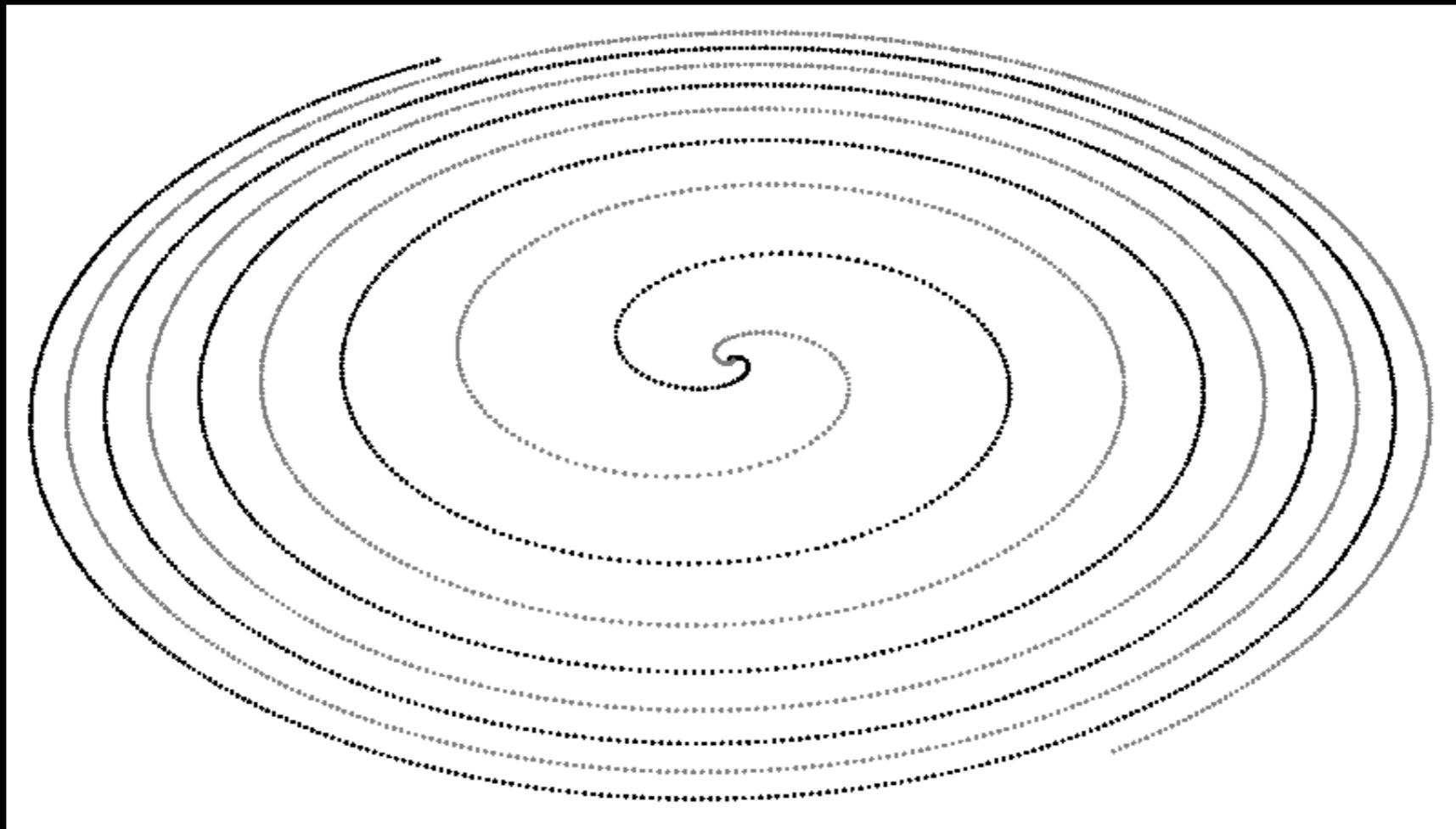


Observer inclined  
 $\pi/3$  to  $J$

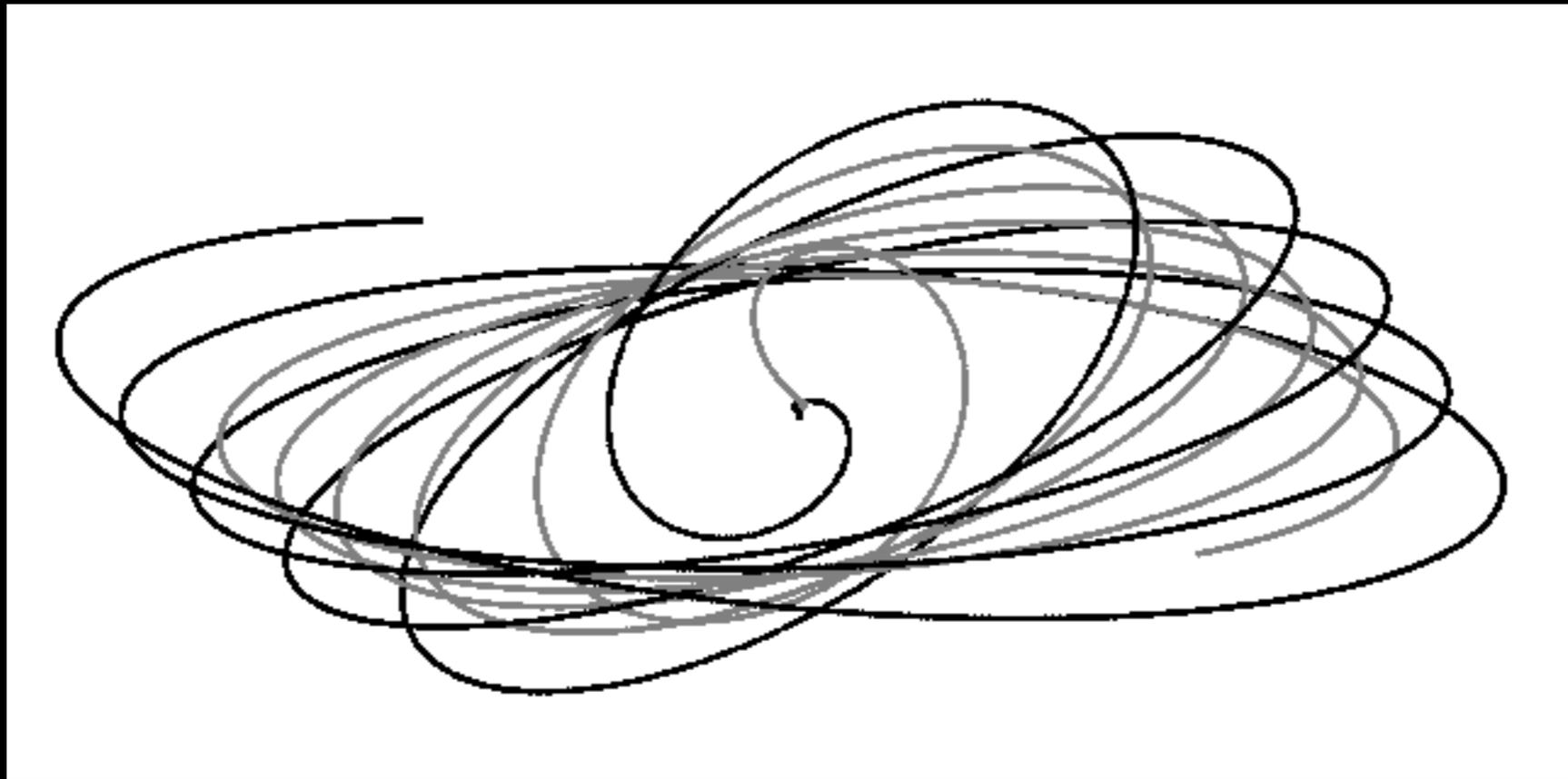


Observer inclined  
 $\pi/2$  to  $J$

# Equal-mass nonspinning BBH consistent with GW150914

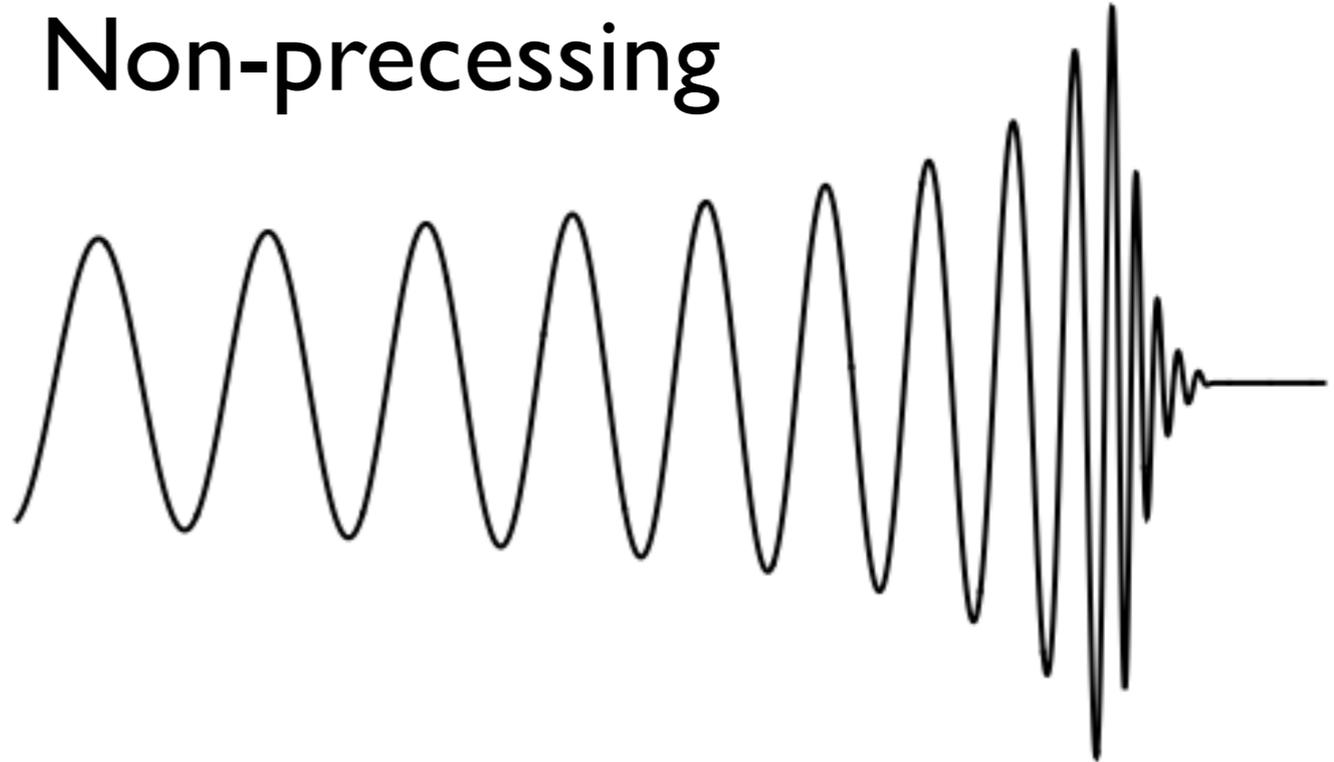


# Unequal-mass precessing BBH consistent with GW150914

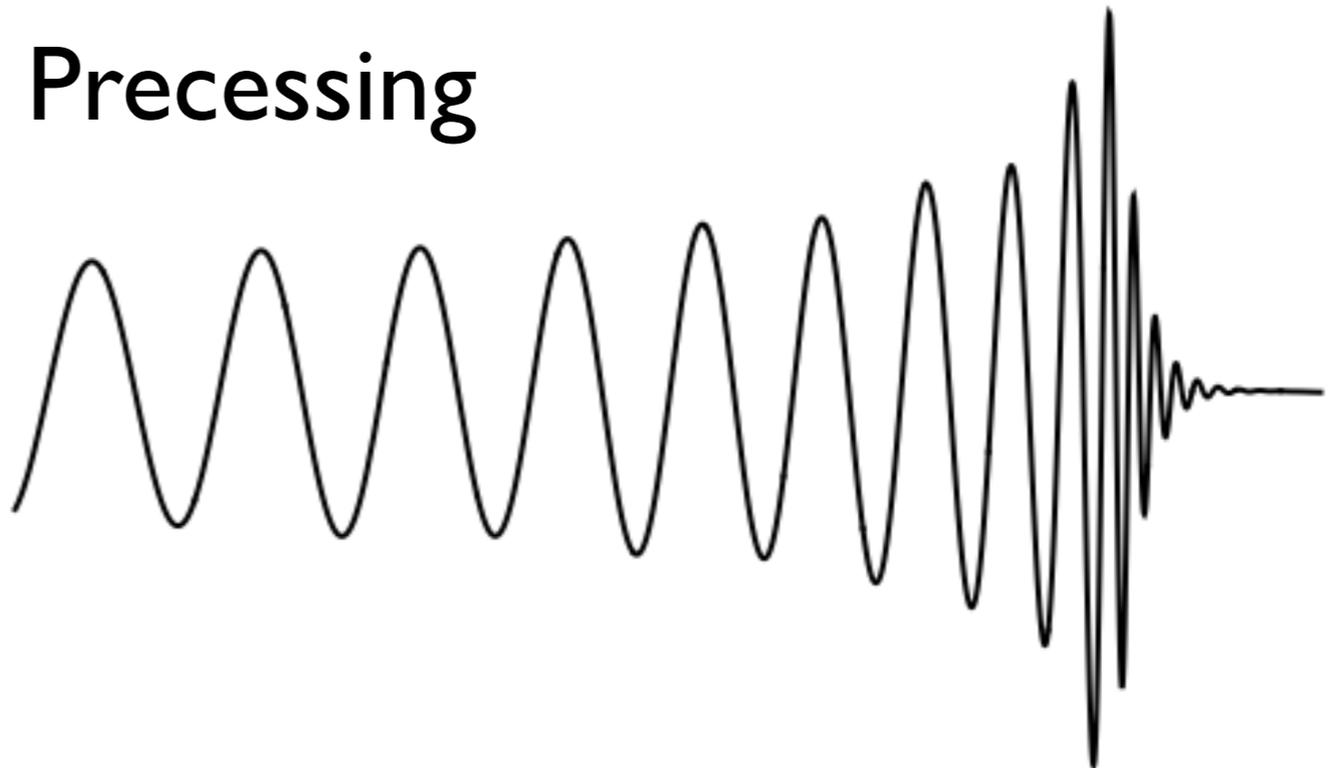


“Face-on”  
to the  
source

Non-precessing

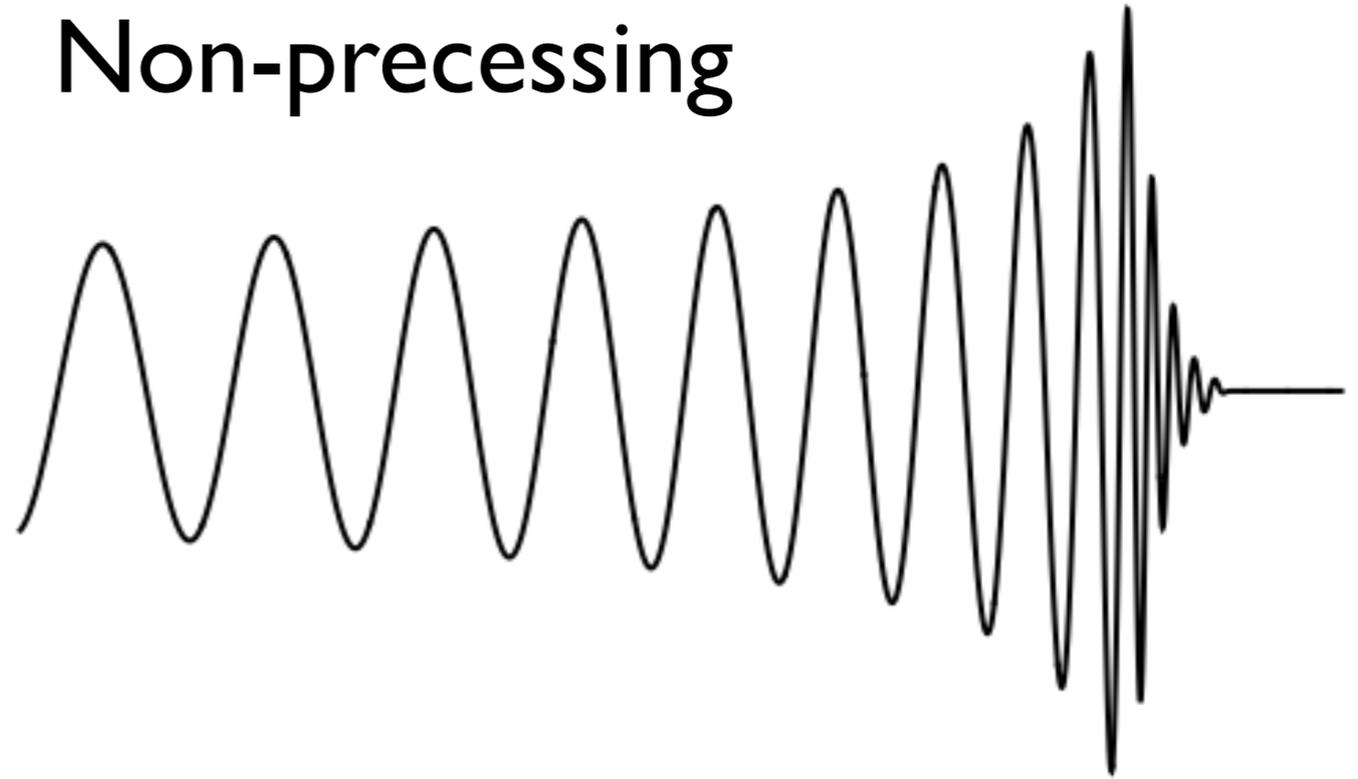


Precessing

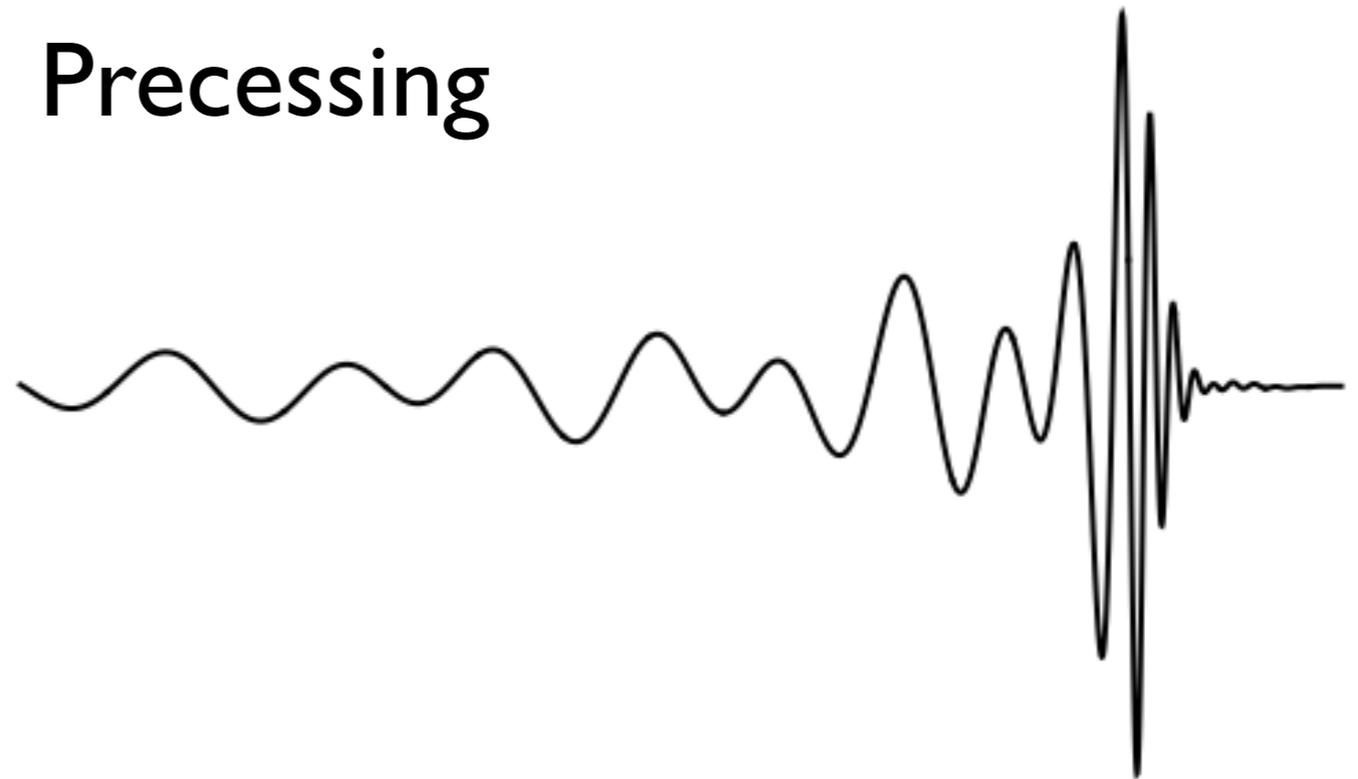


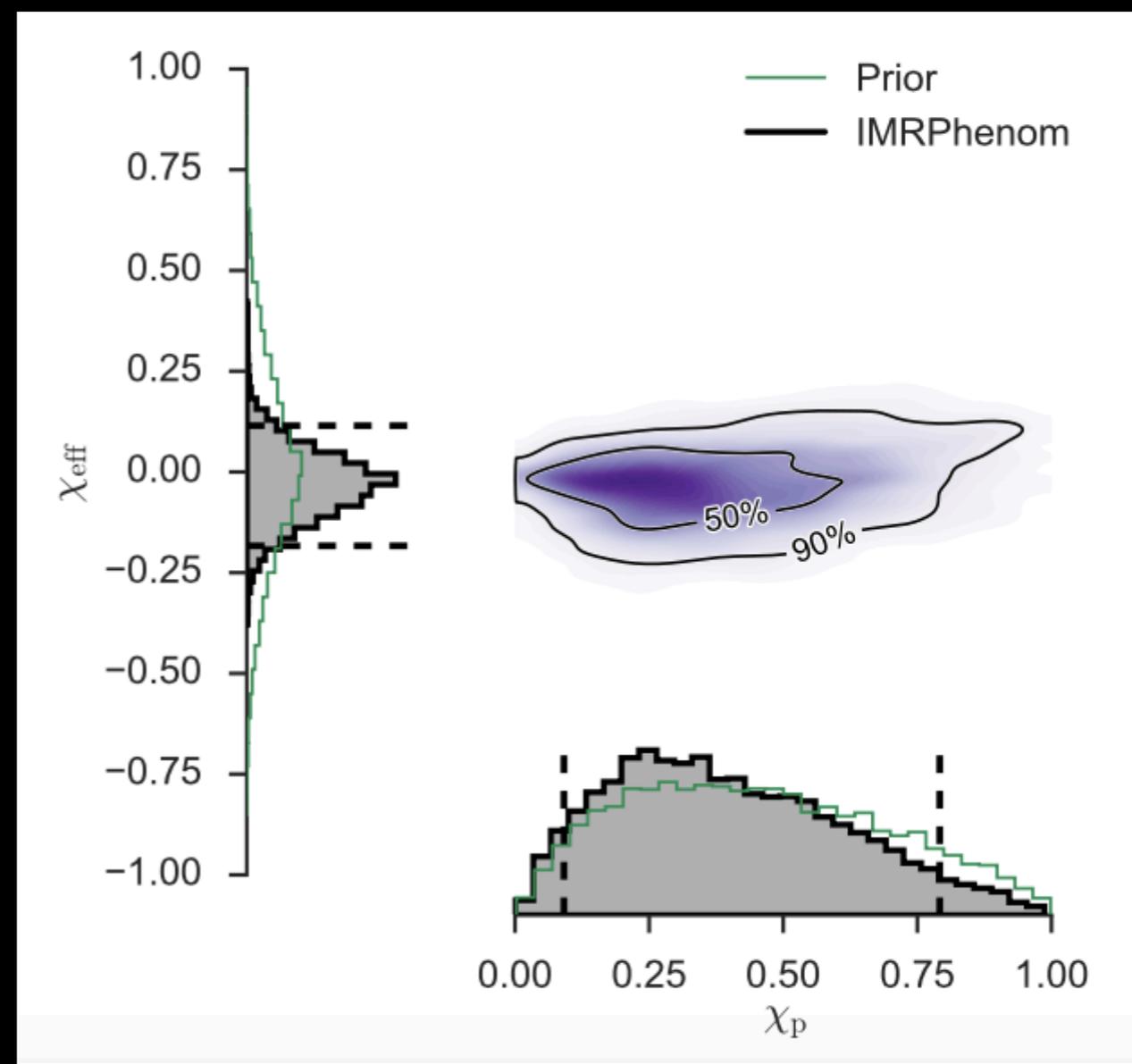
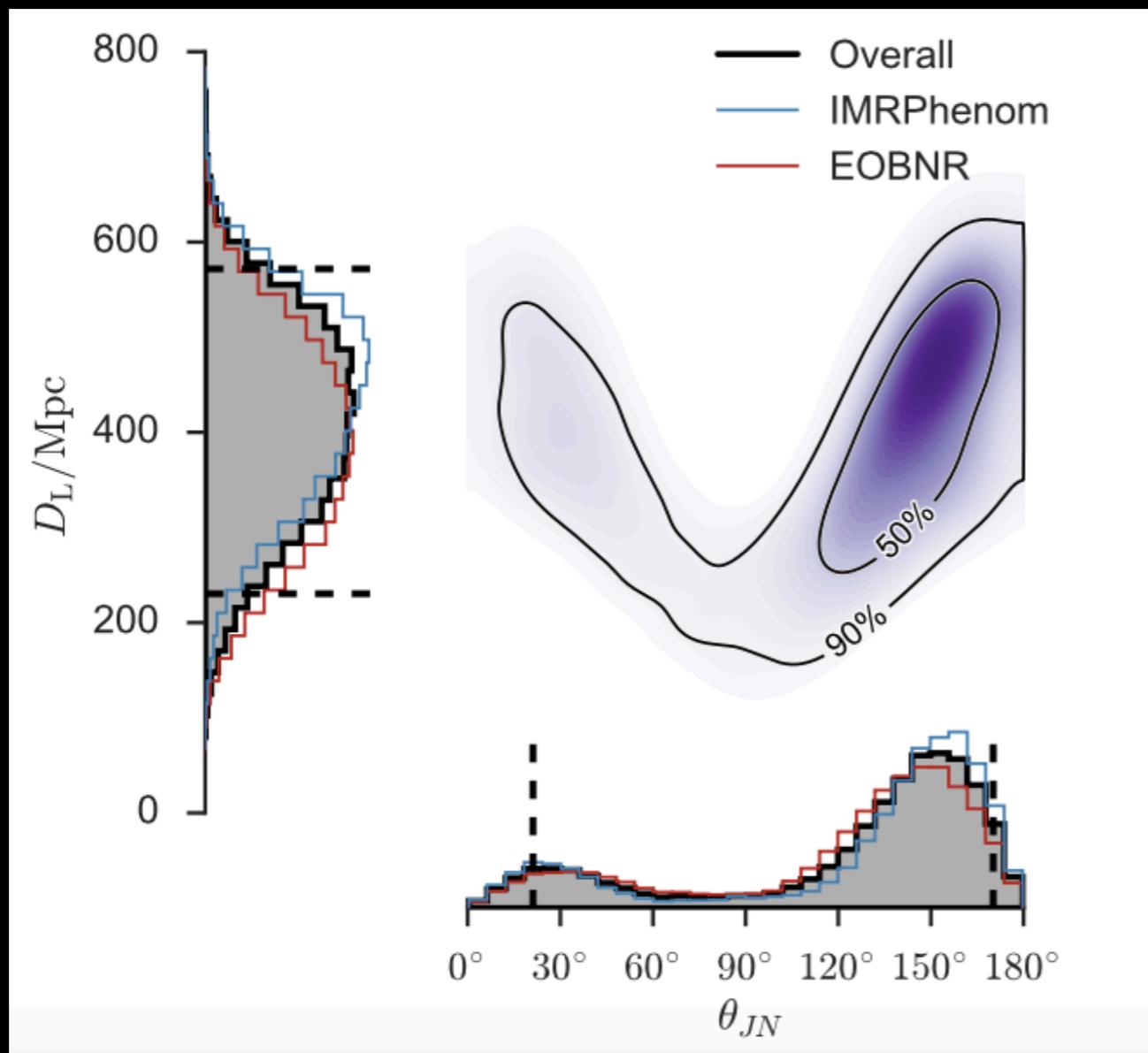
“Edge-on”  
to the  
source

Non-precessing

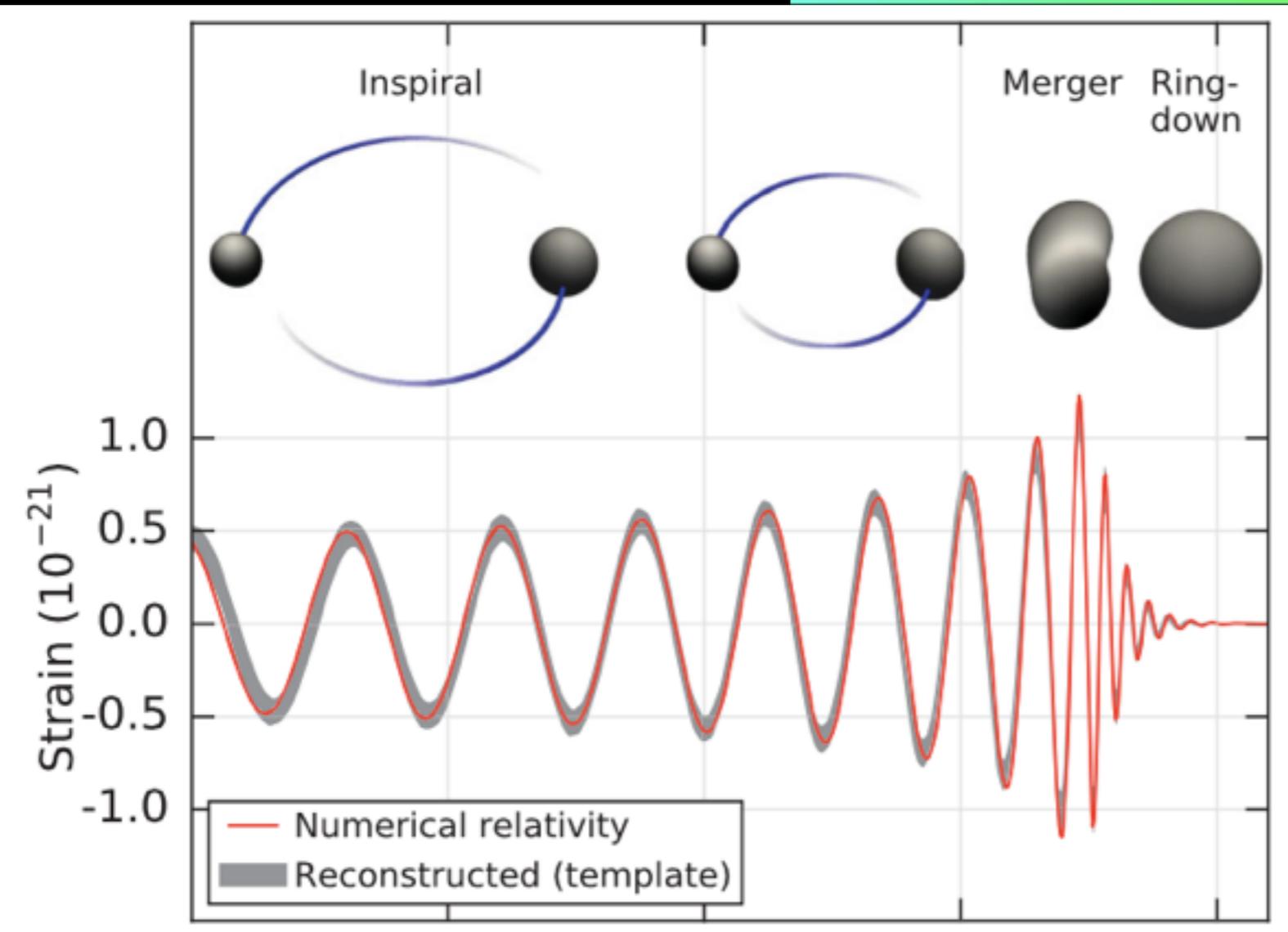
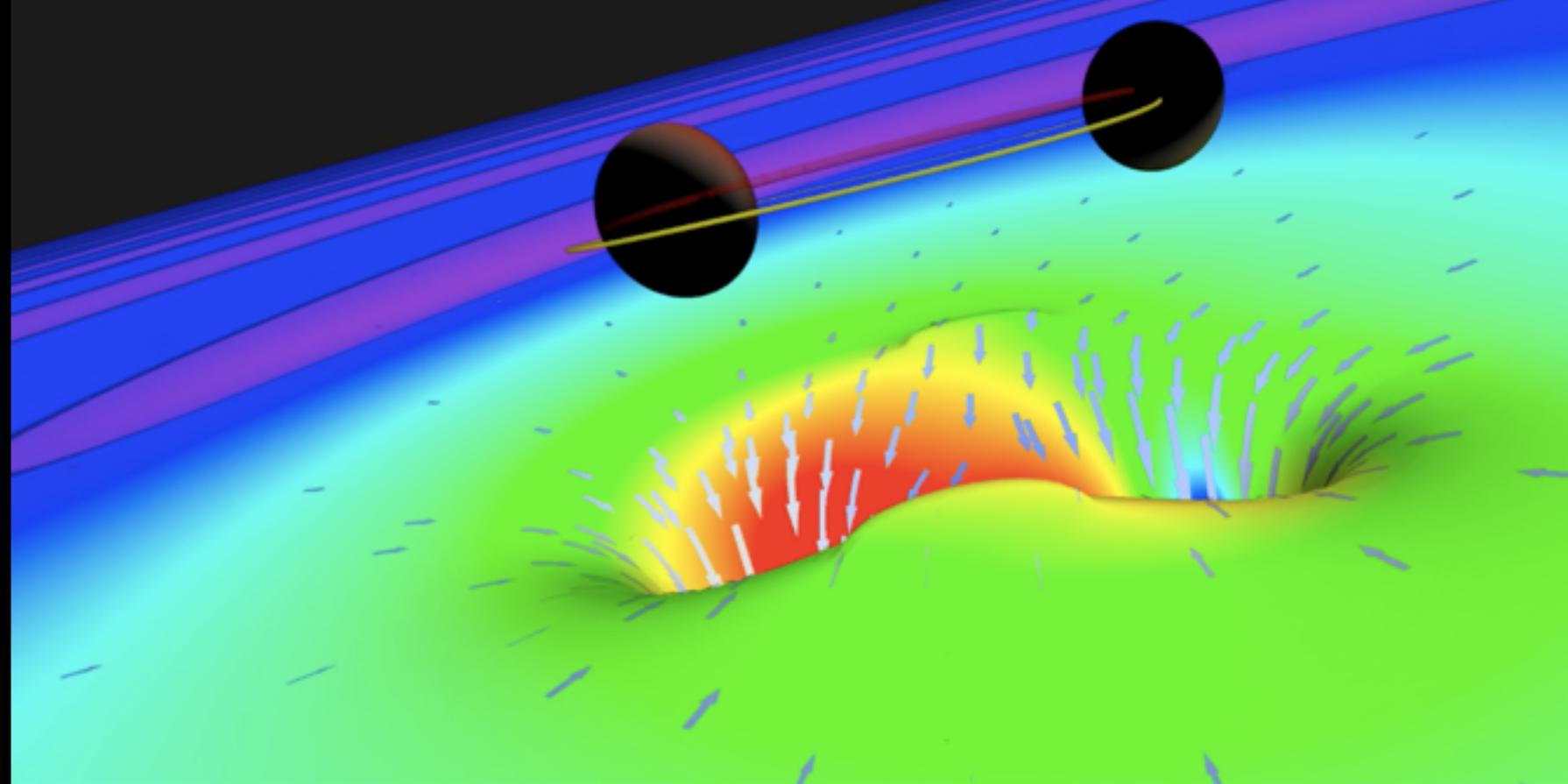


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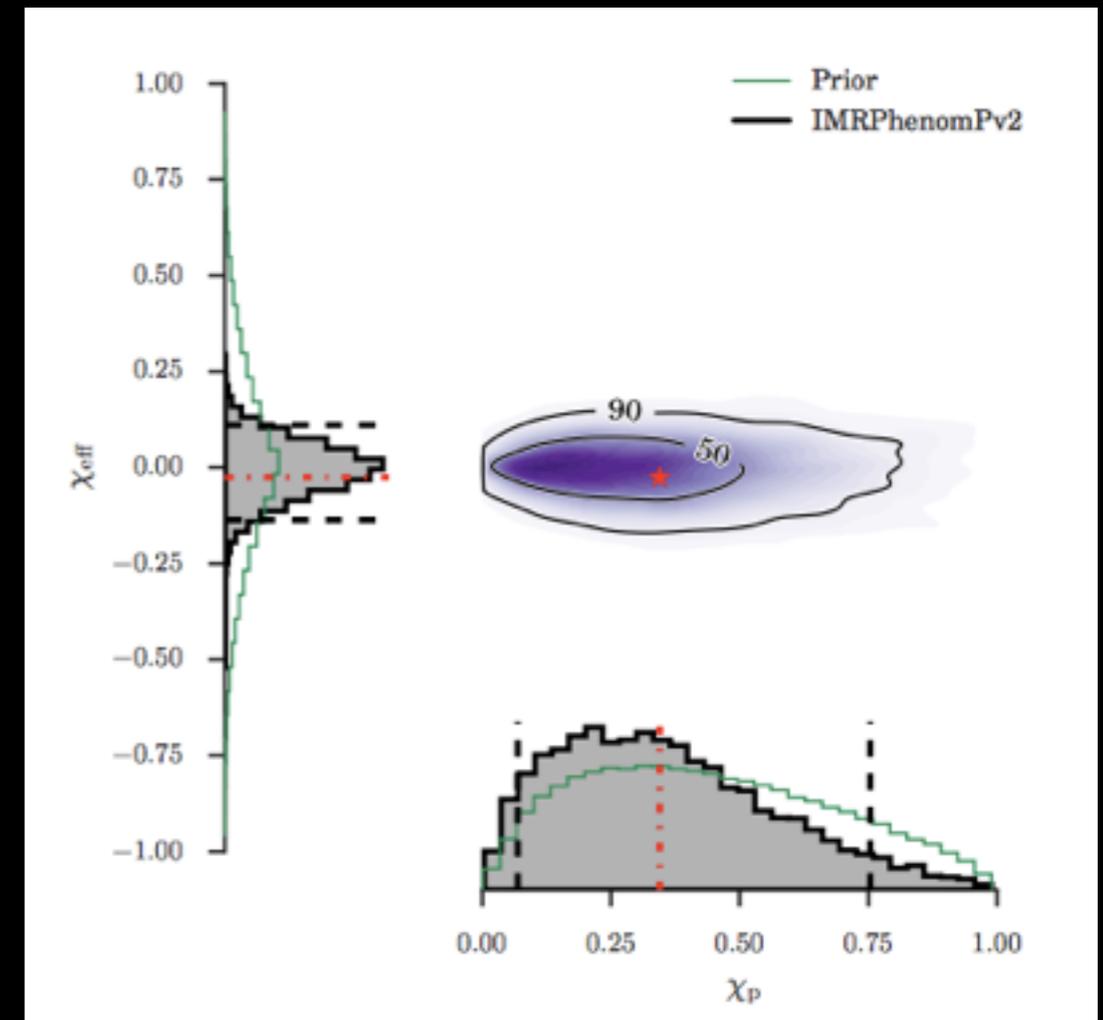
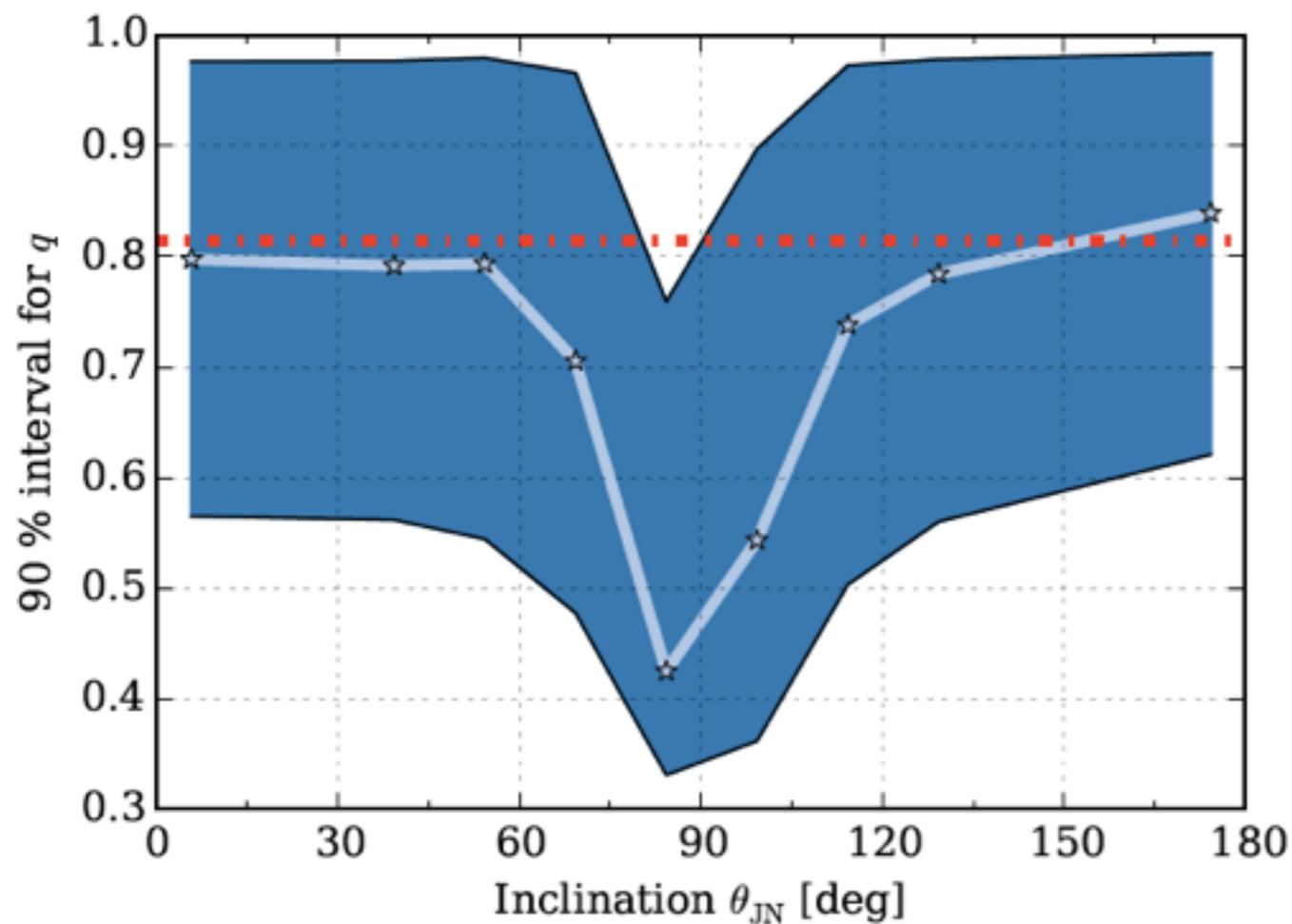




# NR simulations around GW150914



# Waveform model systematics



[Preliminary]

# Local modelling

- Several hundred NR simulations performed
- Cross-check of parameter estimates
- Reduced-order-quadrature model
- PhenomP tuned through merger

# Future observations

- SNR 25 at the accuracy limit of current models
- GW150914 was in the best-modelled region
- Better models need
  - Higher harmonics
  - Precession physics through merger
  - More accuracy (!)
- BUT: degeneracies will not evaporate!